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Tame majorant analyticity for the Birkhoff map of the defocusing Nonlinear Schrödinger equation on the circle

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Abstract

For the defocusing Nonlinear Schrödinger equation on the circle, we construct a Birkhoff map Φ which is tame majorant analytic in a neighborhood of the origin. Roughly speaking, majorant analytic means that replacing the coefficients of the Taylor expansion of Φ by their absolute values gives rise to a series (the majorant map) which is uniformly and absolutely convergent, at least in a small neighborhood. Tame majorant analytic means that the majorant map of Φ fulfills tame estimates.

The proof is based on a new tame version of the Kuksin-Perelman theorem [KP10], which is an infinite dimensional Vey type theorem.

1 Introduction and statement of the main result

1.1 Introduction

It is well known that the cubic defocusing Nonlinear Schrödinger equation (dNLS) on the circle

$$i\dot{\varphi} = -\partial_{xx}\varphi + 2|\varphi|^2\varphi, \quad x \in \mathbb{T} := \mathbb{R}/\mathbb{Z} \quad (1.1)$$

is an integrable system [ZS71, ZM74]. The actual construction of action-angle coordinates is quite complicated, and it has been studied analytically in the last decade by Grébert, Kappeler and collaborators in a series of works culminating in [GK14]. In particular these authors showed that there exists a globally defined map $\Phi : \varphi \mapsto (z_k, \bar{z}_k)_{k \in \mathbb{Z}}$, the *Birkhoff map*, which introduces *Birkhoff coordinates*, namely complex conjugates canonical coordinates $(z_k, \bar{z}_k)_{k \in \mathbb{Z}}$, with the property that the dNLS Hamiltonian, once expressed in such coordinates, is a real analytic function of the actions $I_k := |z_k|^2$ alone. As a consequence, in the Birkhoff coordinates the flow (1.1) is conjugated to an infinite chain of nonlinearly coupled oscillators:

$$i\dot{z}_k = \omega_k(I)z_k \quad \forall k \in \mathbb{Z}, \quad (1.2)$$

where the $\omega_k(I)$ are frequencies depending only on the actions $\{I_k\}_{k \in \mathbb{Z}}$.

Recently much effort has been made to understand various analytic properties of the Birkhoff map which are useful in applications. Such properties include for example the 1-smoothing of the nonlinear part of Φ [KP10, KSTa], two-sided polynomial estimates on the norm of Φ [Mo15], extension of Φ to spaces of low regularity [Mo16].

In this paper we contribute to such analysis by investigating the property of tame majorant analyticity of the Birkhoff map. Roughly speaking, an analytic map is majorant analytic if replacing the coefficients of its Taylor expansion by their absolute value gives rise to a series (the majorant map) which is uniformly and absolutely convergent, at least in a small neighborhood. Then tame majorant analytic means that the majorant map fulfills tame estimates.

Here we prove that this is indeed true for the Birkhoff map of dNLS, at least in a small neighborhood of the origin and in appropriate topologies. Our construction of the Birkhoff map is

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quite different and less explicit than the one of Grébert-Kappeler, however the two Birkhoff maps coincide (up to some normalization) as they are both perturbations of the Fourier transform.

While Grébert and Kappeler provide explicit and globally defined formulas for the action-angle coordinates of dNLS, our construction is valid only close to zero and it is based on the Kuksin-Perelman theorem [KP10], which is an infinite dimensional Vey-type theorem [Vey78, Eli90]. The Kuksin-Perelman theorem states that given a set of non-canonical coordinates, it is possible, under certain circumstances, to deform them into canonical Birkhoff coordinates. Therefore the main idea of our proof is to construct a starting set of non-canonical coordinates (essentially following the construction of Bättig, Grébert, Guillot and Kappeler [BGGK]), and then show that they fulfill the assumptions of the Kuksin-Perelman theorem, so that they can be deformed into Birkhoff coordinates. Actually we need a little bit more, since our aim is to construct *tame* Birkhoff coordinates. Therefore we prove that the starting coordinates are tame majorant analytic, and then we develop a tame version of the Kuksin-Perelman theorem, which guarantees that if the starting coordinates are tame-majorant analytic, so are the final Birkhoff coordinates. We think that the tame version of Kuksin-Perelman theorem could be interesting in itself.

Majorant analyticity and tameness of vector fields are properties extremely useful in perturbation theory and when one wants to apply Birkhoff normal form techniques. Indeed such properties are closed under composition, generation of flows and solution of homological equations, which are the typical operations needed in a perturbative scheme. This makes tame majorant analyticity an extremely robust tool when investigating stability of solutions. For example, majorant analyticity was used by Nikolenko [Nik86] to obtain Poincaré normal forms for some dissipative PDEs. Similarly, tame majorant analyticity was exploited by Bambusi and Grébert [BG06] to develop Birkhoff normal form theory for a wide class of Hamiltonian PDEs, and by Cong, Liu and Yuan [CLY] to study long time stability of small KAM tori of NLS with external potential; see also Berti, Biasco and Procesi [BBP13] for applications to KAM theory.

Concerning the usefulness of tame properties in perturbation theory, the idea is essentially the following. Tame estimates are estimates linear in the highest norm, a typical example being $\|u^n\|_{H^s} \leq C\|u\|_{H^s}\|u\|_{H^1}^{n-1}$; such estimates allow to control the size of a nonlinear term in a high regularity norm by conditions on the size of the function in a lower regularity norm. In [BG06] this property is exploited to show that, in the algorithm of Birkhoff normal form, large parts of the nonlinearity are actually very small in size and therefore harmless. A different applications of tame estimates is in differentiable Nash-Moser scheme, see e.g. [BBP10]; in this case the employ of tame estimates is one of the necessary ingredients for the convergence of the quadratic scheme. Also our interest in tame majorant analyticity of the Birkhoff map of dNLS was first motivated by applications: in the paper [MP17], we discuss the stability of small finite gap solutions of (1.1) when they are considered as solutions of the defocusing NLS on \mathbb{T}^2 . We first introduce Birkhoff coordinates and then perform a few steps of Birkhoff normal form. As in [BG06] this requires the *majorant analyticity* of the Hamiltonian.

Furthermore we think that properties of tame majorant analyticity of the Birkhoff map might be useful in the study of long time stability of perturbed dNLS.

1.2 Main result

As usual it is convenient to augment (1.1) with the conjugated equation for $\bar{\varphi}$ and to consider $(\varphi, \bar{\varphi})$ as independent variables belonging to the phase space $L_c^2 := L^2(\mathbb{T}, \mathbb{C}) \times L^2(\mathbb{T}, \mathbb{C})$ with elements $\varphi = (\varphi_1, \varphi_2)$. More generally we denote $H_c^s := H^s(\mathbb{T}, \mathbb{C}) \times H^s(\mathbb{T}, \mathbb{C})$ for any $s \in \mathbb{R}$. The dNLS Hamiltonian is given by

$$H_{\text{NLS}}(\varphi_1, \varphi_2) = \int_{\mathbb{T}} \left(\partial_x \varphi_1(x) \partial_x \varphi_2(x) + \varphi_1^2(x) \varphi_2^2(x) \right) dx$$

and the corresponding Hamiltonian system is

$$\begin{cases} i\dot{\varphi}_1 = \partial_{\varphi_2} H_{\text{NLS}} = -\partial_{xx} \varphi_1 + 2\varphi_2 \varphi_1^2 \\ i\dot{\varphi}_2 = -\partial_{\varphi_1} H_{\text{NLS}} = \partial_{xx} \varphi_2 - 2\varphi_1 \varphi_2^2 \end{cases} . \quad (1.3)$$

In such a way equation (1.1) is obtained by restricting the system above to the *real* invariant subspace

$$H_r^s := \{(\varphi_1, \varphi_2) \in H_c^s : \varphi_2 = \bar{\varphi}_1\} \quad (1.4)$$

of states of *real* type. We denote $L_r^2 := H_r^0$ and in such space we introduce the *real* scalar product $\langle \cdot, \cdot \rangle$ and the symplectic form Ω_0 defined for $\varphi_1 \equiv (\varphi_1, \bar{\varphi}_1)$ and $\varphi_2 \equiv (\varphi_2, \bar{\varphi}_2)$ by

$$\langle \varphi_1, \varphi_2 \rangle := 2 \operatorname{Re} \int_{\mathbb{T}} \varphi_1(x) \bar{\varphi}_2(x) dx, \quad \Omega_0(\varphi_1, \varphi_2) := \langle E \varphi_1, \varphi_2 \rangle, \quad (1.5)$$

where $E := i$.

It is useful to identify functions (φ_1, φ_2) with their Fourier coefficients. Thus we denote by $\mathcal{F} : L^2(\mathbb{T}, \mathbb{C}) \rightarrow \ell^2(\mathbb{Z}, \mathbb{C})$ the Fourier transform and associate to φ_1 its sequence of Fourier coefficients $\{\xi_j\}_{j \in \mathbb{Z}} = \mathcal{F}(\varphi_1)$ and to φ_2 its sequence of Fourier coefficients $\{\eta_{-j}\}_{j \in \mathbb{Z}} = \mathcal{F}(\varphi_2)$:

$$\varphi_1(x) = \sum_{j \in \mathbb{Z}} \xi_j e^{-i2j\pi x}, \quad \varphi_2(x) = \sum_{j \in \mathbb{Z}} \eta_j e^{i2j\pi x}. \quad (1.6)$$

Clearly

$$(\varphi_1, \varphi_2) \in L_c^2 \iff (\xi, \eta) \in \ell_c^2 := \ell^2(\mathbb{Z}, \mathbb{C}) \times \ell^2(\mathbb{Z}, \mathbb{C}),$$

and

$$(\varphi_1, \varphi_2) \in L_r^2 \iff (\xi, \eta) \in \ell_r^2 := \{(\xi, \eta) \in \ell_c^2 : \bar{\xi}_j = \eta_j, \forall j \in \mathbb{Z}\}.$$

We endow ℓ_r^2 with the *real* scalar product $\langle \cdot, \cdot \rangle$ and symplectic form ω_0 defined for $\xi^1 \equiv (\xi^1, \bar{\xi}^1)$ and $\xi^2 \equiv (\xi^2, \bar{\xi}^2)$ by

$$\langle \xi^1, \xi^2 \rangle := 2 \operatorname{Re} \sum_{j \in \mathbb{Z}} \xi_j^1 \bar{\xi}_j^2, \quad \omega_0(\xi^1, \xi^2) := \langle E \xi^1, \xi^2 \rangle, \quad (1.7)$$

and one has $\omega_0 := (\mathcal{F}^{-1})^* \Omega_0$.

We are interested also in more general spaces which we now introduce. It is more convenient to define such spaces in term of the Fourier coefficients (ξ, η) of (φ_1, φ_2) . So for any real $1 \leq p \leq 2$, $s \geq 0$ define

$$\ell_c^{p,s} := \{(\xi, \eta) \in \ell_c^2 : \|(\xi, \eta)\|_{p,s} < \infty\}, \quad (1.8)$$

where $\|(\xi, \eta)\|_{p,s} := \|\xi\|_{p,s} + \|\eta\|_{p,s}$ and

$$\|\xi\|_{p,s} := \left(\sum_{j \in \mathbb{Z}} \langle j \rangle^{ps} |\xi_j|^p \right)^{1/p}; \quad (1.9)$$

here $\langle j \rangle := 1 + |j|$. Correspondingly $\ell_r^{p,s} := \{(\xi, \bar{\xi}) \in \ell_c^{p,s}\}$ with the induced norm. Note that when $s = 0$, then the norm (1.9) is simply the ℓ^p norm of the Fourier coefficients of (φ_1, φ_2) ; therefore the spaces $\ell_c^{p,s}$ can also be thought as *weighted Fourier Lebesgue spaces*. Furthermore

$$(\varphi_1, \varphi_2) \in H_c^s \iff (\xi, \eta) \in \ell_c^{2,s}, \quad (\varphi_1, \varphi_2) \in H_r^s \iff (\xi, \eta) \in \ell_r^{2,s}.$$

We denote by $B^{p,s}(\rho)$ the ball with center 0 and radius ρ in the topology of $\ell_c^{p,s}$, and by $B_r^{p,s}(\rho)$ the same ball in $\ell_r^{p,s}$. For $s = 0$, we will write simply $\ell_c^p \equiv \ell_c^{p,0}$ and $B^p(\rho) \equiv B^{p,0}(\rho)$.

In order to state our main theorem we need to introduce the concept of tame majorant analytic map more precisely. Given a $\rho > 0$, $1 \leq p \leq 2$, let $F : B_r^{p,s}(\rho) \rightarrow \ell_r^{p,s'}$ be a real analytic map in a neighborhood of the origin¹. Write $F(\xi, \bar{\xi}) = (F_j(\xi, \bar{\xi}), \overline{F_j(\xi, \bar{\xi})})_{j \in \mathbb{Z}}$ in components and expand each component $F_j(\xi, \eta)$ into its uniformly convergent Taylor series in a neighborhood of the origin:

$$F_j(\xi, \eta) = \sum_{|K|+|L| \geq 0} F_{KL}^j \xi^K \eta^L.$$

¹here real analytic in a neighborhood of the origin means that there exists an analytic map $\tilde{F} : B^{p,s}(\rho) \rightarrow \ell^{p,s'}$ (defined in a complex ball) which coincides with F on the real subspace $B^{p,s}(\rho) \cap \ell_r^{p,s} \equiv B_r^{p,s}$

Define

$$\underline{F}_j(\xi, \eta) := \sum_{|K|+|L| \geq 0} |F_{KL}^j| \xi^K \eta^L$$

and the majorant map \underline{F} component-wise by $\underline{F}(\xi, \bar{\xi}) = (\underline{F}_j(\xi, \bar{\xi}), \underline{F}_j(\bar{\xi}, \xi))_{j \in \mathbb{Z}}$. Then F will be said to be *majorant analytic* if $\exists \rho_* > 0$ s.t. \underline{F} defines a real analytic map in a neighborhood of the origin mapping $B_r^{p,s}(\rho_*) \rightarrow \ell_r^{p,s'}$.

Given $0 \leq s \leq s' \leq s''$, F will be said to be (p, s, s', s'') - *tame majorant analytic* if it is majorant analytic $B_r^{p,s}(\rho_*) \rightarrow \ell_r^{p,s'}$ and furthermore \underline{F} restricts to a real analytic map $B_r^{p,s}(\rho_*) \cap \ell_r^{p,s'} \rightarrow \ell_r^{p,s''}$ fulfilling

$$\sup_{\zeta \in B_r^{p,s}(\rho_*) \cap \ell_r^{p,s'}} \frac{\|\underline{F}(\zeta)\|_{p,s''}}{\|\zeta\|_{p,s'}} < \infty. \quad (1.10)$$

Note that, in the estimate (1.10), the supremum is taken over $B_r^{p,s}(\rho_*) \cap \ell_r^{p,s'}$, namely on all the elements of $\ell_r^{p,s'}$ which belong to a ball of fix radius in the weaker topology of $\ell_r^{p,s}$. As we will show below (see Remark 2.12), (1.10) implies that each polynomial of the Taylor expansion of F is tame in the sense of polynomial maps.

Our main theorem is the following one:

Theorem 1.1. *There exists $\rho_0 > 0$ and a real analytic and majorant analytic map $\Phi : B_r^2(\rho_0) \rightarrow \ell_r^2$ s.t. the following is true:*

- (i) Φ is canonical: $\Phi^* \omega_0 = \omega_0$.
- (ii) The map Φ is a perturbation of the identity; more precisely $d\Phi(0, 0) = \mathbb{1}$, with $\mathbb{1}$ the identity map.
- (iii) For any reals $1 \leq p \leq 2$, $s \geq 1$, $\exists 0 < \rho_s < \rho_0$ s.t. $\Phi - \mathbb{1}$ restricts to a $(p, 0, s, s)$ - tame majorant analytic map $B^p(\rho_s) \cap \ell_c^{p,s} \rightarrow \ell_c^{p,s}$. Moreover there exists $C > 0$, independent of s , s.t. $\forall 0 < \rho \leq \rho_s$ one has

$$\|\Phi - \mathbb{1}(\xi, \eta)\|_{p,s} \leq C 4^s \rho^2 \|(\xi, \eta)\|_{p,s}, \quad \forall (\xi, \eta) \in B^p(\rho) \cap \ell_c^{p,s}.$$

The same is true for $\Phi^{-1} - \mathbb{1}$, with different constants.

- (iv) Φ is a local Birkhoff map in the following sense: for any $(\xi, \bar{\xi}) \in B_r^2(\rho_1) \cap \ell_r^{2,1}$, define $(z_j, \bar{z}_j) := \Phi_j(\xi, \bar{\xi})$. Then the integrals of motion of dNLS are real analytic functions of the actions $I_j = |z_j|^2$. In particular, the Hamiltonian $H_{NLS}(\varphi) \equiv \int_{\mathbb{T}} |\partial_x \varphi(x)|^2 dx + \int_{\mathbb{T}} |\varphi(x)|^4 dx$, the mass $M(\varphi) := \int_{\mathbb{T}} |\varphi(x)|^2 dx$ and the momentum $P(\varphi) := \int_{\mathbb{T}} \bar{\varphi}(x) i \partial_x \varphi(x) dx$ have the form

$$(H_{NLS} \circ \mathcal{F}^{-1} \circ \Phi^{-1})(z, \bar{z}) = h_{nls}(\dots, I_{-1}, I_0, I_1, \dots), \quad (1.11)$$

$$(M \circ \mathcal{F}^{-1} \circ \Phi^{-1})(z, \bar{z}) = \sum_{j \in \mathbb{Z}} I_j, \quad (1.12)$$

$$(P \circ \mathcal{F}^{-1} \circ \Phi^{-1})(z, \bar{z}) = \sum_{j \in \mathbb{Z}} j I_j. \quad (1.13)$$

Finally $\Phi - \mathbb{1}$ is 1-smoothing, in the following sense:

- (v) For any reals $1 \leq p \leq 2$, $s \geq 1$, $\exists 0 < \rho'_s < \rho_0$ s.t. $\Phi - \mathbb{1}$ restricts to a $(p, 1, s, s+1)$ -tame majorant analytic map $B^{p,1}(\rho'_s) \cap \ell_c^{p,s} \rightarrow \ell_c^{p,s+1}$. Moreover there exists $C' > 0$, independent of s , s.t. $\forall 0 < \rho \leq \rho'_s$ one has

$$\|\Phi - \mathbb{1}(\xi, \eta)\|_{p,s+1} \leq C' 4^s \rho^2 \|(\xi, \eta)\|_{p,s}, \quad \forall (\xi, \eta) \in B^{p,1}(\rho) \cap \ell_c^{p,s}.$$

The same is true for $\Phi^{-1} - \mathbb{1}$, with different constants.

The main novelty of Theorem 1.1 are the tame majorant analytic properties of the Birkhoff map illustrated in item (iii) and (v). In particular item (iii) shows that Φ is convergent provided (ξ, η) are small in the low regularity space ℓ_c^2 , despite having large norm in higher regularity spaces. This turns out to be useful in applications (e.g. in perturbation theory), since in such a way one has typically to control only the low regularity norms of the solution. Finally item (v) shows that the nonlinear part of Φ is 1-smoothing provided the variables (ξ, η) are chosen at least in $\ell_c^{p,1}$. In such a way, one recovers (in a neighborhood of the origin), the 1-smoothing property of the Birkhoff map proved in [KSTa].

We are actually able to prove convergence of the Birkhoff map in spaces more general than $\ell_c^{p,s}$; for example we are able to deal with weighted Fourier Lebesgue spaces where the weight $\langle j \rangle^s$ in (1.9) is replaced by a more general weight w , e.g. by an analytic weight of the form

$$w_j := \langle j \rangle^s e^{a|j|}, \quad a > 0, \quad j \in \mathbb{Z}; \quad (1.14)$$

in such a case the norm is defined by $\|\xi\|_{p,s,a} := \left(\sum_{j \in \mathbb{Z}} \langle j \rangle^{ps} e^{pa|j|} |\xi_j|^p \right)^{1/p}$ and the space by $\ell_c^{p,s,a} := \{(\xi, \eta) \in \ell_c^2 : \|\xi\|_{p,s,a} + \|\eta\|_{p,s,a} < \infty\}$. Then we have the following theorem

Theorem 1.2. *With ρ_0 as in Theorem 1.1, $\forall 1 \leq p \leq 2$, $s, a \geq 0$, the map $\Phi - \mathbb{1}$ of Theorem 1.1 restricts to a majorant analytic map $B^{p,s,a}(\rho_0) \rightarrow \ell_c^{p,s,a}$. Moreover there exists $C > 0$, independent of s, a , s.t. $\forall 0 < \rho \leq \rho_0$ one has*

$$\sup_{\|(\xi, \eta)\|_{p,s,a} \leq \rho} \|\Phi - \mathbb{1}(\xi, \eta)\|_{p,s,a} \leq C\rho^3.$$

The same is true for $\Phi^{-1} - \mathbb{1}$, with different constants.

In this case we prove just majorant analyticity (and not tameness), but in spaces of analytic functions. Note that the domain of (majorant) analyticity of the Birkhoff map does not shrink to 0 as s, a go to infinity. This is a consequence of an explicit control of every constant in the proof of the quantitative Kuksin-Perelman theorem 2.19. Finally we mention that we are able to treat even more general weighted Fourier Lebesgue spaces, giving sufficient conditions for the weight, see Section 3.

An immediate corollary of Theorem 1.2 concerns the dNLS dynamics in $\ell_r^{p,s,a}$. Recall that the Cauchy problem for (1.1) is well posed in L_r^2 [Bou]. In Birkhoff coordinates, the flow of (1.1) is given by

$$(z_j(t), \bar{z}_j(t)) = \left(e^{-i\omega_j(I)t} z_j(0), e^{i\omega_j(I)t} \bar{z}_j(0) \right), \quad \forall j \in \mathbb{Z}, \quad (1.15)$$

where $\omega_j := \partial_{I_j} H_{NLS} \circ \Phi^{-1}$ is the j^{th} frequency, which depends only on the actions $(I_k)_{k \in \mathbb{Z}}$. Then in the original Fourier coordinates $\xi = \mathcal{F}(\varphi)$, provided $\xi(0)$ is small enough to belong to the domain of the Birkhoff map, one has $(\xi(t), \bar{\xi}(t)) = \Phi^{-1}(z(t), \bar{z}(t))$, where $z(t) := (z_j(t))_{j \in \mathbb{Z}}$. Since the norm $\|\cdot\|_{p,s,a}$ is invariant by the dynamics (1.15), one gets the following result:

Corollary 1.3. *There exist constants $\rho_*, C_* > 0$ s.t. for any $1 \leq p \leq 2$, $s, a \geq 0$ the following holds true. Consider the solution $\xi(t) = \mathcal{F}(\varphi(t))$ of (1.1) corresponding to initial data $\xi_0 = \mathcal{F}(\varphi_0) \in B_r^{p,s,a}(\rho)$, $\rho \leq \rho_*$. Then one has*

$$\sup_{t \in \mathbb{R}} \|\xi(t)\|_{p,s,a} \leq \rho(1 + C_*\rho^2). \quad (1.16)$$

Note that in Corollary 1.3 there is no loss of analyticity of the solution (as it happens in [KP09] for example), in the sense that exponential decay of the initial datum is preserved by the flow. The point is that we work only with small initial datum, for which we know that the Birkhoff map and its inverse map $B^{p,s,a}(\rho) \rightarrow \ell_r^{p,s,a}$ with the same a .

Before closing this introduction, we recall some previous works on analytic properties of the Birkhoff map of infinite dimensional integrable systems.

Concerning majorant analyticity of the Birkhoff map, the first result was proved by Kuksin and Perelman in the case of KdV on \mathbb{T} [KP10]. In particular, these authors proved that in a small

neighborhood of the origin in H_r^s , $\forall s \geq 0$, the nonlinear part of the Birkhoff map is both majorant analytic and 1-smoothing. The techniques of this paper were extended by Bambusi and the author [BM16] in order to deal with the Toda lattice with N particles, N arbitrary large.

Later on, it was proved in [KST13] that the nonlinear part of the Birkhoff map of the KdV on \mathbb{T} is globally 1-smoothing, and the same is true also for the Birkhoff map of the dNLS on \mathbb{T} [KSTa] (see also [MS16] for the case of KdV on \mathbb{R}). However, none of these papers addresses the question of tameness.

Also the use of Fourier-Lebesgue spaces (namely spaces with norms like (1.9) with $p \neq 2$) is not new in this context; e.g. in [KMMT] and [Mo16] the Birkhoff map of KdV and of dNLS were extended to weighted Fourier-Lebesgue spaces in order to study analytic properties of the action-to-frequency map $I \mapsto \omega(I)$.

Concerning tameness properties, recently Kappeler and Montalto [KM16] constructed real analytic, canonical coordinates for the dNLS on \mathbb{T} , which are defined in neighborhoods of families of finite dimensional invariant tori, and which satisfy tame estimates. However such coordinates are not Birkhoff coordinates, and the dNLS Hamiltonian, once expressed in such coordinates, is in normal form only up to order three. On the contrary, our coordinates are well defined only in a neighborhood of the origin, but the dNLS Hamiltonian, written in such coordinates, is in normal form at every order.

Finally we want to comment on Corollary 1.3, which shows that weighted Fourier-Lebesgue norms of the solution are uniformly bounded in time. As a consequence, there is no growth of Sobolev norms. The problem of giving upper bounds of the form (1.16) has been widely studied both for linear time dependent and nonlinear Schrödinger equations (see e.g. [MR17, BGMR2] for the linear case, [Soh97, PTV] for the nonlinear one and references therein). In case of linear Schrödinger equations quasiperiodic in time, $i\dot{\psi} = -\Delta\psi + V(\omega t, x)\psi$, it is known that the Sobolev norms of the solution can be uniformly bounded, provided the frequency vector ω is well chosen, see e.g. [EK09] for bounded perturbations on \mathbb{T}^d (see also [BGMR1] for some special perturbations on \mathbb{R}^d , and reference therein).

In case of dNLS, the inequality (1.16) is well known for data in H_r^s , $s \in \mathbb{N}$, and can be proved using the conservation laws of the dNLS hierarchy. In H_r^s , $s > 1$ real, inequality (1.16) is proved in [KSTb]. The novelty of inequality (1.16) is to treat the case $1 \leq p < 2$ and weighted spaces. We point out that the uniform bound (1.16) is not a mere consequence of integrability, but of the stronger property that the Birkhoff map preserves the topology, see Theorem 1.1. Indeed Gérard and Grellier proved that the cubic Szegő equation on \mathbb{T} is integrable [GG10, GG12], and nevertheless there are phenomenons of growth of Sobolev norms [GG17].

1.3 Scheme of the proof

In order to prove Theorem 1.1, we apply a tame version of the Kuksin-Perelman theorem [KP10] to the dNLS equation. We recall that the starting point of the Kuksin-Perelman theorem is to construct a map $\zeta \mapsto \Psi(\zeta)$ (not symplectic and locally defined), s.t. the quantities $|\Psi_j(\zeta)|^2$ are in involution, the level sets $|\Psi_j(\zeta)|^2 = c_j$ give a foliation in invariant tori, and Ψ and $d\Psi^*$ are majorant analytic maps. Then Kuksin and Perelman [KP10] showed that it is possible to deform Ψ into a new map Φ which is symplectic, majorant analytic and it is a Birkhoff map, in the sense that $(z, \bar{z}) := \Phi(\xi, \bar{\xi})$ are complex Birkhoff coordinates.

Therefore the first step of our proof is to prove a tame version of the Kuksin-Perelman theorem, which tells that if Ψ (namely the original map) is a tame majorant analytic map, so is the new map Φ . In order to prove such a theorem, we revisit the proof of the Kuksin-Perelman theorem (actually, of the quantitative version of the Kuksin-Perelman theorem proved in [BM16]), and prove that the algorithm of construction of Φ can be made tame, in the sense that at each step of the construction we can control quantities as in (1.10) for every object involved. This turns out to be true since the Kuksin-Perelman algorithm is based on a combination of some basic operations (like composition of functions, inversion of functions, generation of flows, and solution of a system of equations) which can be made tame.

Then the second step of our proof is to apply the tame Kuksin-Perelman theorem to the dNLS.

This amounts to construct the starting map $\zeta \mapsto \Psi(\zeta)$ and to prove that it fulfills the assumptions of the tame Kuksin-Perelman theorem (in particular, that Ψ is tame majorant analytic). Here we adapt to dNLS the ideas already employed in [KP10] for the KdV on \mathbb{T} and in [BM16] for the Toda lattice (see also the pioneering work of Bättig, Grébert, Guillot and Kappeler [BGGK]). The strategy is to construct Ψ by exploiting the integrable structure of dNLS, and in particular the Lax pair of dNLS. More precisely, starting from the spectral data of the Lax operator, one constructs perturbatively a map $\zeta \mapsto \Psi(\zeta)$ s.t. the quantities $|\Psi_j|^2$ equal the spectral gaps γ_j^2 , which are real analytic functions in involution. The main technical challenge is to show that the map Ψ is tame majorant analytic. This is proved by computing explicitly every polynomial in the Taylor expansion of Ψ , in order to have a precise formula for a majorant map.

The paper is structured in the following way: in Section 2 we recall the setup of weighted Sobolev spaces and state the tame Kuksin-Perelman theorem. Its proof is a variant of the proof written in [BM16], therefore we postpone it to Appendix A. In Section 3 we consider the dNLS equation and construct the map Ψ required by the tame Kuksin-Perelman theorem, and show that it is tame majorant analytic.

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2 The tame Kuksin-Perelman theorem

We prefer to work in the setting of abstract weighted Fourier-Lebesgue spaces, which we now recall. First we define *weight* a function $w : \mathbb{Z} \rightarrow \mathbb{R}$ such that $w_j > 0 \ \forall j \in \mathbb{Z}$. A weight will be said to be *symmetric* if $w_{-j} = w_j \ \forall j \in \mathbb{Z}$ and *sub-multiplicative* if $w_{j+i} \leq w_i w_j \ \forall i, j \in \mathbb{Z}$. Given two weights w and v , we will say that $v \leq w$ iff $v_j \leq w_j, \ \forall j \in \mathbb{Z}$. Given a weight w we define for any $\mathbb{R} \ni p \geq 1$ the space $\ell^{p,w}$ of complex sequences $\xi = \{\xi_j\}_{j \in \mathbb{Z}}$ with norm

$$\|\xi\|_{p,w} := \left(\sum_{j \in \mathbb{Z}} w_j^p |\xi_j|^p \right)^{1/p} < \infty. \quad (2.1)$$

We denote by $\ell_c^{p,w}$ the complex Banach space $\ell_c^{p,w} := \ell^{p,w} \oplus \ell^{p,w} \ni (\xi, \eta) \equiv \zeta$ endowed with the norm

$$\|\zeta\|_{p,w} \equiv \|(\xi, \eta)\|_{p,w} := \|\xi\|_{p,w} + \|\eta\|_{p,w}.$$

We denote by $\ell_r^{p,w}$ the *real* subspace of $\ell_c^{p,w}$ defined by

$$\ell_r^{p,w} := \{(\xi, \eta) \in \ell_c^{p,w} : \eta_j = \bar{\xi}_j \ \forall j \in \mathbb{Z}\}. \quad (2.2)$$

We endow such a space with the real scalar product and symplectic form (1.7). We will denote by $B^{p,w}(\rho)$ (respectively $B_r^{p,w}(\rho)$) the ball in the topology of $\ell_c^{p,w}$ (respectively $\ell_r^{p,w}$) with center 0 and radius $\rho > 0$. Clearly if $w_j = 1 \ \forall j$ one has $\ell_c^{p,w} \equiv \ell_c^p := \ell^p(\mathbb{Z}, \mathbb{C}) \times \ell^p(\mathbb{Z}, \mathbb{C})$. In this case we denote the norm simply by $\|\cdot\|_p$ and the ball of radius ρ by $B^p(\rho)$. Similarly we write $\ell_r^p \equiv \ell_r^{p,w}$ and $B_r^p(\rho) \equiv B_r^{p,w}(\rho)$.

As for any $1 \leq p \leq 2$ and weight $w > 0$ one has the inclusion $\ell_r^{p,w} \hookrightarrow \ell_r^2$, the scalar product and the symplectic form (1.7) are well defined on $\ell_r^{p,w}$ as well.

Remark 2.1. The space $\ell_c^{p,s}$ defined in (1.8) coincides with the weighted space $\ell_c^{p,w}$ choosing the weight $w = \{\langle j \rangle^s\}_{j \in \mathbb{Z}}$.

Given a smooth function $F : \ell_r^{p,w} \rightarrow \mathbb{R}$, we denote by X_F the Hamiltonian vector field of F , given by $X_F = J \nabla F$, where $J = E^{-1}$. For $F, G : \ell_r^{p,w} \rightarrow \mathbb{R}$ we denote by $\{F, G\}$ the Poisson bracket (with respect to ω_0): $\{F, G\} := \langle \nabla F, J \nabla G \rangle$ (provided it exists). We say that the functions

F, G commute if $\{F, G\} = 0$.

For \mathcal{X}, \mathcal{Y} Banach spaces, we shall write $\mathcal{L}(\mathcal{X}, \mathcal{Y})$ to denote the set of linear and bounded operators from \mathcal{X} to \mathcal{Y} . For $\mathcal{X} = \mathcal{Y}$ we will write just $\mathcal{L}(\mathcal{X})$.

Recall that a map $\tilde{P}^n : (\ell_c^{p,w})^n \rightarrow \mathcal{B}$, with \mathcal{B} a Banach space, is said to be *n-multilinear* if $\tilde{P}^n(\zeta^{(1)}, \dots, \zeta^{(n)})$ is \mathbb{C} -linear in each variable $\zeta^{(j)} \equiv (\xi^{(j)}, \eta^{(j)})$; a *n-multilinear map* is said to be *bounded* if there exists a constant $C > 0$ such that

$$\|\tilde{P}^n(\zeta^{(1)}, \dots, \zeta^{(n)})\|_{\mathcal{B}} \leq C \|\zeta^{(1)}\|_{p,w} \dots \|\zeta^{(n)}\|_{p,w}, \quad \forall \zeta^{(1)}, \dots, \zeta^{(n)} \in \ell_c^{p,w}.$$

Correspondingly its norm is defined by

$$\|\tilde{P}^n\| := \sup_{\|\zeta^{(1)}\|_{p,w}, \dots, \|\zeta^{(n)}\|_{p,w} \leq 1} \|\tilde{P}^n(\zeta^{(1)}, \dots, \zeta^{(n)})\|_{\mathcal{B}}.$$

A map $P^n : \ell_c^{p,w} \rightarrow \mathcal{B}$ is a *homogeneous polynomial* of order n if there exists a *n-multilinear map* $\tilde{P}^n : (\ell_c^{p,w})^n \rightarrow \mathcal{B}$ such that

$$P^n(\zeta) = \tilde{P}^n(\zeta, \dots, \zeta), \quad \forall \zeta \in \ell_c^{p,w}. \quad (2.3)$$

A *n-homogeneous polynomial* is bounded if it has finite norm

$$\|P^n\| := \sup_{\|\zeta\|_{p,w} \leq 1} \|P^n(\zeta)\|_{\mathcal{B}}.$$

Remark 2.2. Clearly $\|P^n\| \leq \|\tilde{P}^n\|$. Furthermore one has $\|\tilde{P}^n\| \leq e^n \|P^n\|$ – cf. [Muj86].

Remark 2.3. It is easy to see that a multilinear map and the corresponding polynomial are continuous (and analytic) if and only if they are bounded.

A map $F : \ell_c^{p,w} \rightarrow \mathcal{B}$ is said to be an *analytic germ* if there exists $\rho > 0$ such that $F : B^{p,w}(\rho) \rightarrow \mathcal{B}$ is analytic. Then F can be written as a power series absolutely and uniformly convergent in $B^{p,w}(\rho)$: $F(\zeta) = \sum_{n \geq 0} F^n(\zeta)$. Here $F^n(\zeta)$ is a homogeneous polynomial of degree n in the variables $\zeta = (\xi, \eta)$. We will write $F = O(\zeta^N)$ if in the previous expansion $F^n(\zeta) = 0$ for every $n < N$.

Let $U \subset \ell_r^{p,w}$ be open. A map $F : U \rightarrow \mathcal{B}$ is said to be a *real analytic germ* on U if for each point $(\xi, \bar{\xi}) \in U$ there exist a neighborhood V of $(\xi, \bar{\xi})$ in $\ell_c^{p,w}$ and an analytic germ which coincides with F on $U \cap V$.

Let now $F : U \subset \ell_c^{p,w^1} \rightarrow \ell_c^{p,w^2}$ be an analytic map. We will say that F is *real for real sequences* if $F(U \cap \ell_r^{p,w^1}) \subseteq \ell_r^{p,w^2}$, namely $F(\xi, \eta) = (F_1(\xi, \eta), F_2(\xi, \eta))$ satisfies $\overline{F_1(\xi, \bar{\xi})} = F_2(\xi, \bar{\xi})$. Clearly, the restriction $F|_{U \cap \ell_r^{p,w^1}}$ is a real analytic map.

2.1 Majorant analytic maps

Let $P^n : \ell_c^{p,w} \rightarrow \mathcal{B}$ be a homogeneous polynomial of order n ; assume \mathcal{B} separable and let $\{\mathbf{b}_m\}_{m \in \mathbb{Z}} \subset \mathcal{B}$ be a basis for the space \mathcal{B} . Expand P^n as follows

$$P^n(\zeta) \equiv P^n(\xi, \eta) = \sum_{\substack{|K|+|L|=n \\ m \in \mathbb{Z}}} P_{K,L}^{n,m} \xi^K \eta^L \mathbf{b}_m, \quad (2.4)$$

where $K, L \in \mathbb{N}_0^{\mathbb{Z}}$, $\mathbb{N}_0 = \mathbb{N} \cup \{0\}$, $|K| := \sum_{j \in \mathbb{Z}} K_j$, $\xi \equiv (\xi_j)_{j \in \mathbb{Z}}$ and $\xi^K := \prod_{j \in \mathbb{Z}} \xi_j^{K_j}$.

Definition 2.4. The modulus of a polynomial P^n is the polynomial \underline{P}^n defined by

$$\underline{P}^n(\xi, \eta) := \sum_{\substack{|K|+|L|=n \\ m \in \mathbb{Z}}} |P_{K,L}^{n,m}| \xi^K \eta^L \mathbf{b}_m. \quad (2.5)$$

A polynomial P^n is said to have bounded modulus if \underline{P}^n is a bounded polynomial.

We generalize now the notion of majorant analytic map given in the introduction.

Definition 2.5. An analytic germ $F : \ell_c^{p,w} \rightarrow \mathcal{B}$ is said to be majorant analytic if there exists $\rho > 0$ such that

$$\underline{F}(\zeta) := \sum_{n \geq 0} \frac{F^n(\zeta)}{n!} \quad (2.6)$$

is absolutely and uniformly convergent in $B^{p,w}(\rho)$. In such a case we will write $F \in \mathcal{N}_\rho(\ell_c^{p,w}, \mathcal{B})$. $\mathcal{N}_\rho(\ell_c^{p,w}, \mathcal{B})$ is a Banach space when endowed by the norm

$$|\underline{F}|_\rho := \sup_{\zeta \in B^{p,w}(\rho)} \|\underline{F}(\zeta)\|_{\mathcal{B}}. \quad (2.7)$$

A map $F : U \rightarrow \mathcal{B}$ is said to be real majorant analytic on U if for each point $(\xi, \bar{\xi}) \in U$ there exist a neighborhood V of $(\xi, \bar{\xi})$ in $\ell_c^{p,w}$ and a majorant analytic germ which coincides with F on $U \cap V$.

Remark 2.6. From Cauchy inequality one has that the Taylor polynomials F^r of F satisfy

$$\|F^r(\zeta)\|_{\mathcal{B}} \leq |\underline{F}|_\rho \frac{\|\zeta\|_{p,w}^r}{\rho^r}, \quad \forall \zeta \in B^{p,w}(\rho). \quad (2.8)$$

Remark 2.7. Since $\forall r \geq 1$ one has $\|F^r\| \leq \|\underline{F}^r\|$, if $F \in \mathcal{N}_\rho(\ell_c^{p,w}, \mathcal{B})$ then the Taylor series of F is uniformly convergent in $B^{p,w}(\rho)$.

We will often consider the case $\mathcal{B} = \ell_c^{p,w}$; in such a case the basis $\{\mathbf{b}_m\}_{m \in \mathbb{Z}}$ coincide with the natural basis $\mathbf{e}_{2m} := (e_m, 0)$, $\mathbf{e}_{2m+1} := (0, e_m)$ of such a space (where e_m is the vector in $\mathbb{C}^{\mathbb{Z}}$ with all components equal to zero except the m^{th} one which is equal to 1). We will consider also the case $\mathcal{B} = \mathcal{L}(\ell_c^{p,w^1}, \ell_c^{p,w^2})$ (bounded linear operators from ℓ_c^{p,w^1} to ℓ_c^{p,w^2}), where w^1 and w^2 are weights. Here the chosen basis is $\mathbf{b}_{jk} = \mathbf{e}_j \otimes \mathbf{e}_k$ (labeled by 2 indexes).

Remark 2.8. For $\zeta \equiv (\xi, \eta) \in \ell_c^{p,w}$, we denote by $|\zeta|$ the vector of the modulus of the components of ζ : $|\zeta| = (|\zeta_j|)_{j \in \mathbb{Z}}$, $|\zeta_j| := (|\xi_j|, |\eta_j|) \in \mathbb{R}^2$. If $F \in \mathcal{N}_\rho(\ell_c^{p,w}, \ell_c^{p,w})$ then for any $\zeta, v \in \ell_c^{p,w}$ one has

$$\underline{dF}(|\zeta|)|v| \leq dF(|\zeta|)|v|$$

(see [KP10]). Thus $\forall 0 < d < 1$, Cauchy estimates imply that $dF \in \mathcal{N}_{(1-d)\rho}(\ell_c^{p,w}, \mathcal{L}(\ell_c^{p,w}, \ell_c^{p,w}))$ with

$$|\underline{dF}|_{\rho(1-d)} \leq \frac{1}{d\rho} |\underline{F}|_\rho, \quad (2.9)$$

where \underline{dF} is computed with respect to the basis $\mathbf{e}_j \otimes \mathbf{e}_k$.

Following Kuksin-Perelman [KP10] we will need also a further property.

Definition 2.9. Let $\mathbb{R} \ni \rho > 0$ and $\mathbb{N} \ni N \geq 2$. A majorant analytic germ $F \in \mathcal{N}_\rho(\ell_c^{p,w}, \ell_c^{p,w})$ will be said to be of class $\mathcal{A}_{w,\rho}^N$ if $F = O(\zeta^N)$ and the map $\zeta \mapsto dF(\zeta)^* \in \mathcal{N}_\rho(\ell_c^{p,w}, \mathcal{L}(\ell_c^{p,w}, \ell_c^{p,w}))$. On $\mathcal{A}_{w,\rho}^N$ we will use the norm

$$\|F\|_{\mathcal{A}_{w,\rho}^N} := |\underline{F}|_\rho + \rho |\underline{dF}|_\rho + \rho |\underline{dF}^*|_\rho. \quad (2.10)$$

Remark 2.10. Assume that for some $\rho > 0$ the map $F \in \mathcal{A}_{w,\rho}^N$, $N \geq 2$, then for every $0 < d \leq \frac{1}{2}$ one has $|\underline{F}|_{d\rho} \leq 2d^N |\underline{F}|_\rho$ and $\|F\|_{\mathcal{A}_{w,d\rho}^N} \leq 6d^N \|F\|_{\mathcal{A}_{w,\rho}^N}$.

A real majorant analytic germ $F : B_r^{p,w}(\rho) \rightarrow \ell_r^{p,w}$ will be said to be of class $\mathcal{N}_\rho(\ell_c^{p,w}, \ell_r^{p,w})$ (respectively $\mathcal{A}_{w,\rho}^N$) if there exists a map of class $\mathcal{N}_\rho(\ell_c^{p,w}, \ell_c^{p,w})$ (respectively $\mathcal{A}_{w,\rho}^N$), which coincides with F on $B_r^{p,w}(\rho)$, namely on the restriction $\bar{\xi}_j = \eta_j$, $\forall j \in \mathbb{Z}$. In this case we will also denote by $|\underline{F}|_\rho$ (respectively $\|F\|_{\mathcal{A}_{w,\rho}^N}$) the norm defined by (2.7) (respectively (2.10)) of the complex extension of F .

2.2 Tame majorant analytic maps

We begin with the following definition

Definition 2.11. Fix $1 \leq p \leq 2$, weights $w^0 \leq w^1 \leq w^2$ and $F \in \mathcal{N}_\rho(\ell_c^{p,w^0}, \ell_c^{p,w^0})$. F is said to be (p, w^0, w^1, w^2) -tame majorant analytic if $\underline{F} : B^{p,w^0}(\rho) \cap \ell_c^{p,w^1} \rightarrow \ell_c^{p,w^2}$ is analytic and

$$|\underline{F}|_\rho^T := \sup \left\{ \frac{\|\underline{F}(\zeta)\|_{p,w^2}}{\|\zeta\|_{p,w^1}} : \zeta \in B^{p,w^0}(\rho) \cap \ell_c^{p,w^1} \right\} < \infty. \quad (2.11)$$

In such a case we will write $F \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \ell_c^{p,w^2})$. We endow such a space with the norm

$$\langle |\underline{F}| \rangle_\rho := |\underline{F}|_\rho + \rho |\underline{F}|_\rho^T, \quad (2.12)$$

where here $|\underline{F}|_\rho := \sup \{ \|\underline{F}(\zeta)\|_{p,w^0} : \|\zeta\|_{p,w^0} \leq \rho \}$.

Remark 2.12. Let $F \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \ell_c^{p,w^2})$. Expand F in Taylor series, $F = \sum_n F^n$. Then it follows by Cauchy estimates that each polynomial F^n is (p, w^0, w^1, w^2) -tame majorant analytic and

$$\|\underline{F}^r(\zeta)\|_{p,w^2} \leq \frac{|\underline{F}|_\rho^T}{\rho^{r-1}} \|\zeta\|_{p,w^0}^{r-1} \|\zeta\|_{p,w^1}. \quad (2.13)$$

Consequently, using also (2.8), one has $\langle |\underline{F}^r| \rangle_\rho \leq \langle |\underline{F}| \rangle_\rho$.

In case of maps with values in the space $\mathcal{L}(\ell_c^{p,w^1}, \ell_c^{p,w^2})$ we give the following

Definition 2.13. Fix $1 \leq p \leq 2$, weights $w^0 \leq w^1 \leq w^2$ and let $\mathcal{G} \in \mathcal{N}_\rho(\ell_c^{p,w^0}, \mathcal{L}(\ell_c^{p,w^0}, \ell_c^{p,w^0}))$. \mathcal{G} is said to be (p, w^0, w^1, w^2) -tame majorant analytic if $\underline{\mathcal{G}} : B^{p,w^0}(\rho) \cap \ell_c^{p,w^1} \rightarrow \mathcal{L}(\ell_c^{p,w^1}, \ell_c^{p,w^2})$ and

$$|\underline{\mathcal{G}}|_\rho^T := \sup \left\{ \frac{\|\underline{\mathcal{G}}(\zeta)v\|_{p,w^2}}{\|\zeta\|_{p,w^1} \|v\|_{p,w^0} + \rho \|v\|_{p,w^1}} : \zeta, v \in \ell_c^{p,w^1}, \|\zeta\|_{p,w^0} \leq \rho \right\} < \infty. \quad (2.14)$$

In such a case we will write $\mathcal{G} \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \mathcal{L}(\ell_c^{p,w^1}, \ell_c^{p,w^2}))$. We endow such a space with the norm

$$\langle |\underline{\mathcal{G}}| \rangle_\rho := |\underline{\mathcal{G}}|_\rho + \rho |\underline{\mathcal{G}}|_\rho^T, \quad (2.15)$$

where here $|\underline{\mathcal{G}}|_\rho := \sup \{ \|\underline{\mathcal{G}}(\zeta)\|_{\mathcal{L}(\ell_c^{p,w^0})} : \|\zeta\|_{p,w^0} \leq \rho \}$.

Remark 2.14. Let $F \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \ell_c^{p,w^2})$. By Cauchy formula one has

$$dF(\zeta)v = \frac{1}{2\pi i} \oint_{|\lambda|=\epsilon} \frac{F(\zeta + \lambda v)}{\lambda^2} d\lambda, \quad (2.16)$$

provided $|\zeta| + \epsilon|v| \in B^{p,w^0}(\rho)$. It follows that for any $0 < d \leq 1/2$, the map $dF \in \mathcal{N}_{(1-d)\rho}(B^{p,w^0} \cap \ell_c^{p,w^1}, \mathcal{L}(\ell_c^{p,w^1}, \ell_c^{p,w^2}))$ with

$$|\underline{dF}|_{\rho(1-d)}^T \leq \frac{1}{d\rho} |\underline{F}|_\rho^T. \quad (2.17)$$

We extend Definition 2.9 to deal with tame majorant analytic maps:

Definition 2.15. Let $\mathbb{R} \ni \rho > 0$ and $\mathbb{N} \ni N \geq 2$. A map $F \in \mathcal{A}_{w^0, \rho}^N$ will be said to be of class $\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}$ if $F \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \ell_c^{p,w^2})$ and the map $\zeta \mapsto dF(\zeta)^* \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \mathcal{L}(\ell_c^{p,w^1}, \ell_c^{p,w^2}))$. On $\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}$ we will use the norm

$$\|F\|_{\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}} := \langle |\underline{F}| \rangle_\rho + \rho \langle |\underline{dF}| \rangle_\rho + \rho \langle |\underline{dF}^*| \rangle_\rho. \quad (2.18)$$

Remark 2.16. Let $F \in \mathcal{T}_{w^0, w^1, \rho}^{w^2, N}$. Then

$$\|\underline{F}(\zeta)\|_{p, w^2} \leq \frac{\|F\|_{\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}}}{\rho} \|\zeta\|_{p, w^1}, \quad \forall \zeta \in B^{p, w^0}(\rho) \cap \ell_c^{p, w^1}.$$

Remark 2.17. Assume that for some $\rho > 0$ the map $F \in \mathcal{T}_{w^0, w^1, \rho}^{w^2, N}$, $N \geq 2$, then for every $0 < d \leq \frac{1}{2}$ one has $\langle |\underline{F}| \rangle_{d\rho} \leq 2d^N \langle |\underline{F}| \rangle_\rho$ and $\|F\|_{\mathcal{T}_{w^0, w^1, d\rho}^{w^2, N}} \leq 6d^N \|F\|_{\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}}$.

2.3 The tame Kuksin-Perelman theorem

We are now able to state a tame version of the Kuksin-Perelman theorem.

Fix $\rho > 0$, $1 \leq p \leq 2$ and let $\Psi : B_r^{p, w^0}(\rho) \rightarrow \ell_r^{p, w^0}$. Write Ψ component-wise, $\Psi = \{(\Psi_j, \bar{\Psi}_j)\}_{j \in \mathbb{Z}}$, and consider the foliation defined by the functions $\left\{ |\Psi_j(\xi, \bar{\xi})|^2 / 2 \right\}_{j \in \mathbb{Z}}$. Given $\xi \equiv (\xi, \bar{\xi}) \in \ell_r^{p, w^0}$ we define the leaf through ξ by

$$\mathcal{T}_\xi := \left\{ (v, \bar{v}) \in \ell_r^{p, w^0} : \frac{|\Psi_j(v, \bar{v})|^2}{2} = \frac{|\Psi_j(\xi, \bar{\xi})|^2}{2}, \quad \forall j \in \mathbb{Z} \right\}. \quad (2.18)$$

Let $\mathcal{T} = \bigcup_{\xi \in \ell_r^{p, w^0}} \mathcal{T}_\xi$ be the collection of all the leaves of the foliation. We denote by $T_\xi \mathcal{T}$ the tangent space to \mathcal{T}_ξ at the point $\xi \in \ell_r^{p, w^0}$. Next we define the function $I = \{I_j\}_{j \in \mathbb{Z}}$ by

$$I_j(\xi) \equiv I_j(\xi, \bar{\xi}) := \frac{|\xi_j|^2}{2} \quad \forall j \in \mathbb{Z}. \quad (2.19)$$

The foliation they define will be denoted by $\mathcal{T}^{(0)}$.

Remark 2.18. Ψ maps the foliation \mathcal{T} into the foliation $\mathcal{T}^{(0)}$, namely $\mathcal{T}^{(0)} = \Psi(\mathcal{T})$.

We state now the tame Kuksin-Perelman theorem:

Theorem 2.19 (Tame Kuksin-Perelman theorem). Let $1 \leq p \leq 2$ be real. Let w^0, w^1 and w^2 be weights with $w^0 \leq w^1 \leq w^2$. Consider the space ℓ_r^{p, w^0} endowed with the symplectic form ω_0 defined in (1.7). Let $\rho > 0$ and assume $\Psi : B_r^{p, w^0}(\rho) \rightarrow \ell_r^{p, w^0}$, $\Psi = \mathbb{1} + \Psi^0$ and $\Psi^0 \in \mathcal{T}_{w^0, w^1, \rho}^{w^2, N}$, $N \geq 2$. Define

$$\epsilon_1 := \|\Psi^0\|_{\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}}. \quad (2.20)$$

Assume that the functionals $\left\{ \frac{1}{2} |\Psi_j(\xi, \bar{\xi})|^2 \right\}_{j \in \mathbb{Z}}$ pairwise commute with respect to the symplectic form ω_0 , and that ρ is so small that

$$\epsilon_1 < 2^{-60} \rho. \quad (2.21)$$

Then there exists a real majorant analytic map $\tilde{\Psi} : B_r^{p, w^0}(\tilde{\mathbf{a}}\rho) \rightarrow \ell_r^{p, w^0}$, $\tilde{\mathbf{a}} = 2^{-120}$, with the following properties:

- i) $\tilde{\Psi}^* \omega_0 = \omega_0$, so that the coordinates $(z, \bar{z}) := \tilde{\Psi}(\xi, \bar{\xi})$ are canonical;
- ii) the functionals $\left\{ \frac{1}{2} |\tilde{\Psi}_j(\xi, \bar{\xi})|^2 \right\}_{j \in \mathbb{Z}}$ pairwise commute with respect to the symplectic form ω_0 ;
- iii) $\mathcal{T}^{(0)} = \tilde{\Psi}(\mathcal{T})$, namely the foliation defined by Ψ coincides with the foliation defined by $\tilde{\Psi}$;
- iv) $\tilde{\Psi} = \mathbb{1} + \tilde{\Psi}^0$ with $\tilde{\Psi}^0 \in \mathcal{T}_{w^0, w^1, \tilde{\mathbf{a}}\rho}^{w^2, N}$ and furthermore $\|\tilde{\Psi}^0\|_{\mathcal{T}_{w^0, w^1, \tilde{\mathbf{a}}\rho}^{w^2, N}} \leq 2^{17} \epsilon_1$.
- v) The inverse map $\tilde{\Psi}^{-1}$ is real majorant analytic $B_r^{p, w^0}(\tilde{\mathbf{a}}\rho) \rightarrow \ell_r^{p, w^0}$, $\tilde{\mathbf{a}} = 2^{-130}$, $\tilde{\Psi}^{-1} - \mathbb{1} \in \mathcal{T}_{w^0, w^1, \tilde{\mathbf{a}}\rho}^{w^2, N}$ with the quantitative estimate $\|\tilde{\Psi}^{-1} - \mathbb{1}\|_{\mathcal{T}_{w^0, w^1, \tilde{\mathbf{a}}\rho}^{w^2, N}} \leq 2^{18} \epsilon_1$.

Finally the theorem holds true also if the class $\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}$ is replaced by the class $\mathcal{A}_{w^0, \rho}^N$.

The novelty of Theorem 2.19 is to prove that the Birkhoff map is *tame majorant analytic*, provided the initial map $\Psi - \mathbb{1}$ is tame majorant analytic. The proof of Theorem 2.19 is actually a variant of the results of [KP10, BM16], so we postpone it to Appendix A.

The following corollary is an immediate application of Theorem 2.19 and shows that $\tilde{\Psi}$ is a Birkhoff map:

Corollary 2.20. *Let $H : \ell_r^{p,w^1} \rightarrow \mathbb{R}$ be a real analytic Hamiltonian function. Let Ψ be as in Theorem 2.19 and assume that for every $j \in \mathbb{Z}$, $|\Psi_j(\xi, \bar{\xi})|^2$ is an integral of motion for H , i.e.*

$$\{H, |\Psi_j|^2\} = 0 \quad \forall j \in \mathbb{Z} . \quad (2.22)$$

Then the coordinates (z_j, \bar{z}_j) defined by $(z_j, \bar{z}_j) := \tilde{\Psi}_j(\xi, \bar{\xi})$ are complex Birkhoff coordinates for H , namely canonical conjugated coordinates in which the Hamiltonian depends only on $|z_j|^2/2$.

Proof. By assumption, Ψ is analytic as a map $B^{p,w^0}(\rho) \cap \ell_c^{p,w^1} \rightarrow \ell_c^{p,w^1}$, therefore the composition $H \circ \Phi^{-1}$ is well defined and real analytic as a map $B_r^{p,w^0}(\rho) \cap \ell_r^{p,w^1} \rightarrow \mathbb{R}$. Thus it admits a convergent Taylor expansion of the form

$$(H \circ \Phi^{-1})(z, \bar{z}) = \sum_{\substack{r \geq 2 \\ |\alpha| + |\beta| = r}} H_{\alpha, \beta}^r z^\alpha \bar{z}^\beta . \quad (2.23)$$

Arguing as in [BM16, Corollary 2.13] one shows that (2.22) implies that in each monomial of the r.h.s. of (2.23), one has $\alpha = \beta$. \square

3 Application to dNLS on \mathbb{T}

In order to construct a tame Birkhoff map for dNLS, we wish to apply the Kuksin-Perelman theorem 2.19. This requires to be able to construct the starting map Ψ and to verify that such a map is tame majorant analytic. To construct Ψ we will exploit the integrable structure of dNLS, following the ideas already employed in the case of the KdV [KP10] and the Toda lattice [BM16] (see also [BGGK]). To prove tame majorant analyticity, we will expand Ψ in Taylor series $\Psi = \sum_{n \in \mathbb{N}} \Psi^n$, compute each polynomial Ψ^n and prove that it belongs to $\mathcal{S}_{w^0, w^1, \rho}^{w^2, n}$ for some weights $w^0 \leq w^1 \leq w^2$. We are able to state sufficient conditions for the choice of weights $w^0 \leq w^1 \leq w^2$ to use. Such conditions depend only on some arithmetic property that we state now. To do so, we need a bit of preparation.

Given $n \geq 3$ odd, $k_1, \dots, k_n \in \mathbb{Z}$, we define the function $\mathfrak{f}_n : \mathbb{Z}^n \times \mathbb{Z} \rightarrow \mathbb{R}$ by

$$\mathfrak{f}_n(k_1, \dots, k_n; j) := \mathbb{1}_{\{k_1 + \dots + k_n = j\}} \prod_{\substack{m=1 \\ m \text{ odd}}}^{n-1} \frac{1}{\langle k_1 + \dots + k_m - j \rangle} \cdot \frac{1}{\langle k_1 + \dots + k_{m+1} \rangle} , \quad (3.1)$$

where here $\mathbb{1}_S$ is the indicator function on the set S . Given an integer $1 \leq r \leq n$ we define $\mathfrak{g}_{n,r} : \mathbb{Z}^n \times \mathbb{Z} \rightarrow \mathbb{R}$ by

$$\mathfrak{g}_{n,r}(k_1, \dots, k_n; j) := \mathfrak{f}_n(k_1, \dots, k_{r-1}, j, k_{r+1}, \dots, k_n; k_r) . \quad (3.2)$$

Note that $\mathfrak{g}_{n,r}$ is \mathfrak{f}_n with indexes k_r and j switched. Its explicit expression is given in Appendix C.

The key point, as we shall see below, is that \mathfrak{f}_n bounds the kernel of the polynomial $\underline{\Psi}^n$, while the $\mathfrak{g}_{n,r}$'s bound the kernel of $[\mathbf{d}\underline{\Psi}^n]^*$. For example, it turns out that boundedness of $\underline{\Psi}^n$ and $[\mathbf{d}\underline{\Psi}^n]^*$ as maps $B^p(\rho) \rightarrow \ell_c^p$ respectively $B^p(\rho) \rightarrow \mathcal{L}(\ell_c^p)$ are implied by the following summability properties of \mathfrak{f}_n and $\mathfrak{g}_{n,r}$:

Lemma 3.1. *Let $1 \leq p \leq 2$ be real. Let p' s.t. $\frac{1}{p} + \frac{1}{p'} = 1$ and define*

$$R_* := \left(\sum_{k \in \mathbb{Z}} \frac{1}{\langle k \rangle^{p'}} \right)^{1/p'} .$$

Then for every $n \geq 3$, n odd

$$\sup_{j \in \mathbb{Z}} \|\mathbf{f}_n(\cdot; j)\|_{\ell^{p'}(\mathbb{Z}^n)} \leq R_*^{n-1}, \quad (3.3)$$

and

$$\max_{1 \leq r \leq n} \sup_{j \in \mathbb{Z}} \|\mathbf{g}_{n,r}(\cdot; j)\|_{\ell^{p'}(\mathbb{Z}^n)} \leq R_*^{n-1}. \quad (3.4)$$

The lemma is proved in Appendix C.1.

In a similar way, we will show below (see Lemma 3.10 and Lemma 3.11) that the maps Ψ^n and $[\mathbf{d}\Psi^n]^*$ are $(p, \mathbf{u}, \mathbf{v}, \mathbf{w})$ -tame majorant analytic if the weights $\mathbf{u} \leq \mathbf{v} \leq \mathbf{w}$ fulfill the following property:

(W)_p Let $\mathbf{u} \leq \mathbf{v} \leq \mathbf{w}$ be symmetric and sub-multiplicative weights. Let $1 \leq p \leq 2$ and $\frac{1}{p} + \frac{1}{p'} = 1$. There exist $R_0 > 0, R_1 \geq R_*$, s.t. for every $n \geq 3$, n odd

$$\sup_{j \in \mathbb{Z}} \mathbf{w}_j \left\| \frac{\mathbf{f}_n(\cdot; j)}{\sum_{l=1}^n \mathbf{v}_{k_l} \prod_{m \neq l} \mathbf{u}_{k_m}} \right\|_{\ell^{p'}(\mathbb{Z}^n)} \leq R_0 R_1^{n-1}, \quad (3.5)$$

and

$$\max_{1 \leq r \leq n} \sup_{j \in \mathbb{Z}} \mathbf{w}_j \left\| \frac{\mathbf{g}_{n,r}(\cdot; j)}{\sum_{l=1}^n \mathbf{v}_{k_l} \prod_{m \neq l} \mathbf{u}_{k_m}} \right\|_{\ell^{p'}(\mathbb{Z}^n)} \leq R_0 R_1^{n-1}. \quad (3.6)$$

We give some examples of weights fulfilling condition (W)_p:

Proposition 3.2. *Let $1 \leq p \leq 2$. Then the following holds true:*

(i) *For any $s \geq 0$, $a \geq 0$ and $0 < b \leq 1$, the weights $\mathbf{u} = \mathbf{v} = \mathbf{w} = \{\langle j \rangle^s e^{a|j|^b}\}_{j \in \mathbb{Z}}$ fulfill (W)_p with constants*

$$R_0 = 1, \quad R_1 = \left(\sum_{k \in \mathbb{Z}} \frac{1}{\langle k \rangle^{p'}} \right)^{1/p'}.$$

(ii) *For any $s \geq 1$, the weights $\mathbf{u} = \{1\}_{j \in \mathbb{Z}}$, $\mathbf{v} = \mathbf{w} = \{\langle j \rangle^s\}_{j \in \mathbb{Z}}$ fulfill (W)_p with constants*

$$R_0 = 1, \quad R_1 = 2^s \left(\sum_{k \in \mathbb{Z}} \frac{1}{\langle k \rangle^{p'}} \right)^{1/p'}.$$

(iii) *For any $s \geq 1$, the weights $\mathbf{u} = \{\langle j \rangle\}_{j \in \mathbb{Z}}$, $\mathbf{v} = \{\langle j \rangle^s\}_{j \in \mathbb{Z}}$ and $\mathbf{w} = \{\langle j \rangle^{s+1}\}_{j \in \mathbb{Z}}$ fulfill (W)_p with constants*

$$R_0 = 1, \quad R_1 = 2^{s+2} \left(\sum_{k \in \mathbb{Z}} \frac{1}{\langle k \rangle^{p'}} \right)^{1/p'}.$$

The proof of the Proposition is postponed in Appendix C.

The main result of this section is the following theorem:

Theorem 3.3. *Fix $1 \leq p \leq 2$. There exist $\mathbf{C}, \varrho_* > 0$ and an analytic map $\Psi : B^p(\varrho_*) \rightarrow \ell_c^p$ s.t. the following is true:*

(i) *The quantities $|\Psi_j(\xi, \bar{\xi})|^2$ are in involution $\forall (\xi, \bar{\xi}) \in B_r^p(\varrho_*)$; more precisely $|\Psi_j(\xi, \bar{\xi})|^2 = \gamma_j^2(\varphi, \bar{\varphi})$, where γ_j is the j^{th} spectral gap (see (3.11)) and $\varphi = \mathcal{F}^{-1}(\xi)$.*

(ii) *For any $\mathbf{u} \leq \mathbf{v} \leq \mathbf{w}$ weights fulfilling (W)_p and*

$$0 < \rho \leq \min \left(\varrho_*, \frac{\varrho_*}{R_1} \right) \quad (3.7)$$

the restriction of Ψ to $B^{p,\mathbf{u}}(\rho)$ is analytic as a map $B^{p,\mathbf{u}}(\rho) \rightarrow \ell_c^{p,\mathbf{u}}$; its nonlinear part $\Psi - \mathbb{1}$ is $(p, \mathbf{u}, \mathbf{v}, \mathbf{w})$ -tame majorant analytic, $\Psi - \mathbb{1} \in \mathcal{T}_{\mathbf{u}, \mathbf{v}, \rho}^{\mathbf{w}, 3}$ with the quantitative estimate

$$\|\Psi - \mathbb{1}\|_{\mathcal{T}_{\mathbf{u}, \mathbf{v}, \rho}^{\mathbf{w}, 3}} \leq \mathbf{C} R_0 R_1^2 \rho^3. \quad (3.8)$$

(iii) Ψ is real for real sequences, i.e. $\Psi: B_r^p(\varrho_*) \rightarrow \ell_r^p$.

Remark 3.4. The constants \mathbf{C}, ϱ_* in Theorem 3.3 do not depend on the regularity of ζ , and it is possible to compute them, see Section 3.1.

Before proving Theorem 3.3, we show how Theorem 3.3 and the Kuksin-Perelman theorem 2.19 imply Theorem 1.1.

Proof of Theorem 1.1. Let $1 \leq p \leq 2$. First take $\mathbf{u} = \mathbf{v} = \mathbf{w} = \{1\}_{j \in \mathbb{Z}}$. By Proposition 3.2 such weights fulfill $(W)_p$ with $R_0 = 1$ and $R_1 = \left(\sum_{k \in \mathbb{Z}} \frac{1}{\langle k \rangle^{p'}} \right)^{1/p'} \leq 2$. For $\rho \leq 2^{-1} \varrho_*$, $\Psi - \mathbb{1} \in \mathcal{T}_{\mathbf{u}, \mathbf{v}, \rho}^{\mathbf{w}, 3}$ with

$$\epsilon_1 := \|\Psi - \mathbb{1}\|_{\mathcal{T}_{\mathbf{u}, \mathbf{v}, \rho}^{\mathbf{w}, 3}} \leq 4\mathbf{C}\rho^3.$$

Hence condition (2.21) is satisfied if $\rho \leq \min(2^{-31}\mathbf{C}^{-1/2}, 2^{-1}\varrho_*) =: \varrho_1$. Applying the tame Kuksin-Perelman theorem, there exists a map $\tilde{\Psi}: B_r^p(\mathbf{a}\varrho_1) \rightarrow \ell_r^p$, $\mathbf{a} = 2^{-120}$, which fulfills $i)-v)$ of Theorem 2.19. In particular $\tilde{\Psi} - \mathbb{1} \in \mathcal{A}_{\mathbf{u}, \mathbf{a}\rho}^3$ for any $\rho \leq \varrho_1$ with

$$\|\tilde{\Psi} - \mathbb{1}\|_{\mathcal{A}_{\mathbf{u}, \mathbf{a}\rho}^3} \leq 2^{17}\epsilon_1 \leq 2^{19}\mathbf{C}\rho^3.$$

Denote now $\Phi := \tilde{\Psi}|_{B_r^p(\mathbf{a}\rho)}$. Such map is majorant analytic as a map $B^p(\rho_0) \rightarrow \ell_c^p$ for $\rho_0 \equiv \mathbf{a}\varrho_1$, and fulfills (i) and (ii) of Theorem 1.1.

We prove now (iii) . Take $\mathbf{u} = \{1\}_{j \in \mathbb{Z}}$, $\mathbf{v} = \mathbf{w} = \{\langle j \rangle^s\}_{j \in \mathbb{Z}}$ with $s \geq 1$. By Proposition 3.2(ii) these weights fulfill $(W)_p$ with constants $R_0 = 1$ and $R_1 = 2^s \left(\sum_{k \in \mathbb{Z}} \frac{1}{\langle k \rangle^{p'}} \right)^{1/p'} \leq 2^{s+1}$. Therefore for $0 < \rho < \varrho_* 2^{-1-s}$, one has that $\Psi - \mathbb{1} \in \mathcal{T}_{\mathbf{u}, \mathbf{v}, \rho}^{\mathbf{w}, 3}$ and $\epsilon_1 \equiv \|\Psi - \mathbb{1}\|_{\mathcal{T}_{\mathbf{u}, \mathbf{v}, \rho}^{\mathbf{w}, 3}} \leq \mathbf{C}2^{2s+2}\rho^3$. Thus, condition (2.21) is fulfilled provided $\rho < \min(2^{-31-s}\mathbf{C}^{-1/2}, 2^{-1}\varrho_*)$, so one applies the Kuksin Perelman theorem obtaining that $\tilde{\Psi} - \mathbb{1} \in \mathcal{T}_{\mathbf{u}, \mathbf{v}, \mathbf{a}\rho}^{\mathbf{w}, 3}$. Then with $\rho' = \mathbf{a}\rho$,

$$\|\Phi - \mathbb{1}\|_{\mathcal{T}_{\mathbf{u}, \mathbf{v}, \rho'}^{\mathbf{w}, 3}} \equiv \|\tilde{\Psi} - \mathbb{1}\|_{\mathcal{T}_{\mathbf{u}, \mathbf{v}, \mathbf{a}\rho}^{\mathbf{w}, 3}} \leq \mathbf{C}'4^s\rho'^3.$$

This estimate and Remark 2.16 implies (iii) . Item (v) is proved analogously using the weights $\mathbf{u} = \{\langle j \rangle\}_{j \in \mathbb{Z}}$, $\mathbf{v} = \{\langle j \rangle^s\}_{j \in \mathbb{Z}}$ and $\mathbf{w} = \{\langle j \rangle^{s+1}\}_{j \in \mathbb{Z}}$.

Item (iii) is a consequence of Corollary 2.20. The explicit form of the mass and the momentum in Birkhoff coordinates follows by the result of [GK14] and the remark that, despite the integrating Birkhoff map is not necessarily unique, it is unique the normal form. \square

Proof of Theorem 1.2 and Corollary 2.20. Theorem 1.2 follows with the same arguments employed in the proof of Theorem 1.1, using that the weights $\mathbf{u} = \mathbf{v} = \mathbf{w} = \{\langle j \rangle^s e^{a|j|}\}_{j \in \mathbb{Z}}$ fulfill $(W)_p$ with constants R_0 and R_1 which do not depend on s, a .

The proof of Corollary 2.20 follows as in [BM16, Corollary 1.6]. \square

3.1 Proof of Theorem 3.3

As we already mentioned, in order to construct the map Ψ we will exploit the integrable structure of the dNLS, which we now recall. It is well known that dNLS on \mathbb{T} admits a Lax pair formulation, where the Lax operator L is the Zakharov-Shabat differential operator given for any $(\varphi_1, \varphi_2) \in L_c^2$ by

$$L(\varphi_1, \varphi_2) = i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \partial_x + \begin{pmatrix} 0 & \varphi_1 \\ \varphi_2 & 0 \end{pmatrix}. \quad (3.9)$$

We consider (3.9) as an operator on the space

$$\mathcal{Y} := L^2(\mathbb{R}/2\mathbb{Z}, \mathbb{C}) \times L^2(\mathbb{R}/2\mathbb{Z}, \mathbb{C})$$

of functions with periodic boundary conditions on the interval $[0, 2]$ (twice the periodicity of φ_1, φ_2). Often we will denote by $\begin{pmatrix} f_1 \\ f_2 \end{pmatrix}$ the elements of \mathcal{Y} . The space \mathcal{Y} is equipped with the complex scalar product

$$\left(\begin{pmatrix} u_1 \\ u_2 \end{pmatrix}, \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} \right)_{\mathcal{Y}} := \int_0^2 (u_1(x) \bar{v}_1(x) + u_2(x) \bar{v}_2(x)) dx. \quad (3.10)$$

The standard theory of Lax pairs guarantees that the eigenvalues of (3.9) are infinitely many commuting constants of motion. More precisely it is well known [GK14] that there exists $U \subset L_c^2$ a complex neighborhood of L_r^2 in L_c^2 s.t. $\forall (\varphi_1, \varphi_2) \in U$ the spectrum of (3.9) is given by a sequence of complex numbers (lexicographically ordered)

$$\cdots \preceq \lambda_0^-(\varphi_1, \varphi_2) \preceq \lambda_0^+(\varphi_1, \varphi_2) \preceq \lambda_1^-(\varphi_1, \varphi_2) \preceq \lambda_1^+(\varphi_1, \varphi_2) < \cdots$$

The $\{\lambda_j^\pm(\varphi_1, \varphi_2)\}_{j \in \mathbb{Z}}$ are not analytic as functions of (φ_1, φ_2) , thus one prefers to use, rather than the eigenvalues, the *spectral gaps*

$$\gamma_j(\varphi_1, \varphi_2) := \lambda_j^+(\varphi_1, \varphi_2) - \lambda_j^-(\varphi_1, \varphi_2), \quad j \in \mathbb{Z} \quad (3.11)$$

which are known to be real analytic commuting constants of motion. As we already mentioned, the map Ψ that we will construct has the property that, for *real* ζ , $|\Psi_j(\xi, \bar{\xi})|^2 = \gamma_j^2(\varphi, \bar{\varphi}) \quad \forall j \in \mathbb{Z}$, where $\xi = \mathcal{F}(\varphi)$.

Before starting the construction of Ψ , it is useful to state some properties of the Lax operator which will be used in the following. First decompose L as the sum $L = L_0 + V$, where

$$L_0 := i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \partial_x, \quad V(\varphi_1, \varphi_2) := \begin{pmatrix} 0 & \varphi_1 \\ \varphi_2 & 0 \end{pmatrix}. \quad (3.12)$$

By (1.6), we identify $(\varphi_1, \varphi_2) \in L_c^2$ with $\zeta = (\xi, \eta) \in \ell_c^2$, thus from now on we will write $V(\zeta) \equiv V(\varphi_1, \varphi_2)$. The following properties are trivially verified:

(H1) **Involution** ι : denote by ι the bounded, antilinear operator $\mathcal{Y} \rightarrow \mathcal{Y}$ defined by

$$\iota \begin{pmatrix} u_1 \\ u_2 \end{pmatrix} = \begin{pmatrix} \bar{u}_2 \\ \bar{u}_1 \end{pmatrix}.$$

Then $\forall f, g \in \mathcal{Y}, \forall \lambda \in \mathbb{C}$ one has

$$\iota^2 f = f, \quad \overline{(f, g)}_{\mathcal{Y}} = (\iota f, \iota g)_{\mathcal{Y}}, \quad \iota(\lambda f) = \bar{\lambda} \iota f.$$

Furthermore

$$\iota L_0 = L_0 \iota, \quad \iota V(\zeta) = V(\zeta)^* \iota. \quad (3.13)$$

(H2) **Spectrum of L_0** : L_0 is a selfadjoint operator with domain $D(L_0)$ dense in \mathcal{Y} . Its spectrum is discrete, $\sigma(L_0) = \{\lambda_j^0\}_{j \in \mathbb{Z}}$, and each eigenvalue $\lambda_j^0 \equiv \pi j$ has multiplicity 2. Remark that

$$\inf_{i \neq j} |\lambda_j^0 - \lambda_i^0| = \pi. \quad (3.14)$$

(H3) **Eigenfunctions of L_0** : for any $j \in \mathbb{Z}$ we denote by $f_{j0}^+, f_{j0}^- \in \mathcal{Y}$ the eigenfunctions corresponding to the eigenvalue λ_j^0 given by

$$f_{j0}^- := \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ e^{i\pi j x} \end{pmatrix}, \quad f_{j0}^+ := \frac{1}{\sqrt{2}} \begin{pmatrix} e^{-i\pi j x} \\ 0 \end{pmatrix}.$$

They fulfill

$$\iota f_{j0}^- = f_{j0}^+, \quad \iota f_{j0}^+ = f_{j0}^-. \quad (3.15)$$

We denote by $E_{j0} := \text{Vect}(f_{j0}^-, f_{j0}^+)$ the vectorial space spanned by f_{j0}^\pm . The vectors $\{f_{j0}^\sigma\}_{j \in \mathbb{Z}, \sigma \in \pm}$ form a basis for \mathcal{Y} .

(H4) **Perturbation $V(\zeta)$:** for any $\zeta \in \ell_c^2$, the operator $V(\zeta)$ has domain $D(V(\zeta)) \supset D(L_0)$.

(H4a) for any *real* $\zeta \in \ell_r^2$, the operator $V(\zeta)$ is symmetric (w.r.t. the scalar product (3.10)) on its domain.

(H4b) for any $i_1, i_2 \in \mathbb{Z}$ we have that

$$\begin{aligned} (V(\zeta) f_{i_1 0}^-, f_{i_2 0}^+)_{\mathcal{Y}} &= \xi_j, & \text{if } i_1 + i_2 = 2j \\ (V(\zeta) f_{i_1 0}^+, f_{i_2 0}^-)_{\mathcal{Y}} &= \eta_j, & \text{if } i_1 + i_2 = 2j \\ (V(\zeta) f_{i_1 0}^{\sigma_1}, f_{i_2 0}^{\sigma_2})_{\mathcal{Y}} &= 0, & \text{otherwise} \end{aligned}$$

Now take $(\xi, \eta) \in \ell_c^2$ with a sufficiently small norm. We will construct perturbatively the spectral data of $L(\zeta)$ (defined in (3.9)) starting from the spectral data of L_0 (defined in (3.12)) which is given in (H2) and (H3). We start with a preliminary result:

Lemma 3.5. *For any $\lambda \in \mathbb{C} \setminus \sigma(L_0)$ the map $\zeta \mapsto V(\zeta)(L_0 - \lambda)^{-1}$ is analytic as a map $\ell_c^2 \rightarrow \mathcal{L}(\mathcal{Y})$ and fulfills*

$$\|V(\zeta)(L_0 - \lambda)^{-1}\|_{\mathcal{L}(\mathcal{Y})} \leq 2c(\lambda) \|\zeta\|_2, \quad (3.16)$$

where

$$c(\lambda) := \left(\sum_{n \in \mathbb{Z}} \frac{1}{|\pi n - \lambda|^2} \right)^{1/2}.$$

Proof. The map $\zeta \mapsto V(\zeta)(L_0 - \lambda)^{-1}$ is a \mathbb{C} -linear map, so it is analytic iff it is bounded. To compute its norm, take $f \in \mathcal{Y}$ and write

$$f = (f_1, f_2) = \frac{1}{\sqrt{2}} \left(\sum_{k \in \mathbb{Z}} \alpha_k e^{-i\pi k x}, \sum_k \beta_k e^{i\pi k x} \right),$$

so that

$$(L_0 - \lambda)^{-1} f \equiv (\tilde{f}_1, \tilde{f}_2) = \frac{1}{\sqrt{2}} \left(\sum_k \frac{\alpha_k}{\pi k - \lambda} e^{-i\pi j x}, \sum_k \frac{\beta_k}{\pi k - \lambda} e^{i\pi k x} \right).$$

Now remark that $\tilde{f}_1, \tilde{f}_2 \in L^\infty[0, 2]$ and

$$\|\tilde{f}_1\|_{L^\infty[0, 2]} \leq \frac{1}{\sqrt{2}} \left(\sum_k |\alpha_k|^2 \right)^{1/2} \left(\sum_k \frac{1}{|\pi k - \lambda|^2} \right)^{1/2} \equiv c(\lambda) \|f_1\|_{L^2[0, 2]}$$

(clearly the same bound holds also for \tilde{f}_2). Thus

$$\begin{aligned} \|V(\zeta)(L_0 - \lambda)^{-1} f\|_{\mathcal{Y}} &= \left(\|\varphi_2 \tilde{f}_1\|_{L^2[0, 2]}^2 + \|\varphi_1 \tilde{f}_2\|_{L^2[0, 2]}^2 \right)^{1/2} \\ &\leq \sqrt{2} \|(\varphi_1, \varphi_2)\|_{L_c^2} \left(\|\tilde{f}_1\|_{L^\infty[0, 2]}^2 + \|\tilde{f}_2\|_{L^\infty[0, 2]}^2 \right)^{1/2} \leq 2c(\lambda) \|\zeta\|_2 \|f\|_{\mathcal{Y}} \end{aligned}$$

which is the claimed inequality. \square

Now it is sufficient to apply classical Kato perturbation theory [Kat66] to get the following:

Lemma 3.6. *Let $0 < \rho < \frac{1}{8}$. Then for any $\|\zeta\|_2 \leq \rho$ the following holds true:*

(i) For any $j \in \mathbb{Z}$, let

$$\Gamma_j := \{\lambda \in \mathbb{C} : |\lambda - \lambda_j^0| = \pi/2\}. \quad (3.17)$$

Then $\Gamma_j \subset \rho(L(\zeta))$, the resolvent set of $L(\zeta)$.

(ii) For any $j \in \mathbb{Z}$, define the projector $P_j(\zeta)$ and the subspace $E_j(\zeta) \subseteq \mathcal{Y}$ as

$$P_j(\zeta) := -\frac{1}{2\pi i} \oint_{\Gamma_j} (L(\zeta) - \lambda)^{-1} d\lambda, \quad E_j(\zeta) := \text{Ran } P_j(\zeta), \quad j \in \mathbb{Z} \quad (3.18)$$

where Γ_j is counterclockwise oriented. Then $\zeta \mapsto P_j(\zeta)$ is analytic as a map $B^2(\rho) \rightarrow \mathcal{L}(\mathcal{Y})$.

(iii) For any $j \in \mathbb{Z}$, define the transformation operator $U_j(\zeta)$ as

$$U_j(\zeta) := (\mathbb{1} - (P_j(\zeta) - P_{j0})^2)^{-1/2} P_j(\zeta) ,$$

where $P_{j0} := P_j(0)$. Then $\zeta \mapsto U_j(\zeta)$ is analytic as a map $B^2(\rho) \rightarrow \mathcal{Y}$ and $\text{Ran } U_j(\zeta) \equiv E_j(\zeta)$. Furthermore for ζ real

$$\|U_j(\zeta)f\|_{\mathcal{Y}} = \|f\|_{\mathcal{Y}} , \quad \forall f \in E_j^0 , \quad (3.19)$$

$$[\iota, U_j(\zeta)] = 0 . \quad (3.20)$$

Proof. By Lemma 3.5, for any $\lambda \in \Gamma_j$, the map $\zeta \mapsto V(\zeta)(L_0 - \lambda)^{-1}$ is analytic as a map $\ell_c^2 \rightarrow \mathcal{L}(\mathcal{Y})$ and

$$\sup_{\lambda \in \Gamma_j} \|V(\zeta)(L_0 - \lambda)^{-1}\|_{\mathcal{L}(\mathcal{Y})} \leq \sup_{\lambda \in \Gamma_j} 2c(\lambda) \|\zeta\|_2 \leq 4\|\zeta\|_2 \quad (3.21)$$

where in the last step we used the explicit formula for Γ_j to estimate $c(\lambda)$. It follows that for $\|\zeta\|_2 \leq \rho < \frac{1}{8}$, the perturbed resolvent $L(\zeta) - \lambda$ is well defined by Neumann series and fulfills the estimate

$$\sup_{\lambda \in \Gamma_j} \|(L(\zeta) - \lambda)^{-1}\|_{\mathcal{L}(\mathcal{Y})} \leq \frac{4}{\pi} . \quad (3.22)$$

Thus for $\|\zeta\|_2 \leq \rho < 1/8$ the projector $P_j(\zeta)$ in (3.18) is well defined and analytic $\forall j \in \mathbb{Z}$. By the resolvent identity

$$P_j(\zeta) - P_j(0) = \frac{1}{2\pi i} \oint_{\Gamma_j} (L(\zeta) - \lambda)^{-1} V(\zeta) (L_0 - \lambda)^{-1} d\lambda , \quad \forall j \in \mathbb{Z}$$

which together with (3.21) and (3.22) gives the estimate

$$\|P_j(\zeta) - P_j(0)\|_{\mathcal{L}(\mathcal{Y})} \leq 8\|\zeta\|_2 < 1$$

provided $\|\zeta\|_2 \leq \rho < 1/8$. Hence also $U_j(\zeta)$ is defined by Neumann series.

We prove now (3.20). First we claim that $\iota P_j(\zeta) = P_j(\zeta)^* \iota$ for every ζ sufficiently small. This follows by a direct computation using that by (H2), (H4a) $\iota(L - \lambda)^{-1} = [(L - \lambda)^{-1}]^* \iota$. Since for ζ real $L(\zeta)$ is self-adjoint, one has $P_j(\zeta)^* = P_j(\zeta)$, hence $[\iota, P_j(\zeta)] = 0$. (3.20) follows easily. \square

For any $j \in \mathbb{Z}$ let us set now

$$f_j^\pm(\zeta) := U_j(\zeta) f_{j0}^\pm \in E_j(\zeta) . \quad (3.23)$$

Remark that the $f_j^\pm(\zeta)$'s do not need to be eigenvectors, but they span the eigenspace $E_j(\zeta)$ and are analytic as functions of ζ .

Finally for any $j \in \mathbb{Z}$ let us define

$$z_j(\zeta) := ((L(\zeta) - \lambda_j^0) f_j^-(\zeta), \iota f_j^-(\zeta))_{\mathcal{Y}} , \quad w_j(\zeta) := ((L(\zeta) - \lambda_j^0) f_j^+(\zeta), \iota f_j^+(\zeta))_{\mathcal{Y}} . \quad (3.24)$$

Such coordinates fulfill the following properties:

Lemma 3.7. *For any $\rho < \frac{1}{8}$ the following holds true:*

(i) $\forall j \in \mathbb{Z}$ the map $B^2(\rho) \rightarrow \mathbb{C}^2$, $\zeta \mapsto (z_j(\zeta), w_j(\zeta))$ is analytic.

(ii) $\forall j \in \mathbb{Z}$, for any real $\zeta \in B_r^2(\rho)$ one has $\overline{z_j(\zeta)} = w_j(\zeta)$.

(iii) $\forall j \in \mathbb{Z}$, for any real $\zeta \in B_r^2(\rho)$ one has $|z_j(\zeta)|^2 = (\lambda_j^+(\zeta) - \lambda_j^-(\zeta))^2$.

Proof. (i) Define $A : \mathcal{Y} \times \mathcal{Y} \rightarrow \mathbb{C}$, $(f, g) \mapsto A(f, g) := (f, \iota g)_{\mathcal{Y}}$. By (H1) A is a \mathbb{C} -multilinear continuous map, hence it is analytic in each components (see e.g. [Muj86]).

Then $z_j(\zeta) = A((L(\zeta) - \lambda_j^0) U_j(\zeta) f_{j0}^-, U_j(\zeta) f_{j0}^-)_{\mathcal{Y}}$ is composition of analytic maps and hence it is

analytic. Analogous for $\zeta \mapsto w_j(\zeta)$.

(ii) We claim that for ζ real

$$\imath f_j^-(\zeta) = f_j^+(\zeta), \quad \imath f_j^+(\zeta) = f_j^-(\zeta), \quad \forall j \in \mathbb{Z}. \quad (3.25)$$

This follows from the following chain of equalities (which hold for ζ real and sufficiently small)

$$f_j^+(\zeta) \stackrel{(3.23)}{=} U_j(\zeta) f_{j0}^+ \stackrel{(3.15)}{=} U_j(\zeta) \imath f_{j0}^- \stackrel{(3.20)}{=} \imath U_j(\zeta) f_{j0}^- \stackrel{(3.23)}{=} \imath f_j^-(\zeta).$$

Thus for ζ real and $\|\zeta\|_2 \leq \rho$ one has

$$\begin{aligned} \overline{z_j(\zeta)} &= \overline{((L(\zeta) - \lambda_j^0) f_j^-(\zeta), \imath f_j^-(\zeta))_{\mathcal{Y}}} = (\imath f_j^-(\zeta), (L(\zeta) - \lambda_j^0) f_j^-(\zeta))_{\mathcal{Y}} \\ &= ((L(\zeta) - \lambda_j^0) \imath f_j^-(\zeta), f_j^-(\zeta))_{\mathcal{Y}} \stackrel{(3.25)}{=} ((L(\zeta) - \lambda_j^0) f_j^+(\zeta), \imath f_j^+(\zeta))_{\mathcal{Y}} = w_j(\zeta), \end{aligned}$$

where in the third equality we used that for ζ real $L(\zeta) - \lambda_j^0$ is self-adjoint.

(iii) By Lemma 3.6, for ζ real and $\|\zeta\|_2 \leq \rho$ the operator $U_j(\zeta)|_{E_j^0}$ is unitary. Since f_{j0}^+, f_{j0}^- are orthogonal in \mathcal{Y} , the vectors $f_j^+(\zeta), f_j^-(\zeta)$ are orthogonal as well, thus form a basis for $E_j(\zeta)$. Let $M_j(\zeta)$ be the matrix of the self-adjoint operator $L(\zeta) - \lambda_j^0$ on this basis. One has

$$M_j(\zeta) = \begin{bmatrix} ((L(\zeta) - \lambda_{j0}) f_j^-(\zeta), f_j^-(\zeta))_{\mathcal{Y}} & ((L(\zeta) - \lambda_{j0}) f_j^-(\zeta), f_j^+(\zeta))_{\mathcal{Y}} \\ ((L(\zeta) - \lambda_{j0}) f_j^+(\zeta), f_j^-(\zeta))_{\mathcal{Y}} & ((L(\zeta) - \lambda_{j0}) f_j^+(\zeta), f_j^+(\zeta))_{\mathcal{Y}} \end{bmatrix} = \begin{bmatrix} a_j^1 & b_j \\ \bar{b}_j & a_j^2 \end{bmatrix}$$

for some $a_j^1, a_j^2, b_j \in \mathbb{C}$. We show now that $a_j^1 = a_j^2 \in \mathbb{R}$. Indeed using the self-adjointness of $L(\zeta) - \lambda_j^0$,

$$\overline{((L(\zeta) - \lambda_j^0) f_j^-(\zeta), f_j^-(\zeta))_{\mathcal{Y}}} = ((L(\zeta) - \lambda_j^0) f_j^-(\zeta), f_j^-(\zeta))_{\mathcal{Y}}$$

while using (3.13), (3.25)

$$\overline{((L(\zeta) - \lambda_j^0) f_j^-(\zeta), f_j^-(\zeta))_{\mathcal{Y}}} = ((L(\zeta) - \lambda_j^0) \imath f_j^-(\zeta), \imath f_j^-(\zeta))_{\mathcal{Y}} = ((L(\zeta) - \lambda_j^0) f_j^+(\zeta), f_j^+(\zeta))_{\mathcal{Y}}.$$

The eigenvalues of $M_j(\zeta)$ are the eigenvalues of $L(\zeta) - \lambda_j^0|_{E_j(\zeta)}$, i.e. $\lambda_j^{\pm}(\zeta) - \lambda_j^0$. Then

$$(\lambda_j^+(\zeta) - \lambda_j^0)^2 = (\text{Tr } M_j(\zeta))^2 - 4 \det M_j(\zeta) = |b_j|^2.$$

Now remark that by (3.25) one has $z_j(\zeta) \equiv b_j, \forall j \in \mathbb{Z}$. \square

Since the maps $\zeta \mapsto z_j(\zeta), \zeta \mapsto w_j(\zeta)$ are analytic, we can expand them in their absolutely and uniformly convergent Taylor series

$$z_j(\zeta) = \sum_{n=1}^{\infty} Z_j^n(\zeta), \quad w_j(\zeta) = \sum_{n=1}^{\infty} W_j^n(\zeta), \quad j \in \mathbb{Z}, \quad (3.26)$$

where Z_j^n and W_j^n are homogeneous polynomials of degree n in ζ . By a direct computation the first terms in the Taylor series are given by

$$\begin{aligned} z_j(\zeta) &= (V(\zeta) f_{j0}^-, f_{j0}^+) + \left(V(\zeta) (L_0 - \lambda_j^0)^{-1} (\mathbb{1} - P_{j0}) V(\zeta) f_{j0}^-, f_{j0}^+ \right) + h.o.t. \\ w_j(\zeta) &= (V(\zeta) f_{j0}^+, f_{j0}^-) + \left(V(\zeta) (L_0 - \lambda_j^0)^{-1} (\mathbb{1} - P_{j0}) V(\zeta) f_{j0}^+, f_{j0}^- \right) + h.o.t. \end{aligned} \quad (3.27)$$

Using (H4b) one verifies that

$$Z_j^1(\zeta) = \xi_j, \quad W_j^1(\zeta) = \eta_j, \quad Z_j^2(\zeta) = W_j^2(\zeta) = 0. \quad (3.28)$$

The expression for the general homogeneous term $Z_j^n(\zeta)$ is much more involved and is given in the following proposition:

Proposition 3.8. For $\|\zeta\|_2 \leq \rho < 1/8$, the homogeneous polynomials $Z_j^n(\zeta)$ and $W_j^n(\zeta)$ are given $\forall j \in \mathbb{Z}, \forall n \in \mathbb{N}, n \geq 3$ by

$$\begin{aligned} Z_j^n(\zeta) &= \sum_{k_1+\dots+k_n=j} \mathcal{K}_j^n(k_1, \dots, k_n) \xi_{k_1} \eta_{-k_2} \xi_{k_3} \cdots \eta_{-k_{n-1}} \xi_{k_n} \\ W_j^n(\zeta) &= \sum_{k_1+\dots+k_n=j} \mathcal{K}_j^n(k_1, \dots, k_n) \eta_{k_1} \xi_{-k_2} \eta_{k_3} \cdots \xi_{-k_{n-1}} \eta_{k_n} \end{aligned} \quad , \quad \text{if } n \text{ odd} \quad (3.29)$$

and by

$$Z_j^n(\zeta) = W_j^n(\zeta) = 0 \quad , \quad \text{if } n \text{ even} .$$

The kernel $\mathcal{K}_j^n(k_1, \dots, k_n)$ has support in $k_1 + \dots + k_n = j$ and there exist $K_0, K_1 > 0$ s.t.

$$|\mathcal{K}_j^n(k_1, \dots, k_n)| \leq K_0 \cdot K_1^{n-1} \prod_{\substack{m=1 \\ m \text{ odd}}}^{n-1} \frac{1}{\langle k_1 + \dots + k_m - j \rangle} \cdot \frac{1}{\langle k_1 + \dots + k_{m+1} \rangle} . \quad (3.30)$$

The proof of the lemma, being quite technical, is postponed in Appendix B.

Remark 3.9. With \mathfrak{f}_n defined in (3.1), one has

$$|\mathcal{K}_j^n(k_1, \dots, k_n)| \leq K_0 \cdot K_1^{n-1} \mathfrak{f}_n(k_1, \dots, k_n; j) \quad (3.31)$$

We are finally ready to define the map Ψ of Theorem 3.3. First for $\zeta \in B^2(\rho)$, $\rho < 1/8$, let

$$Z(\zeta) := (z_j(\zeta))_{j \in \mathbb{Z}} \quad , \quad W(\zeta) := (w_j(\zeta))_{j \in \mathbb{Z}} .$$

Now for any $\zeta \in B^2(\rho)$ we define the map $\Psi : B^2(\rho) \rightarrow \mathbb{C}^{\mathbb{Z}} \times \mathbb{C}^{\mathbb{Z}}$ by

$$\Psi(\zeta) := \left(Z(\zeta) , W(\zeta) \right) . \quad (3.32)$$

In the rest of the section we show that Ψ fulfills the properties claimed in Theorem 3.3.

First note that, at least formally, by (3.28) $\Psi = \mathbb{1} + \Psi_3$, where $\Psi_3(\zeta) := (Z_3(\zeta), W_3(\zeta))$ and $Z_3 := Z - \mathbb{1}$, $W_3 := W - \mathbb{1}$ are $O(\zeta^3)$.

In the next proposition we show that, provided $\mathfrak{u} \leq \mathfrak{v} \leq \mathfrak{w}$ are weights fulfilling $(W)_p$, Ψ_3 is analytic and tame majorant analytic as a map $B^{p,\mathfrak{u}}(\rho) \cap \ell_c^{p,\mathfrak{v}} \rightarrow \ell_c^{p,\mathfrak{w}}$ (in the sense of Definition 2.5).

Lemma 3.10. Fix $1 \leq p \leq 2$ and let $\mathfrak{u} \leq \mathfrak{v} \leq \mathfrak{w}$ be weights fulfilling $(W)_p$ with constants R_0, R_1 . Then for any

$$0 < \rho < \min \left(\frac{1}{8}, \frac{1}{8 K_1 R_1} \right)$$

the map $\zeta \mapsto \Psi_3(\zeta)$ is analytic and tame majorant analytic, $\Psi_3 \in \mathcal{N}_\rho^T(B^{p,\mathfrak{u}} \cap \ell_c^{p,\mathfrak{v}}; \ell_c^{p,\mathfrak{w}})$ and

$$\langle |\underline{\Psi}_3| \rangle_\rho \leq 2^5 K_0 K_1^2 R_0 R_1^2 \rho^3 . \quad (3.33)$$

Proof. By formally Taylor expanding the map Ψ_3 , one has that $\Psi_3 = \sum_{n \geq 3} \Psi^n$, where

$$\Psi^n(\zeta) := \left(Z^n(\zeta) , W^n(\zeta) \right) , \quad Z^n(\zeta) := (Z_j^n(\zeta))_{j \in \mathbb{Z}} , \quad W^n(\zeta) := (W_j^n(\zeta))_{j \in \mathbb{Z}} .$$

It is sufficient to show that $\underline{\Psi}^n(\zeta) := \left(\underline{Z}^n(\zeta) , \underline{W}^n(\zeta) \right)$ fulfills

$$\|\underline{Z}^n(\zeta)\|_{p,\mathfrak{w}} + \|\underline{W}^n(\zeta)\|_{p,\mathfrak{w}} \leq 2 K_0 R_0 (2 K_1 R_1)^{n-1} \|\zeta\|_{p,\mathfrak{u}}^{n-1} \|\zeta\|_{p,\mathfrak{v}} , \quad \forall n \geq 3 . \quad (3.34)$$

$$\|\underline{Z}^n(\zeta)\|_{p,\mathfrak{u}} + \|\underline{W}^n(\zeta)\|_{p,\mathfrak{u}} \leq 2 K_0 R_0 (2 K_1 R_1)^{n-1} \|\zeta\|_{p,\mathfrak{u}}^n , \quad \forall n \geq 3 . \quad (3.35)$$

Indeed (3.34) implies that $\underline{\Psi}_3 = \sum_{n \geq 3} \underline{\Psi}^n$ fulfills

$$\|\underline{\Psi}_3(\zeta)\|_{p,\mathfrak{w}} \leq \sum_{n \geq 3} \|\underline{\Psi}^n(\zeta)\|_{p,\mathfrak{w}} \leq 2 K_0 R_0 \|\zeta\|_{p,\mathfrak{v}} \sum_{n \geq 3} (2 K_1 R_1)^{n-1} \|\zeta\|_{p,\mathfrak{u}}^{n-1} ,$$

from which one deduces that

$$|\underline{\Psi}_3|_\rho^T \equiv \sup_{\|\zeta\|_{p,u} \leq \rho} \frac{\|\underline{\Psi}_3(\zeta)\|_{p,w}}{\|\zeta\|_{p,v}} \leq 2^4 K_0 K_1^2 R_0 R_1^2 \rho^2 .$$

Analogously, using (3.35) one proves that $|\underline{\Psi}_3|_\rho \equiv \sup_{\|\zeta\|_{p,u} \leq \rho} \|\underline{\Psi}_3(\zeta)\|_{p,u} \leq 2^4 K_0 K_1^2 R_0 R_1^2 \rho^3$. Estimate (3.33) follows by the definition of $\langle |\underline{\Psi}_3| \rangle_\rho := |\underline{\Psi}_3|_\rho + \rho |\underline{\Psi}_3|_\rho^T$. Note that by Remark 2.7 one has that Ψ_3 is analytic.

To prove (3.34), (3.35) we use the explicit formulas for Z_j^n and W_j^n given in Proposition 3.8. We perform the computations only for Z^n , since for W^n they are identical. We begin by proving (3.34). Multiplying and dividing the kernel of Z_j^n by $\sum_{l=1}^n v_{k_l} \prod_{m \neq l} u_{k_m}$ and using Cauchy-Schwartz one gets for $1/p + 1/p' = 1$

$$|Z_j^n(\zeta)| \leq \left\| \frac{\mathcal{K}_j^n(\mathbf{k})}{\sum_{l=1}^n v_{k_l} \prod_{m \neq l} u_{k_m}} \right\|_{\ell^{p'}(\mathbb{Z}^n)} \cdot \sum_{l=1}^n \beta_{j,l} ,$$

$$\beta_{j,l} := \left(\sum_{k_1 + \dots + k_n = j} v_{k_l}^p |\zeta_{k_l}|^p \prod_{m \neq l} u_{k_m}^p |\zeta_{k_m}|^p \right)^{1/p} ;$$

by Remark 3.9

$$|Z_j^n(\zeta)| \leq K_0 \cdot K_1^{n-1} \left\| \frac{f_n(\cdot; j)}{\sum_{l=1}^n v_{k_l} \prod_{m \neq l} u_{k_m}} \right\|_{\ell^{p'}(\mathbb{Z}^n)} \cdot \sum_{l=1}^n \beta_{j,l} .$$

Now we have

$$\begin{aligned} \|\underline{Z}^n(|\zeta|)\|_{p,w} &\leq K_0 \cdot K_1^{n-1} \sup_{j \in \mathbb{Z}} \left[w_j \left\| \frac{f_n(\cdot; j)}{\sum_{l=1}^n v_{k_l} \prod_{m \neq l} u_{k_m}} \right\|_{\ell^{p'}(\mathbb{Z}^n)} \right] \left\| \sum_{l=1}^n \beta_{j,l} \right\|_{\ell_j^p} \\ &\leq K_0 R_0 \cdot (K_1 R_1)^{n-1} \sum_{l=1}^n \|\beta_{j,l}\|_{\ell_j^p} \leq K_0 R_0 \cdot (K_1 R_1)^{n-1} \cdot n \|\zeta\|_{p,v} \|\zeta\|_{p,u}^{n-1} \\ &\leq K_0 R_0 \cdot (2 K_1 R_1)^{n-1} \cdot \|\zeta\|_{p,v} \|\zeta\|_{p,u}^{n-1} \end{aligned}$$

where to pass to the second line we used (3.5), while in the third inequality we used that

$$\|\beta_{j,l}\|_{\ell_j^p} \equiv \left(\sum_j |\beta_{j,l}|^p \right)^{1/p} \leq \left(\sum_j \sum_{k_1 + \dots + k_n = j} v_{k_l}^p |\zeta_{k_l}|^p \prod_{m \neq l} u_{k_m}^p |\zeta_{k_m}|^p \right)^{1/p} \leq \|\zeta\|_{p,v} \|\zeta\|_{p,u}^{n-1} .$$

To prove (3.35) it is enough to repeat the computations above with $w = v = u$ and use (3.3). \square

Next we study the operator $d\Psi_3(\zeta)^t$, which is the transposed of $d\Psi_3(\zeta)$ w.r.t. the (complex) scalar product:

$$\left(d\Psi_3(\zeta)(\xi^1, \eta^1), \overline{(\xi^2, \eta^2)} \right) = \left((\xi^1, \eta^1), \overline{d\Psi_3(\zeta)^t(\xi^2, \eta^2)} \right) , \quad (\xi^1, \eta^1), (\xi^2, \eta^2) \in \ell_c^2 .$$

We prove the following result:

Lemma 3.11. *With the same assumptions of Lemma 3.10, let*

$$0 < \rho < \min \left(\frac{1}{8}, \frac{1}{4 K_1 R_1} \right) . \quad (3.36)$$

Then the following holds true:

(i) *For any $n \geq 3$, $\zeta \mapsto d\Psi^n(\zeta)^t \in \mathcal{N}_\rho^T(B^{p,u} \cap \ell_c^{p,v}; \mathcal{L}(\ell_c^{p,v}, \ell_c^{p,w}))$ and fulfills*

$$\langle |d\Psi^n|^t | \rangle_\rho \leq 2^5 K_0 R_0 (2 K_1 R_1)^{n-1} \rho^{n-1} , \quad \forall n \geq 3 . \quad (3.37)$$

(ii) The map $d\Psi_3^* \in \mathcal{N}_\rho^T(B^{p,u} \cap \ell_c^{p,v}, \mathcal{L}(\ell_c^{p,v}, \ell_c^{p,w}))$ and moreover

$$\left\langle \underline{|d\Psi_3^*|} \right\rangle_\rho \leq 2^8 K_0 K_1^2 R_0 R_1^2 \rho^2 . \quad (3.38)$$

Proof. We will prove that for any $n \geq 3$

$$\|\underline{d\Psi^n}(\zeta)^t v\|_{p,w} \leq 2^4 K_0 R_0 (2K_1 R_1)^{n-1} (\|v\|_{p,v} \|\zeta\|_{p,u} + \|v\|_{p,u} \|\zeta\|_{p,v}) \|\zeta\|_{p,u}^{n-2} , \quad (3.39)$$

$$\|\underline{d\Psi^n}(\zeta)^t v\|_{p,u} \leq 2^4 K_0 R_0 (2K_1 R_1)^{n-1} \|v\|_{p,u} \|\zeta\|_{p,u}^{n-1} , \quad (3.40)$$

from which item (i) follows.

The j^{th} -component of $d\Psi^n(\zeta)^t v$ is given by $[d\Psi^n(\zeta)^t v]_j = \left(A_j^n(\zeta)v , B_j^n(\zeta)v \right)$ where, letting $v = (\tilde{\xi}, \tilde{\eta})$,

$$A_j^n(\zeta)(\tilde{\xi}, \tilde{\eta}) := \sum_\ell \frac{\partial Z_\ell^n(\zeta)}{\partial \xi_j} \tilde{\xi}_\ell + \frac{\partial W_\ell^n(\zeta)}{\partial \xi_j} \tilde{\eta}_\ell , \quad B_j^n(\zeta)(\tilde{\xi}, \tilde{\eta}) := \sum_\ell \frac{\partial Z_\ell^n(\zeta)}{\partial \eta_j} \tilde{\xi}_\ell + \frac{\partial W_\ell^n(\zeta)}{\partial \eta_j} \tilde{\eta}_\ell .$$

To compute such terms explicitly we use Proposition 3.8. One has for example (we compute only $A_j^n(\zeta)(\tilde{\xi}, \tilde{\eta})$, the other is analogous)

$$\begin{aligned} A_j^n(\zeta)(\tilde{\xi}, \tilde{\eta}) &= I_j + II_j , \\ I_j &:= \sum_{\ell \in \mathbb{Z}} \frac{\partial Z_\ell^n(\zeta)}{\partial \xi_j} \tilde{\xi}_\ell = \sum_{\substack{r=1 \\ r \text{ odd}}}^n \sum_{\mathbf{k} \in \mathfrak{S}_{-j}^{n,r}} \mathcal{K}_{k_r}^n(k_1, \dots, k_{r-1}, j, k_{r+1}, \dots, k_n) \tilde{\xi}_{k_r} \cdot \frac{\xi_{k_1} \eta_{-k_2} \dots \xi_{k_n}}{\xi_{k_r}} \\ II_j &:= \sum_{\ell \in \mathbb{Z}} \frac{\partial W_\ell^n(\zeta)}{\partial \xi_j} \tilde{\eta}_\ell = \sum_{\substack{r=1 \\ r \text{ even}}}^n \sum_{\mathbf{k} \in \mathfrak{S}_j^{n,r}} \mathcal{K}_{k_r}^n(k_1, \dots, k_{r-1}, -j, k_{r+1}, \dots, k_n) \tilde{\eta}_{k_r} \cdot \frac{\eta_{k_1} \xi_{-k_2} \dots \eta_{k_n}}{\xi_{k_r}} \end{aligned}$$

where $\mathbf{k} := (k_1, \dots, k_n) \in \mathbb{Z}^n$ and $\mathfrak{S}_a^{n,r}$ is the set defined by

$$\mathfrak{S}_a^{n,r} := \left\{ \mathbf{k} \in \mathbb{Z}^n : \sum_{\substack{i=1 \\ i \neq r}}^n k_i - k_r = a \right\} . \quad (3.41)$$

Then $|\underline{A_j^n}(\zeta)(\tilde{\xi}, \tilde{\eta})| \leq \underline{I_j} + \underline{II_j}$, where

$$\begin{aligned} \underline{I_j} &:= \sum_{\substack{r=1 \\ r \text{ odd}}}^n \sum_{\mathbf{k} \in \mathfrak{S}_{-j}^{n,r}} |\mathcal{K}_{k_r}^n(k_1, \dots, k_{r-1}, j, k_{r+1}, \dots, k_n)| |\tilde{\xi}_{k_r}| \cdot \frac{|\xi_{k_1}| |\eta_{-k_2}| \dots |\xi_{k_n}|}{|\xi_{k_r}|} , \\ \underline{II_j} &:= \sum_{\substack{r=1 \\ r \text{ even}}}^n \sum_{\mathbf{k} \in \mathfrak{S}_j^{n,r}} |\mathcal{K}_{k_r}^n(k_1, \dots, k_{r-1}, -j, k_{r+1}, \dots, k_n)| |\tilde{\eta}_{k_r}| \cdot \frac{|\eta_{k_1}| |\xi_{-k_2}| \dots |\eta_{k_n}|}{|\xi_{k_r}|} . \end{aligned}$$

To estimate $|\underline{I_j}|$, we first multiply and divide the kernel of $\underline{I_j}$ by $\sum_{l=1}^n \mathbf{v}_{k_l} \prod_{m \neq l} \mathbf{u}_{k_m}$, then we use by Cauchy-Schwartz to obtain

$$\begin{aligned} |\underline{I_j}| &\leq \sum_{\substack{r=1 \\ r \text{ odd}}}^n \left\| \frac{\mathcal{K}_{k_r}^n(k_1, \dots, k_{r-1}, j, k_{r+1}, \dots, k_n)}{\sum_{l=1}^n \mathbf{v}_{k_l} \prod_{m \neq l} \mathbf{u}_{k_m}} \right\|_{\ell^{p'}(\mathbb{Z}^n)} \sum_{l=1}^n \beta_{j,l,r} \\ \beta_{j,l,r} &:= \left(\sum_{\mathbf{k} \in \mathfrak{S}_{-j}^{n,r}} \frac{|\tilde{\xi}_{k_r}|^p |\xi_{k_1}|^p |\eta_{-k_2}|^p \dots |\xi_{k_n}|^p}{|\xi_{k_r}|^p} \mathbf{v}_{k_l}^p \prod_{m \neq l} \mathbf{u}_{k_m}^p \right)^{1/p} ; \end{aligned}$$

then since $|\mathcal{K}_{k_r}^n(k_1, \dots, k_{r-1}, \pm j, k_{r+1}, \dots, k_n)| \leq K_0 \cdot K_1^{n-1} \mathfrak{g}_{n,r}(\mathbf{k}; \pm j)$ for any $\mathbf{k} \in \mathfrak{S}_{\mp j}^{n,r}$ one gets

$$|\underline{I_j}| \leq K_0 \cdot K_1^{n-1} \sum_{\substack{r=1 \\ r \text{ odd}}}^n \left\| \frac{\mathfrak{g}_{n,r}(\cdot; j)}{\sum_{l=1}^n \mathbf{v}_{k_l} \prod_{m \neq l} \mathbf{u}_{k_m}} \right\|_{\ell^{p'}(\mathbb{Z}^n)} \sum_{l=1}^n \beta_{j,l,r} .$$

Finally

$$\begin{aligned}
\left(\sum_{j \in \mathbb{Z}} \mathfrak{w}_j^p I_j^p \right)^{1/p} &\leq K_0 \cdot K_1^{n-1} \max_{1 \leq r \leq n} \sup_{j \in \mathbb{Z}} \mathfrak{w}_j \left\| \frac{\mathfrak{g}_{n,r}(\cdot; j)}{\sum_{l=1}^n \mathfrak{v}_{kl} \prod_{m \neq l} \mathfrak{u}_{km}} \right\|_{\ell^{p'}(\mathbb{Z}^n)} \left\| \sum_r \sum_l \beta_{j,l,r} \right\|_{\ell_j^p} \\
&\leq K_0 R_0 \cdot (K_1 R_1)^{n-1} \sum_{r,l} \|\beta_{j,l,r}\|_{\ell_j^p} \\
&\leq 4 K_0 R_0 \cdot (2 K_1 R_1)^{n-1} (\|v\|_{p,u} \|\zeta\|_{p,v} + \|v\|_{p,v} \|\zeta\|_{p,u}) \|\zeta\|_{p,u}^{n-2}
\end{aligned}$$

where to pass to the second inequality we used (3.6), while to pass to the third inequality we used the explicit expression of $\beta_{j,l,r}$ and the inequality $\sum_{l,r} 1 \leq 4 \cdot 2^{n-1}$.

Clearly $\left(\sum_j \mathfrak{w}_j^p |\underline{II}_j|^p \right)^{1/p}$ is bounded by the same quantity. It follows that

$$\begin{aligned}
\|\underline{A}^n(|\zeta|)|v\|_{p,\mathfrak{w}} &\leq \|(I_j)_{j \in \mathbb{Z}}\|_{p,\mathfrak{w}} + \|(\underline{II}_j)_{j \in \mathbb{Z}}\|_{p,\mathfrak{w}} \\
&\leq 2^3 K_0 R_0 (2 K_1 R_1)^{n-1} \|\zeta\|_{p,u}^{n-2} (\|v\|_{p,v} \|\zeta\|_{p,u} + \|v\|_{p,u} \|\zeta\|_{p,v}) . \quad (3.42)
\end{aligned}$$

One verifies that $\|\underline{B}^n(\zeta)v\|_{p,\mathfrak{w}}$ admits the same bound (3.42). Thus estimate (3.39) follows. The estimate for (3.40) is obtained similarly putting $\mathfrak{w} = \mathfrak{v} = \mathfrak{u}$. Note that in such a case it is enough to use the sub-multiplicative property of the weight and estimate (3.4).

We prove now (ii). By the results of item (i) we need just to check that $\zeta \mapsto d\Psi(\zeta)^* \in \mathcal{N}_\rho^T(B^{p,\mathfrak{u}} \cap \ell_c^{p,\mathfrak{v}}, \mathcal{L}(\ell_c^{p,\mathfrak{v}}, \ell_c^{p,\mathfrak{w}}))$. However this follows from the identity

$$\underline{d\Psi_3}(|\zeta|)^*|v| = \underline{d\Psi_3}(|\zeta|)^t|v|$$

and the previous result. Finally $\left\langle |\underline{d\Psi_3}^t| \right\rangle_\rho \leq \sum_{n \geq 3} \left\langle |\underline{d\Psi_n}^t| \right\rangle_\rho \leq 2^8 K_0 R_0 K_1^2 R_1^2 \rho^2$. \square

Then one obtains immediately the following

Lemma 3.12. *Fix $1 \leq p \leq 2$ and let $\mathfrak{u} \leq \mathfrak{v} \leq \mathfrak{w}$ be weights fulfilling $(W)_p$ with constants R_0, R_1 . Then for any*

$$0 < \rho < \min\left(\frac{1}{2^4}, \frac{1}{8 K_1 R_1}\right)$$

the map $\Psi_3 \in \mathcal{T}_{\mathfrak{u},\mathfrak{v},\rho}^{\mathfrak{w},3}$ and moreover

$$\|\Psi_3\|_{\mathcal{T}_{\mathfrak{u},\mathfrak{v},\rho}^{\mathfrak{w},3}} \leq 2^{10} K_0 R_0 K_1^2 R_1^2 \rho'^3 .$$

Proof. Let $\rho' = \rho/2$. Then by Lemma 3.10, Lemma 3.11 and Cauchy estimates (2.9),(2.16)

$$\|\Psi_3\|_{\mathcal{T}_{\mathfrak{u},\mathfrak{v},\rho'}^{\mathfrak{w},3}} \leq 2 \langle |\underline{\Psi}_3| \rangle_{2\rho'} + \rho' \left\langle |\underline{d\Psi_3}^*| \right\rangle_{\rho'} \leq 2^{10} K_0 R_0 K_1^2 R_1^2 \rho'^3 ,$$

then we denote again $\rho' \equiv \rho$. \square

We conclude the section with the proof of Theorem 3.3.

Proof of Theorem 3.3. Fix $1 \leq p \leq 2$. By Proposition 3.2(i) the weights $\mathfrak{u} = \mathfrak{v} = \mathfrak{w} = \{1\}_{j \in \mathbb{Z}}$ fulfill $(W)_p$ with $R_0 = 1$ and $1 \leq R_1 \equiv \left(\sum_{k \in \mathbb{Z}} \frac{1}{\langle k \rangle^{p'}} \right)^{1/p'} \leq 2$. Then by Lemma 3.12, $\Psi - \mathbb{1} \in \mathcal{T}_{\mathfrak{u},\mathfrak{v},\rho}^{\mathfrak{w},3}$ for any $\rho < \min(2^{-4}, 2^{-4} K_1^{-1}) \equiv \varrho_*$. Item (i) of Theorem 3.3 follows by Lemma 3.7 (iii). Item (ii) of Theorem 3.3 follows by Lemma 3.12, with $C = 2^{10} K_0 K_1^2$. Item (iii) follows by Lemma 3.7(ii). \square

A Proof of tame Kuksin-Perelman theorem

A.1 Properties of tame majorant analytic functions

In this section we show that the class of tame majorant analytic maps is closed under several operations like composition, inversion and flow-generation, and provide new quantitative estimates which will be used during the proof of Theorem 2.19. In the rest of the section denote by $S := \sum_{n=1}^{\infty} 1/n^2$ and by

$$\mu := 1/e^2 32S \approx 0.0025737 > 2^{-10}. \quad (\text{A.1})$$

Lemma A.1. *Let $F \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \ell_c^{p,w^2})$. Write $F = \sum_{n \geq 0} F^n$ and denote by \tilde{F}^n the symmetric multilinear map associated to F^n . Then each multilinear map \tilde{F}^n is a (p, w^0, w^1, w^2) -tame modulus map in the sense that*

$$\|\tilde{F}^n(\zeta^{(1)}, \dots, \zeta^{(n)})\|_{p,w^2} \leq e^n \frac{|F|_\rho^T}{\rho^{n-1}} \frac{1}{n} \sum_{l=1}^n \|\zeta^{(l)}\|_{p,w^1} \prod_{m \neq l} \|\zeta^{(m)}\|_{p,w^0}, \quad \forall \zeta^{(1)}, \dots, \zeta^{(n)} \in \ell_c^{p,w^1}. \quad (\text{A.2})$$

Proof. By Cauchy formula

$$\tilde{F}^n(\zeta^{(1)}, \dots, \zeta^{(n)}) = \frac{1}{(2\pi i)^n n!} \oint_{|\lambda_1|=\epsilon_1} \dots \oint_{|\lambda_n|=\epsilon_n} \frac{F(\lambda_1 \zeta^{(1)} + \dots + \lambda_n \zeta^{(n)})}{\lambda_1^2 \dots \lambda_n^2} d\lambda_1 \dots d\lambda_n;$$

such formula is well defined provided $\sum_j \lambda_j \zeta^{(j)} \in B^{p,w^0}(\rho)$: this is true e.g. choosing $\epsilon_i = \rho/n \|\zeta^{(i)}\|_{p,w^0} \forall 1 \leq i \leq n$. Then

$$\begin{aligned} \|\tilde{F}^n(\zeta^{(1)}, \dots, \zeta^{(n)})\|_{p,w^2} &\leq \frac{|F|_\rho^T}{n! \epsilon_1 \dots \epsilon_n} \sum_j \epsilon_j \|\zeta^{(j)}\|_{p,w^1} \leq \frac{|F|_\rho^T}{n!} \sum_j \frac{\|\zeta^{(j)}\|_{p,w^1}}{\prod_{l \neq j} \epsilon_l} \\ &\leq \frac{|F|_\rho^T}{\rho^{n-1}} \frac{n^n}{n!} \frac{1}{n} \sum_j \|\zeta^{(j)}\|_{p,w^1} \prod_{l \neq j} \|\zeta^{(l)}\|_{p,w^0} \end{aligned}$$

and the claimed estimate follows. \square

Lemma A.2. *Let $1 \leq p \leq 2$. Let $w^0 \leq w^1 \leq w^2 \leq w^3$ be weights.*

(i) *Let $F \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \ell_c^{p,w^2})$ and $\mathcal{G} \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \mathcal{L}(\ell_c^{p,w^1}, \ell_c^{p,w^2}))$. Define $H(\zeta) = \mathcal{G}(\zeta)F(\zeta)$. Then $H \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \ell_c^{p,w^2})$ and $\langle |H| \rangle_\rho \leq \langle |\mathcal{G}| \rangle_\rho \langle |F| \rangle_\rho$.*

(ii) *Let $\mathcal{H}, \mathcal{G} \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \mathcal{L}(\ell_c^{p,w^1}, \ell_c^{p,w^2}))$. Define $\mathcal{I}(\zeta)v := \mathcal{H}(\zeta)\mathcal{G}(\zeta)v$ for $v \in \ell_c^{p,w^1}$. Then $\mathcal{I} \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \mathcal{L}(\ell_c^{p,w^1}, \ell_c^{p,w^2}))$ and $\langle |\mathcal{I}| \rangle_\rho \leq \langle |\mathcal{H}| \rangle_\rho \langle |\mathcal{G}| \rangle_\rho$.*

Proof. (i) One has that $\mathcal{H}(|\zeta|) \leq \mathcal{G}(|\zeta|)F(|\zeta|)$, which implies immediately that $|H|_\rho \leq |\mathcal{G}|_\rho |F|_\rho$ and $|H|_\rho^T \leq |\mathcal{G}|_\rho^T (|F|_\rho + \rho |F|_\rho^T)$. The claim follows.

(ii) One has that $\mathcal{H}(\zeta)\mathcal{G}(\zeta)v \leq \mathcal{H}(|\zeta|)\mathcal{G}(|\zeta|)|v|$, then the claim follows as above. \square

Lemma A.3. *Let $1 \leq p \leq 2$. Let $w^0 \leq w^1 \leq w^2 \leq w^3$ be weights.*

(i) *Let $G \in \mathcal{N}_\rho(\ell_c^{p,w^0}, \ell_c^{p,w^1})$ with $|G|_\rho \leq \sigma$ and $F \in \mathcal{N}_\sigma(\ell_c^{p,w^1}, \ell_c^{p,w^2})$. Then $F \circ G$ belongs to $\mathcal{N}_\rho(\ell_c^{p,w^0}, \ell_c^{p,w^2})$ and $|F \circ G|_\rho \leq |F|_\sigma$.*

(ii) *Let $G \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \ell_c^{p,w^2})$ with $\langle |G| \rangle_\rho \leq \sigma$ and $F \in \mathcal{N}_\sigma^T(B^{p,w^0} \cap \ell_c^{p,w^2}, \ell_c^{p,w^3})$. Then $F \circ G \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \ell_c^{p,w^3})$ and $\langle |F \circ G| \rangle_\rho \leq \langle |F| \rangle_\sigma$.*

(iii) *Let $G \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \ell_c^{p,w^2})$ with $\langle |G| \rangle_\rho \leq \sigma$ and $\mathcal{E} \in \mathcal{N}_\sigma^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \mathcal{L}(\ell_c^{p,w^1}, \ell_c^{p,w^3}))$. Define $\mathcal{H}(\zeta) = \mathcal{E}(G(\zeta))$. Then $\mathcal{H} \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \mathcal{L}(\ell_c^{p,w^1}, \ell_c^{p,w^3}))$ and $\langle |\mathcal{H}| \rangle_\rho \leq \langle |\mathcal{E}| \rangle_\sigma$.*

Proof. (i) It follows as in [BM16, Lemma A.1].

(ii) Recall that $\underline{F} \circ \underline{G}(|\zeta|) \leq \underline{F} \circ \underline{G}(|\zeta|)$ (cf. [KP10]). Provided $|\underline{G}|_\rho \leq \sigma$ one has $\|\underline{F}(\underline{G}(|\zeta|))\|_{p,w^3} \leq |\underline{F}|_\sigma^T |\underline{G}|_\rho^T \|\zeta\|_{p,w^1}$, hence $|\underline{F} \circ \underline{G}|_\rho^T \leq |\underline{F} \circ \underline{G}|_\rho^T \leq |\underline{F}|_\sigma^T |\underline{G}|_\rho^T$. Thus

$$\langle |\underline{F} \circ \underline{G}| \rangle_\rho \equiv |\underline{F} \circ \underline{G}|_\rho + \rho |\underline{F} \circ \underline{G}|_\rho^T \leq |\underline{F}|_\sigma + |\underline{F}|_\sigma^T |\underline{G}|_\rho^T \rho \leq |\underline{F}|_\sigma + |\underline{F}|_\sigma^T \sigma \equiv \langle |\underline{F}| \rangle_\sigma.$$

(iii) First $\underline{H}(|\zeta|)|v| \leq \underline{E}(\underline{G}(|\zeta|))|v|$, which implies $|\underline{H}|_\rho \leq |\underline{E}|_\sigma$. Furthermore, using also $|\underline{G}|_\rho^T \leq \sigma/\rho$,

$$\begin{aligned} \|\underline{H}(|\zeta|)|v|\|_{p,w^3} &\leq |\underline{E}|_\sigma^T (\|\underline{G}(|\zeta|)\|_{p,w^1} \|v\|_{p,w^0} + \sigma \|v\|_{p,w^1}) \leq |\underline{E}|_\sigma^T (|\underline{G}|_\rho^T \|\zeta\|_{p,w^1} \|v\|_{p,w^0} + \sigma \|v\|_{p,w^1}) \\ &\leq |\underline{E}|_\sigma^T \frac{\sigma}{\rho} (\|\zeta\|_{p,w^1} \|v\|_{p,w^0} + \rho \|v\|_{p,w^1}) \end{aligned}$$

therefore $|\underline{H}|_\rho^T \leq |\underline{E}|_\sigma^T \frac{\sigma}{\rho}$. Finally $\langle |\underline{H}| \rangle_\rho = |\underline{H}|_\rho + \rho |\underline{H}|_\rho^T \leq |\underline{E}|_\sigma + \sigma |\underline{E}|_\sigma^T \equiv \langle |\underline{E}| \rangle_\sigma$. \square

Lemma A.4. Fix $1 \leq p \leq 2$ and weights $w^0 \leq w^1 \leq w^2$.

(i) Let $F \in \mathcal{N}_\rho(\ell_c^{p,w^0}, \ell_c^{p,w^0})$, $F = O(\zeta^N)$ for some $N \geq 2$ and $|\underline{F}|_\rho \leq \rho/e$. Then the map $\mathbb{1} + F$ is invertible in $B^{p,w^0}(\mu\rho)$, μ as in (A.1). Moreover there exists $G \in \mathcal{N}_{\mu\rho}(\ell_c^{p,w^0}, \ell_c^{p,w^0})$, $G = O(\zeta^N)$, such that $(\mathbb{1} + F)^{-1} = \mathbb{1} - G$, and

$$|\underline{G}|_{\mu\rho} \leq \frac{|\underline{F}|_\rho}{8}. \quad (\text{A.3})$$

(ii) Assume that $F \in \mathcal{N}_\rho^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \ell_c^{p,w^2})$, $F = O(\zeta^N)$ for some $N \geq 2$ and $\langle |\underline{F}| \rangle_\rho \leq \rho/e$, then $G \in \mathcal{N}_{\mu\rho}^T(B^{p,w^0} \cap \ell_c^{p,w^1}, \ell_c^{p,w^2})$ and

$$\langle |\underline{G}| \rangle_{\mu\rho} \leq \frac{\langle |\underline{F}| \rangle_\rho}{8}. \quad (\text{A.4})$$

Proof. Item (i) follows as in [BM16, Lemma A.2]. We prove item (ii). Following the scheme of [BM16], G is given by the power series $G = \sum_{n \geq 2} G^n$, where the homogeneous polynomial $G^n = 0$ for $1 < n < N$ and

$$G^n(\zeta) = \sum_{r=2}^n \sum_{k_1 + \dots + k_r = n} \tilde{F}^r(G^{k_1}(\zeta), \dots, G^{k_r}(\zeta)), \quad \forall n \geq N. \quad (\text{A.5})$$

In the formula above $k_1, \dots, k_r \in \mathbb{N}$, and we wrote $F = \sum_{r \geq N} F^r$, where F^r is a homogeneous polynomial of degree r and \tilde{F}^r is its associated multilinear map (see (2.3)). Moreover we write $G^1(\zeta) := \zeta$. We show now that the formal series $G = \sum_{n \geq N} G^n$ with G^n defined by (A.5) is a tame majorant analytic map. Note that

$$\underline{G}^n(|\zeta|) \leq \sum_{r=N}^n \sum_{k_1 + \dots + k_r = n} \underline{\tilde{F}}^r(\underline{G}^{k_1}(|\zeta|), \dots, \underline{G}^{k_r}(|\zeta|)). \quad (\text{A.6})$$

We prove that there exists a constant $A > 0$ such that

$$\|\underline{G}^n(|\zeta|)\|_{p,w^2} \leq \frac{|\underline{F}|_\rho^T}{8Sn^2} A^{n-1} \|\zeta\|_{p,w^0}^{n-1} \|\zeta\|_{p,w^1}, \quad \forall n \geq N, \quad (\text{A.7})$$

$$\|\underline{G}^n(|\zeta|)\|_{p,w^0} \leq \frac{|\underline{F}|_\rho}{8Sn^2} A^n \|\zeta\|_{p,w^0}^n, \quad \forall n \geq N \quad (\text{A.8})$$

The proof is by induction on n . For $n = N$, by (A.5) it follows that $G^N(\zeta) = \tilde{F}^N(\zeta, \dots, \zeta)$. Using also Lemma A.1 one has

$$\|\underline{G}^N(|\zeta|)\|_{p,w^2} \leq e^N \frac{|\underline{F}|_\rho^T}{\rho^{N-1}} \|\zeta\|_{p,w^0}^{N-1} \|\zeta\|_{p,w^1},$$

thus (A.7) holds for $n = N$ with $A = e^2 32S/\rho$. The proof of (A.8) is analogous, and we skip it. We prove now the inductive step $n - 1 \rightsquigarrow n$. Assume therefore that (A.7), (A.8) hold up to order $n - 1$. Then one has

$$\begin{aligned}
\|G^n(|\zeta|)\|_{p,w^2} &\leq \sum_{r=N}^n \sum_{k_1+\dots+k_r=n} \frac{|F|_\rho^T}{\rho^{r-1}} \frac{e^r}{r} \sum_{\ell=1}^r \|G^{k_\ell}(|\zeta|)\|_{p,w^2} \prod_{m \neq \ell} \|G^{k_m}(|\zeta|)\|_{p,w^0} \\
&\leq \frac{|F|_\rho^T}{8S} A^{n-1} \|\zeta\|_{p,w^0}^{n-1} \|\zeta\|_{p,w^1} e |F|_\rho^T \sum_{r=N}^n \left(\frac{e|F|_\rho}{8S\rho} \right)^{r-1} \sum_{k_1+\dots+k_r=n} \frac{1}{k_1^2 \dots k_r^2} \\
&\leq \frac{|F|_\rho^T}{8Sn^2} A^{n-1} \|\zeta\|_{p,w^0}^{n-1} \|\zeta\|_{p,w^1} e |F|_\rho^T \sum_{r=1}^\infty \left(\frac{e|F|_\rho}{2\rho} \right)^r \\
&\leq \frac{|F|_\rho^T}{8Sn^2} A^{n-1} \|\zeta\|_{p,w^0}^{n-1} \|\zeta\|_{p,w^1}
\end{aligned}$$

where in the first inequality we used $w^1 \leq w^2$, in the second the inductive assumption and in the last we used the hypothesis $|F|_\rho + \rho|F|_\rho^T \leq \rho/e$. Finally to pass from the second to the third line we used the following standard inequality (see e.g. [BM16, Lemma A.5])

$$n^2 \sum_{k_1+\dots+k_r=n} \frac{1}{k_1^2 \dots k_r^2} \leq (4S)^{r-1}, \quad n \geq 1. \quad (\text{A.9})$$

In such a way we proved (A.7). The proof of (A.8) is analogous (for details see [BM16]). Finally from (A.7) one has $|G|_{\mu\rho}^T \leq \sum_{n \geq N} \frac{|F|_\rho^T}{8Sn^2} (A\mu\rho)^{n-1} \leq \frac{|F|_\rho^T}{8}$ choosing $\mu\rho = 1/A = \rho/e^2 32S$. \square

Next we have closedness of the classes $\mathcal{A}_{w^0,\rho}^N$ and $\mathcal{T}_{w^0,w^1,\rho}^{w^2,N}$ under different operations.

Lemma A.5. *Fix $1 \leq p \leq 2$, weights $w^0 \leq w^1 \leq w^2$, $\mathbb{N} \ni N \geq 2$ and let μ be as in (A.1). Then the following holds true:*

(i) *If $F \in \mathcal{T}_{w^0,w^1,\rho}^{w^2,N}$ and $G \in \mathcal{T}_{w^0,w^1,\mu\rho}^{w^2,N}$ with $\|G\|_{\mathcal{T}_{w^0,w^1,\mu\rho}^{w^2,N}} < \frac{\rho\mu}{e}$, then $H := F(\zeta + G(\zeta)) \in \mathcal{T}_{w^0,w^1,\mu\rho}^{w^2,N}$ and*

$$\|H\|_{\mathcal{T}_{w^0,w^1,\mu\rho}^{w^2,N}} \leq 2\|F\|_{\mathcal{T}_{w^0,w^1,\rho}^{w^2,N}}.$$

(ii) *If $F \in \mathcal{T}_{w^0,w^1,\rho}^{w^2,N}$ and $\|F\|_{\mathcal{T}_{w^0,w^1,\rho}^{w^2,N}} \leq \rho/e$, then $(1 + F)^{-1} = 1 + G$ with $G \in \mathcal{T}_{w^0,w^1,\mu\rho}^{w^2,N}$ with*

$$\|G\|_{\mathcal{T}_{w^0,w^1,\mu\rho}^{w^2,N}} \leq 2\|F\|_{\mathcal{T}_{w^0,w^1,\rho}^{w^2,N}}. \quad (\text{A.10})$$

(iii) *If $F \in \mathcal{T}_{w^0,w^1,\rho}^{w^2,N}$, then $H(\zeta) := dF(\zeta)\zeta$ is in the class $\mathcal{T}_{w^0,w^1,\mu\rho}^{w^2,N}$ and*

$$\|H\|_{\mathcal{T}_{w^0,w^1,\mu\rho}^{w^2,N}} \leq 2\|F\|_{\mathcal{T}_{w^0,w^1,\rho}^{w^2,N}}.$$

(iv) *If $F^0, G^0 \in \mathcal{T}_{w^0,w^1,\rho}^{w^2,N}$ with $\|F^0\|_{\mathcal{T}_{w^0,w^1,\rho}^{w^2,N}} \leq \frac{\rho}{e}$, then $H_0 := dG^0(\zeta)^*(F^0(\zeta)) \in \mathcal{T}_{w^0,w^1,\rho/2}^{w^2,N}$ and*

$$\|H_0\|_{\mathcal{T}_{w^0,w^1,\rho/2}^{w^2,N}} \leq \|G^0\|_{\mathcal{T}_{w^0,w^1,\rho}^{w^2,N}} \frac{4}{\rho} \|F^0\|_{\mathcal{T}_{w^0,w^1,\rho/2}^{w^2,N}} \leq 2\|G^0\|_{\mathcal{T}_{w^0,w^1,\rho}^{w^2,N}}.$$

(v) *If $F^0, G^0 \in \mathcal{T}_{w^0,w^1,\rho}^{w^2,N}$ with $\|F^0\|_{\mathcal{T}_{w^0,w^1,\rho}^{w^2,N}} \leq \frac{\rho}{e}$, then $H_0 := dG^0(\zeta)F^0(\zeta) \in \mathcal{T}_{w^0,w^1,\rho/2}^{w^2,N}$ and*

$$\|H_0\|_{\mathcal{T}_{w^0,w^1,\rho/2}^{w^2,N}} \leq \|G^0\|_{\mathcal{T}_{w^0,w^1,\rho}^{w^2,N}} \frac{4}{\rho} \|F^0\|_{\mathcal{T}_{w^0,w^1,\rho/2}^{w^2,N}} \leq 2\|G^0\|_{\mathcal{T}_{w^0,w^1,\rho}^{w^2,N}}.$$

Finally all the results hold true replacing everywhere $\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}$ with $\mathcal{A}_{w^0, \rho}^N$.

Proof. Again we prove only the tame estimates, since the results for the class $\mathcal{A}_{w^0, \rho}^N$ are already proved in [BM16, Lemma A.3].

- (i) One has $\mathbb{1} + G \in \mathcal{N}_{\mu\rho}^T(B^{p, w^0} \cap \ell_c^{p, w^1}, \ell_c^{p, w^1})$ with $\langle |\mathbb{1} + \underline{G}| \rangle_{\mu\rho} \leq \rho$, then by Lemma A.3

$$\langle |\underline{H}| \rangle_{\mu\rho} \leq \langle |\underline{F}| \rangle_{\rho}.$$

Now $dH(\zeta) = dF(\zeta + G(\zeta))(\mathbb{1} + dG(\zeta))$ therefore

$$\underline{dH}(|\zeta|) \leq \underline{dF}(|\zeta| + \underline{G}(|\zeta|)) + \underline{dF}(|\zeta| + \underline{G}(|\zeta|))\underline{dG}(|\zeta|).$$

Then exploiting Lemma A.3 (iii) and Lemma A.2

$$\langle |\underline{dH}| \rangle_{\mu\rho} \leq \langle |\underline{dF}| \rangle_{\rho}(1 + \langle |\underline{dG}| \rangle_{\mu\rho}).$$

Finally the adjoint $dH(\zeta)^* = (\mathbb{1} + dG(\zeta))^*dF(\zeta + G(\zeta))^*$, thus analogously one finds

$$\langle |\underline{dH}^*| \rangle_{\mu\rho} \leq \langle |\underline{dF}^*| \rangle_{\rho}(1 + \langle |\underline{dG}^*| \rangle_{\mu\rho}).$$

Therefore since $\mu\rho(\langle |\underline{dG}| \rangle_{\mu\rho} + \langle |\underline{dG}^*| \rangle_{\mu\rho}) \leq \|G\|_{\mathcal{T}_{w^0, w^1, \mu\rho}^{w^2, N}} \leq \mu\rho/e$, one has

$$\begin{aligned} \|H\|_{\mathcal{T}_{w^0, w^1, \mu\rho}^{w^2, N}} &\leq \langle |\underline{F}| \rangle_{\mu\rho} + \langle |\underline{dF}| \rangle_{\rho}(\mu\rho + \mu\rho\langle |\underline{dG}| \rangle_{\mu\rho}) + \langle |\underline{dF}^*| \rangle_{\rho}(\mu\rho + \mu\rho\langle |\underline{dG}^*| \rangle_{\mu\rho}) \\ &\leq 2\|F\|_{\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}}. \end{aligned}$$

- (ii) By Lemma A.4 we know that $\langle |\underline{G}| \rangle_{\mu\rho} \leq \frac{\langle |\underline{F}| \rangle_{\rho}}{8}$. Differentiating the identity $G = F \circ (\mathbb{1} - G)$ one gets

$$dG(\zeta) = [\mathbb{1} + dF(\zeta - G(\zeta))]^{-1}dF(\zeta - G(\zeta)),$$

thus by Lemma A.3(iii), Lemma A.2 and the assumption $\|F\|_{\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}} \leq \rho/e$, it follows that

$$\langle |\underline{dG}| \rangle_{\mu\rho} \leq \langle |\underline{dF}| \rangle_{\rho} \sum_{k \geq 0} \langle |\underline{dF}| \rangle_{\rho}^k \leq \frac{e}{e-1} \langle |\underline{dF}| \rangle_{\rho}.$$

Analogously $\langle |\underline{dG}^*| \rangle_{\mu\rho} \leq \frac{e}{e-1} \langle |\underline{dF}^*| \rangle_{\rho}$. Estimate (A.10) then follows.

- (iii) Clearly $\langle |\underline{H}| \rangle_{\rho} \leq \langle |\underline{dF}| \rangle_{\rho}$. Then $dH(\zeta)v = dF(\zeta)v + d^2F(\zeta)(v, \zeta)$, and the claim follows arguing as in item (i) and using Cauchy estimate (see also the proof of (iv)).

- (iv) Consider first $\underline{H}_0(|\zeta|)$. By Lemma A.2(i), $\langle |\underline{H}_0| \rangle_{\rho} \leq \langle |\underline{dG}^{0*}| \rangle_{\rho} \langle |\underline{F}^0| \rangle_{\rho}$. Consider now $dH_0(\zeta)v = dG^0(\zeta)^*dF^0(\zeta)v + d_{\zeta}(dG^0(\zeta)^*U)v$, $U = F^0(\zeta)$. One has

$$\underline{dH}_0(|\zeta|)|v| \leq \underline{dG}^0(|\zeta|)^*\underline{dF}^0(|\zeta|)|v| + d_{|\zeta|}(\underline{dG}^0(|\zeta|)^*\underline{U})|v|, \quad \underline{U} = \underline{F}^0(|\zeta|).$$

Using also the Cauchy estimates (2.16), one has $\langle |\underline{dH}_0| \rangle_{\rho/2} \leq \frac{2}{\rho} \langle |\underline{dG}^{0*}| \rangle_{\rho} \|F^0\|_{\mathcal{T}_{w^0, w^1, \rho/2}^{w^2, N}}$.

Finally in order to estimate $\underline{dH}_0(|\zeta|)^*$ remark that (see [KP10, Lemma 3.6]) $dH_0(\zeta)^*v = dF^0(\zeta)^*dG^0(\zeta)v + d_{\zeta}(dG^0(\zeta)^*U)v$, thus

$$\underline{dH}_0(|\zeta|)^*|v| \leq \underline{dF}^0(|\zeta|)^*\underline{dG}^0(|\zeta|)|v| + d_{|\zeta|}(\underline{dG}^0(|\zeta|)^*\underline{U})|v|.$$

The $\langle |\cdot| \rangle_{\rho/2}$ norm of first term in the r.h.s. is estimated by $\langle |\underline{dF}^{0*}| \rangle_{\rho/2} \langle |\underline{dG}^0| \rangle_{\rho/2}$. To estimate the second term we use Cauchy formula (2.15), obtaining that its $\langle |\cdot| \rangle_{\rho/2}$ norm is controlled by $\frac{2}{\rho} \langle |\underline{dG}^{0*}| \rangle_{\rho} \langle |\underline{F}^0| \rangle_{\rho/2}$. The claim follows.

(v) The proof is similar to (iv), and we skip it. \square

Finally we analyze the flow generated by a vector field of class $\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}$. Given a time dependent vector field $Y_t(v)$, consider the differential equation

$$\begin{cases} \dot{u}(t) = Y_t(u(t)) \\ u(0) = \zeta \in \ell_c^{p, w^0} \end{cases} \quad . \quad (\text{A.11})$$

We will denote by $\phi^t(\zeta)$ the corresponding flow map whose existence and properties are given in the next lemma.

Lemma A.6. *Assume that $[0, 1] \ni t \mapsto Y_t \in \mathcal{T}_{w^0, w^1, \rho}^{w^2, N}$ is continuous and $\sup_{t \in [0, 1]} \|Y_t\|_{\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}} \leq \rho/e$; then for each $t \in [0, 1]$, $\phi^t - \mathbb{1} \in \mathcal{T}_{w^0, w^1, \mu\rho}^{w^2, N}$ and*

$$\|\phi^t - \mathbb{1}\|_{\mathcal{T}_{w^0, w^1, \mu\rho}^{w^2, N}} \leq 2 \sup_{t \in [0, 1]} \|Y_t\|_{\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}} \quad . \quad (\text{A.12})$$

The same holds true if the class $\mathcal{T}_{w^0, w^1, \mu\rho}^{w^2, N}$ is replaced everywhere by $\mathcal{A}_{w^0, \mu\rho}^N$.

Proof. The claim for the class $\mathcal{A}_{w^0, \mu\rho}^N$ follows as in [BM16], thus we consider only the tame class. We look for a solution $u(t, \zeta) = \sum_{j \geq 1} u^j(t, \zeta)$ in power series of ζ , with $u^j(t, \zeta)$ a homogeneous polynomial of degree j in ζ . Expanding the vector field $Y_t(\zeta) = \sum_{r \geq N} Y_t^r(\zeta)$ in Taylor series, one obtains the recursive formula for the solution

$$u^1(t, \zeta) = \zeta, \quad u^n(t, \zeta) = \sum_{r=2}^n \sum_{k_1 + \dots + k_r = n} \int_0^t \tilde{Y}_s^r(u^{k_1}(s, \zeta), \dots, u^{k_r}(s, \zeta)) ds \quad \forall n \geq 2, \quad (\text{A.13})$$

where \tilde{Y}_s^r is the multilinear map associated to Y_s^r (see (2.3)). Arguing as in the proof of (A.4) one gets that $u^n(t, \zeta) = 0$ if $1 < n < N$, while

$$\|\underline{u}^n(t, \zeta)\|_{p, w^2} \leq \frac{\sup_{t \in [0, 1]} |\underline{Y}_t|_\rho^T}{8Sn^2} A^{n-1} \|\zeta\|_{p, w^0}^{n-1} \|\zeta\|_{p, w^1} \quad \forall n \geq N, \quad (\text{A.14})$$

$$\|\underline{u}^n(t, \zeta)\|_{p, w^0} \leq \frac{\sup_{t \in [0, 1]} |\underline{Y}_t|_\rho}{8Sn^2} A^n \|\zeta\|_{p, w^0}^n \quad \forall n \geq N, \quad (\text{A.15})$$

with $A = \frac{e^2}{\rho} 32S$, from which it follows that $\langle |\phi^t - \mathbb{1}| \rangle_{\mu\rho} \leq \sup_{t \in [0, 1]} \langle |\underline{Y}_t| \rangle_\rho / 8$. We come to the estimate of the differential of $u(t, \zeta)$ and of its adjoint. To do this, remark that $du(t, \zeta)v$ is the solution of the linearized equation

$$\dot{w}(t) = dY_t(u(t, \zeta))w(t), \quad w(0) = v,$$

whose solution can be written by Picard iteration as

$$du(t, \zeta)v = v + \sum_{n=1}^{\infty} \int_0^t \cdots \int_0^{t_{n-1}} dY_{t_1}(u(t_1, \zeta)) \cdots dY_{t_n}(u(t_n, \zeta)) v dt_n \dots dt_1.$$

The series is absolutely and uniformly convergent for $\|\zeta\|_{p, w^0}$ sufficiently small; moreover

$$(\underline{du}(t, |\zeta|) - \mathbb{1})|v| \leq \sum_{n=1}^{\infty} \int_0^t \cdots \int_0^{t_{n-1}} \underline{dY}_{t_1}(\underline{u}(t_1, |\zeta|)) \cdots \underline{dY}_{t_n}(\underline{u}(t_n, |\zeta|)) |v| dt_n \dots dt_1, \quad (\text{A.16})$$

so one has

$$\langle |\underline{du} - \mathbb{1}| \rangle_{\mu\rho} \leq \sum_{n=1}^{\infty} \left(\sup_{t \in [0, 1]} \langle |\underline{dY}_t| \rangle_\rho \right)^n \leq \frac{e}{e-1} \sup_{t \in [0, 1]} \langle |\underline{dY}_t| \rangle_\rho$$

Finally one has to estimate $[\underline{du}^n]^*$, but this is done by simply taking the adjoint of (A.16) and estimate the r.h.s. using the bounds on $\sup_{t \in [0, 1]} \langle |\underline{dY}_t^*| \rangle_\rho$. \square

A.2 Proof Theorem 2.19

The map $\tilde{\Psi}$ of Theorem 2.19 will be constructed in two steps based on the Darboux theorem. Such a theorem states that in order to construct a coordinate transformation ψ transforming the closed nondegenerate form Ω_1 into a closed nondegenerate form Ω_0 , then it is convenient to look for ψ as the time 1 flow ψ^t of a time-dependent vector field Y^t . To construct Y^t one defines $\Omega_t := \Omega_0 + t(\Omega_1 - \Omega_0)$ and imposes that

$$0 = \frac{d}{dt} \Big|_{t=0} \psi^{t*} \Omega_t = \psi^{t*} (\mathcal{L}_{Y^t} \Omega_t + \Omega_1 - \Omega_0) = \psi^{t*} (d(Y^t \lrcorner \Omega_t) + d(\alpha_1 - \alpha_0))$$

where α_1, α_0 are potential forms for Ω_1 and Ω_0 (namely $d\alpha_i = \Omega_i$, $i = 0, 1$) and \mathcal{L}_{Y^t} is the Lie derivative of Y^t . Then one gets

$$Y^t \lrcorner \Omega_t + \alpha_1 - \alpha_0 = df \quad (\text{A.17})$$

for each f smooth; then, if Ω_t is nondegenerate, this defines Y^t . If Y^t generates a flow ψ^t defined up to time 1, the map $\psi := \psi^t|_{t=1}$ satisfies $\psi^* \Omega_1 = \Omega_0$ (see also [BM16b] for another application of Darboux theorem in a different model).

We prove the Kuksin-Perelman theorem only in the case $\Psi^0 \in \mathcal{T}_{w^0, w^1, \rho}^{w^2, N}$. In case $\Psi^0 \in \mathcal{A}_{w^0, \rho}^N$, then it is sufficient to replace in the following the class $\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}$ with the class $\mathcal{A}_{w^0, \rho}^N$ (thanks to the results of Section A.1). Actually the proof is exactly as the one of [BM16]; therefore we only state the main lemmas with the new classes $\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}$, while for the proofs we refer the reader to the corresponding proofs in [BM16].

From now on we will work always in the real subspace $\ell_r^{p, w}$, so abusing notation, we will write just $\xi \equiv (\xi, \bar{\xi})$. Correspondingly, for a map $F : B_r^{p, w}(\rho) \rightarrow \ell_r^{p, w}$ we will write simply $v = F(\xi)$ rather than $(v, \bar{v}) \equiv F(\xi, \bar{\xi})$. Furthermore we fix a real $1 \leq p \leq 2$ and weights $w^0 \leq w^1 \leq w^2$. We will meet maps in $\mathcal{T}_{w^0, w^1, \rho}^{w^2, N}$; to simplify notation, we will write only \mathcal{T}_ρ if nothing else is specified.

We recall the setup of [BM16]. A non-constant symplectic form Ω is represented through a linear skew-symmetric invertible operator E_Ω as follows:

$$\Omega(v)(\xi^{(1)}; \xi^{(2)}) = \langle E_\Omega(v) \xi^{(1)}; \xi^{(2)} \rangle, \quad \forall \xi^{(1)}, \xi^{(2)} \in T_v \ell_r^{p, w} \simeq \ell_r^{p, w}, \quad \forall v \in \ell_r^{p, w}. \quad (\text{A.18})$$

Here the scalar product is the one defined in (1.7). We denote by $\{F, G\}_\Omega$ the Poisson bracket with respect to $\Omega : \{F, G\}_\Omega := \langle \nabla F, J_\Omega \nabla G \rangle$, $J_\Omega := E_\Omega^{-1}$. Similarly we will represent 1-forms through the vector field A such that

$$\alpha(v)(\xi) = \langle A(v), \xi \rangle, \quad \forall v \in T_v \ell_r^{p, w}, \quad \forall v \in \ell_r^{p, w}. \quad (\text{A.19})$$

Consider the Hamiltonian vector fields $X_{-I_l}^0$ of the functions $I_l \equiv \frac{|\xi_l|^2}{2}$ through the symplectic form ω_0 ; they are given by

$$[X_{-I_l}^0(\xi)]_j = \delta_{l,j} i \xi_l, \quad \forall l, j \in \mathbb{Z}, \quad \forall \xi \in \ell_r^{p, w}. \quad (\text{A.20})$$

For every $l \in \mathbb{Z}$ the corresponding flow $\phi_l^t \equiv \phi_{X_{-I_l}^0}^t$ is given by

$$\phi_l^t(\xi) = (\dots, \xi_{l-1}, e^{it} \xi_l, \xi_{l+1}, \dots) .$$

Remark that the map ϕ_l^t is linear in ξ , 2π periodic in t and its adjoint satisfies $(\phi_l^t)^* = \phi_l^{-t}$. Given a k -form α on $\ell_r^{p, w}$ ($k \geq 0$), we define its average by

$$M_l \alpha(\xi) = \frac{1}{2\pi} \int_0^{2\pi} ((\phi_l^t)^* \alpha)(\xi) dt, \quad l \in \mathbb{Z}, \quad \text{and} \quad M \alpha(\xi) = \int_{\mathbb{T}^\infty} [(\phi^\theta)^* \alpha] d\theta \quad (\text{A.21})$$

where \mathbb{T}^∞ is the infinite dimensional torus, the map $\phi^\theta = (\dots \circ \phi_{-1}^{\theta_{-1}} \circ \phi_0^{\theta_0} \circ \phi_2^{\theta_2} \dots)$ and $d\theta$ is the Haar measure on \mathbb{T}^∞ .

Remark A.7. In the particular cases of 1 and 2-forms it is useful to compute the average in term of the representations (A.18) and (A.19). Thus, for $v, \xi^{(1)}, \xi^{(2)} \in \ell_r^{p, w}$, if

$$\alpha(v) \xi^{(1)} = \langle A(v); \xi^{(1)} \rangle, \quad \omega(v)(\xi^{(1)}, \xi^{(2)}) = \langle E(v) \xi^{(1)}; \xi^{(2)} \rangle,$$

one has

$$(M\alpha)(v)\xi^{(1)} = \langle (MA)(v); \xi^{(1)} \rangle, \quad \text{with} \quad MA(v) = \int_{\mathbb{T}^\infty} \phi^{-\theta} A(\phi^\theta(v)) d\theta \quad (\text{A.22})$$

and

$$(M\omega)(v)(\xi^{(1)}, \xi^{(2)}) = \langle (ME)(v)\xi^{(1)}; \xi^{(2)} \rangle, \quad \text{with} \quad ME(v) = \int_{\mathbb{T}^\infty} \phi^{-\theta} E(\phi^\theta(v)) \phi^\theta d\theta. \quad (\text{A.23})$$

Remark A.8. The operator M commutes with the differential operator d and the rotations ϕ^θ . In particular $MA(v)$ and $ME(v)$ as in (A.22), (A.23) satisfy

$$\phi^\theta MA(v) = MA(\phi^\theta v), \quad \phi^\theta ME(v)\xi = ME(\phi^\theta v)\phi^\theta \xi, \quad \forall \theta \in \mathbb{T}^\infty.$$

Remark A.9. By condition (2.21) and (A.1), one has

$$\epsilon_1 < \mu^6 \rho. \quad (\text{A.24})$$

Define $\omega_1 := (\Psi^{-1})^* \omega_0$, and let E_{ω_1} be the operator representing the symplectic form ω_1 .

Lemma A.10. Let $\Phi := \Psi^{-1}$ and ω_1 be as above. Assume that $\epsilon_1 \leq \rho/e$. Then the following holds:

(i) $E_{\omega_1} = i + \Upsilon_{\omega_1}$, with $\Upsilon_{\omega_1} \in \mathcal{N}_{\mu\rho}^T(B^{p,w^0} \cap \ell_r^{p,w^1}, \mathcal{L}(\ell_r^{p,w^1}, \ell_r^{p,w^2}))$ and

$$\left\langle |\Upsilon_{\omega_1}| \right\rangle_{\mu\rho} \leq \frac{8\epsilon_1}{\mu\rho}. \quad (\text{A.25})$$

Furthermore Υ_{ω_1} is antisymmetric, $\Upsilon_{\omega_1}(\xi)^* = -\Upsilon_{\omega_1}(\xi)$.

(ii) Define

$$W_{\omega_1}(\xi) := \int_0^1 \Upsilon_{\omega_1}(t\xi) t\xi dt, \quad (\text{A.26})$$

then $W_{\omega_1} \in \mathcal{T}_{\mu^3\rho}$ and $\|W_{\omega_1}\|_{\mathcal{T}_{\mu^3\rho}} \leq 8\epsilon_1$. Moreover the 1-form $\alpha_{W_{\omega_1}} := \langle W_{\omega_1}; \cdot \rangle$ satisfies $d\alpha_{W_{\omega_1}} = \omega_1 - \omega_0$.

Proof. The proof follows [BM16, Lemma 2.16], using the results of Lemma A.5 and A.3. \square

Remark A.11. One has $M\alpha_{\omega_1} - \alpha_0 = M\alpha_{W_{\omega_1}} = \langle MW_{\omega_1}, \cdot \rangle$ and $\|MW_{\omega_1}\|_{\mathcal{T}_{\mu^3\rho}} \leq \|W_{\omega_1}\|_{\mathcal{T}_{\mu^3\rho}}$.

We are ready now for the first step.

Lemma A.12. There exists a map $\hat{\psi} : B_r^{p,w^0}(\mu^5\rho) \rightarrow \ell_r^{p,w^0}$ such that $(1 - \hat{\psi}) \in \mathcal{T}_{\mu^5\rho}$ and

$$\|1 - \hat{\psi}\|_{\mathcal{T}_{\mu^5\rho}} \leq 2^5 \epsilon_1. \quad (\text{A.27})$$

Moreover $\hat{\psi}$ satisfies the following properties:

(i) $\hat{\psi}$ commutes with the rotations ϕ^θ , namely $\phi^\theta \hat{\psi}(\xi) = \hat{\psi}(\phi^\theta \xi)$ for every $\theta \in \mathbb{T}^\infty$.

(ii) Denote $\hat{\omega}_1 := \hat{\psi}^* \omega_1$, then $M\hat{\omega}_1 = \omega_0$.

Proof. It follows as in [BM16, Lemma 2.18]. We apply the Darboux procedure described at the beginning of this section with $\Omega_0 = \omega_0$ and $\Omega_1 = M\omega_1$. Then Ω_t is represented by the operator $\hat{E}_{\omega_1}^t := (i + t(ME_{\omega_1} - i))$. Write equation (A.17), with $f \equiv 0$, in terms of the operators defining the symplectic forms, getting the equation $\hat{E}_{\omega_1}^t \hat{Y}^t = -MW_{\omega_1}$ (see also Remark A.11). This equation can be solved by inverting the operator $\hat{E}_{\omega_1}^t$ by Neumann series:

$$\hat{Y}^t := -(i + tM\Upsilon_{\omega_1})^{-1} MW_{\omega_1}. \quad (\text{A.28})$$

In order to estimate it, we expand the r.h.s. of (A.28) in Neumann series and estimate each piece. First note that for any $G \in \mathcal{T}_{\mu^3\rho}$, by Lemma A.5, one has

$$\|\Upsilon_{\omega_1} G\|_{\mathcal{T}_{\mu^3\rho/2}} \leq \frac{1}{2} \|G\|_{\mathcal{T}_{\mu^3\rho/2}} ;$$

then by induction one has that $\|[\Upsilon_{\omega_1}]^k MW_{\omega_1}\|_{\mathcal{T}_{\mu^3\rho/2}} \leq 2^{-k} \|MW_{\omega_1}\|_{\mathcal{T}_{\mu^3\rho}}$. Therefore the Neumann series converges, \hat{Y}^t is of class $\mathcal{T}_{\mu^4\rho}$ and fulfills

$$\sup_{t \in [0,1]} \|\hat{Y}^t\|_{\mathcal{T}_{\mu^4\rho}} \leq 2 \|MW_{\omega_1}\|_{\mathcal{T}_{\mu^3\rho}} \leq 2^4 \epsilon_1 . \quad (\text{A.29})$$

By Lemma A.6 the vector field \hat{Y}^t generates a flow $\hat{\psi}^t : B_r^{p,w^0}(\mu^5\rho) \rightarrow \ell_r^{p,w^0}$ such that $\hat{\psi}^t - \mathbb{1}$ is of class $\mathcal{T}_{\mu^5\rho}$ and satisfies

$$\|\hat{\psi}^t - \mathbb{1}\|_{\mathcal{T}_{\mu^5\rho}} \leq 2 \sup_{t \in [0,1]} \|\hat{Y}^t\|_{\mathcal{T}_{\mu^4\rho}} \leq 2^5 \epsilon_1 .$$

Therefore the map $\hat{\psi} \equiv \hat{\psi}^t|_{t=1}$ exists, satisfies the claimed estimate (A.27) and furthermore $\hat{\psi}^* M\omega_1 = \omega_0$.

Item (i) is a geometric property and it follows exactly as in [BM16, Lemma 2.18]. \square

The analytic properties of the symplectic form $\hat{\omega}_1$ can be studied in the same way as in Lemma A.10; we get therefore the following corollary:

Corollary A.13. *Denote by $E_{\hat{\omega}_1}$ the symplectic operator describing $\hat{\omega}_1 = \hat{\psi}^* \omega_1$. Then*

$$(i) \ E_{\hat{\omega}_1} = \mathbb{I} + \Upsilon_{\hat{\omega}_1}, \text{ with } \Upsilon_{\hat{\omega}_1} \in \mathcal{N}_{\mu^5\rho}^T(B_r^{p,w^0} \cap \ell_r^{p,w^1}, \mathcal{L}(\ell_r^{p,w^1}, \ell_r^{p,w^2})) \text{ and } \left\langle |\Upsilon_{\hat{\omega}_1}| \right\rangle_{\mu^5\rho} \leq 2^7 \frac{\epsilon_1}{\mu\rho} .$$

$$(ii) \ \text{Define } W(\xi) := \int_0^1 \Upsilon_{\hat{\omega}_1}(t\xi) t\xi \, dt, \text{ then } W \in \mathcal{T}_{\mu^7\rho} \text{ and } \|W\|_{\mathcal{T}_{\mu^7\rho}} \leq 2^7 \epsilon_1 .$$

Furthermore the 1-form $\alpha_W := \langle W, \cdot \rangle$ satisfies $d\alpha_W = \hat{\omega}_1 - \omega_0$.

Finally we will need also some analytic and geometric properties of the map

$$\check{\Psi} := \hat{\psi}^{-1} \circ \Psi . \quad (\text{A.30})$$

The functions $\{\check{\Psi}(\xi)\}_{j \in \mathbb{Z}}$ forms a new set of coordinates in a suitable neighborhood of the origin whose properties are given by the following corollary:

Corollary A.14. *The map $\check{\Psi} : B_r^{p,w^0}(\mu^8\rho) \rightarrow \ell_r^{p,w^0}$, defined in (A.30), satisfies the following properties:*

$$(i) \ d\check{\Psi}(0) = \mathbb{1} \text{ and } \check{\Psi}^0 := \check{\Psi} - \mathbb{1} \in \mathcal{T}_{\mu^8\rho} \text{ with } \|\check{\Psi}^0\|_{\mathcal{T}_{\mu^8\rho}} \leq 2^8 \epsilon_1 .$$

$$(ii) \ \mathcal{T}^{(0)} = \check{\Psi}(\mathcal{T}), \text{ namely the foliation defined by } \check{\Psi} \text{ coincides with the foliation defined by } \Psi .$$

$$(iii) \ \text{The functionals } \{\frac{1}{2} |\check{\Psi}_j|^2\}_{j \in \mathbb{Z}} \text{ pairwise commute with respect to the symplectic form } \omega_0 .$$

Proof. As in [BM16, Corollary 2.20]. \square

The second step consists in transforming $\hat{\omega}_1$ into the symplectic form ω_0 while preserving the functions I_l . In order to perform this transformation, we apply once more the Darboux procedure with $\Omega_1 = \hat{\omega}_1$ and $\Omega_0 = \omega_0$. However, we require each leaf of the foliation to be invariant under the transformation. In practice, we look for a change of coordinates ψ satisfying

$$\psi^* \Omega_1 = \Omega_0 , \quad (\text{A.31})$$

$$I_l(\psi(\xi)) = I_l(\xi), \quad \forall l \in \mathbb{Z} . \quad (\text{A.32})$$

In order to fulfill the second equation, we take advantage of the arbitrariness of f in equation (A.17). It turns out that if f satisfies the set of differential equations given by

$$df(X_{-I_l}^0) - (\alpha_1 - \alpha_0)(X_{-I_l}^0) = 0, \quad \forall l \in \mathbb{Z} , \quad (\text{A.33})$$

then equation (A.32) is satisfied (as it will be proved below). Here α_1 is the potential form of $\hat{\omega}_1$ and is given by $\alpha_1 := \alpha_0 + \alpha_W$, where α_W is defined in Corollary A.13. However, (A.33) is essentially a system of equations for the potential of a 1-form on a torus, so there is a solvability condition. In Lemma A.16 below we will prove that the system (A.33) has a solution if the following conditions are satisfied:

$$d(\alpha_1 - \alpha_0)|_{T\mathcal{T}^{(0)}} = 0, \quad (\text{A.34})$$

$$M(\alpha_1 - \alpha_0)|_{T\mathcal{T}^{(0)}} = 0. \quad (\text{A.35})$$

In order to show that these two conditions are fulfilled, we need a preliminary result. First, for $\xi \in \ell_r^{p,w}$ fixed, define the symplectic orthogonal of $T_\xi \mathcal{T}^{(0)}$ with respect to the form $\omega^t := \omega_0 + t(\hat{\omega}_1 - \omega_0)$ by

$$(T_\xi)^{\angle t} := \left\{ h \in \ell_r^{p,w} : \omega^t(\xi)(u, h) = 0 \quad \forall u \in T_\xi \mathcal{T}^{(0)} \right\}. \quad (\text{A.36})$$

Lemma A.15. (i) For $\xi \in B_r^{p,w^0}(\mu^5 \rho)$, one has $T_\xi \mathcal{T}^{(0)} = (T_\xi \mathcal{T}^{(0)})^{\angle t}$.

(ii) The solvability conditions (A.34), (A.35) are fulfilled.

Proof. It follows with the same arguments of [BM16, Lemma 2.21, Lemma 2.22]. \square

We show now that the system (A.33) can be solved and its solution has good analytic properties:

Lemma A.16. If conditions (A.34) and (A.35) are fulfilled, then equation (A.33) has a solution f . Moreover, denoting $h_j := (\alpha_1 - \alpha_0)(X_{-I_j}^0)$, the solution f is given by the explicit formula $f(\xi) = \sum_{k=0}^{\infty} f_k(\xi)$,

$$f_0(\xi) = L_0 h_0, \quad f_{2i-1}(\xi) = M_0 \prod_{\ell=1}^{i-1} (M_\ell M_{-\ell}) L_i h_i, \quad f_{2i}(\xi) = M_0 \prod_{\ell=1}^{i-1} (M_\ell M_{-\ell}) M_i L_{-i} h_{-i}, \quad (\text{A.37})$$

where

$$L_j g(\xi) = \frac{1}{2\pi} \int_0^{2\pi} t g(\phi_j^t(\xi)) dt, \quad \forall j \in \mathbb{Z}.$$

Finally $f \in \mathcal{N}_{\mu^7 \rho}(\ell_r^{p,w^1}, \mathbb{C})$, $\nabla f \in \mathcal{N}_{\mu^7 \rho}^T(B^{p,w^0} \cap \ell_r^{p,w^1}, \ell_r^{p,w^2})$, $f = O(\xi^{N+1})$ and

$$\|f\|_{\mu^7 \rho} \leq 2^{10} \epsilon_1 \mu^7 \rho, \quad \langle |\nabla f| \rangle_{\mu^7 \rho} \leq 2^{11} \epsilon_1. \quad (\text{A.38})$$

Proof. Denote by θ_j the time along the flow generated by $X_{-I_j}^0$, then one has $dg(X_{-I_j}^0) = \frac{\partial g}{\partial \theta_j}$, so that the equations to be solved take the form

$$\frac{\partial f}{\partial \theta_j} = h_j, \quad \forall j \in \mathbb{Z}. \quad (\text{A.39})$$

Clearly $\frac{\partial}{\partial \theta_j} M_j h_j = 0$, and by (A.34) it follows that

$$\frac{\partial}{\partial \theta_l} M_j h_j = M_j \frac{\partial h_j}{\partial \theta_l} = M_j \frac{\partial h_l}{\partial \theta_j} = \frac{\partial}{\partial \theta_j} M_j h_l = 0, \quad \forall l, j \in \mathbb{Z},$$

which shows that $M_j h_j$ is independent of all the θ 's, thus $M_j h_j = M h_j$. Furthermore, by (A.35) one has $M h_j = 0$, $\forall j \in \mathbb{Z}$. Now, using that $\frac{\partial}{\partial \theta_j} L_j g = L_j \frac{\partial}{\partial \theta_j} g = g - M_j g$, one verifies that f_i defined in (A.37) satisfies

$$\begin{aligned} \frac{\partial f_0}{\partial \theta_l} &= h_l - M_0 h_l, \\ \frac{\partial f_{2i-1}}{\partial \theta_l} &= \begin{cases} 0 & \text{if } |l| < i \\ M_0 (\prod_{\ell=1}^{i-1} M_\ell M_{-\ell}) h_l & \text{if } l = i \\ M_0 (\prod_{\ell=1}^{i-1} M_\ell M_{-\ell}) h_l - M_0 (\prod_{\ell=1}^{i-1} M_\ell M_{-\ell}) M_i h_l & \text{otherwise} \end{cases}, \\ \frac{\partial f_{2i}}{\partial \theta_l} &= \begin{cases} 0 & \text{if } |l| < i \text{ or } l = i \\ M_0 (\prod_{\ell=1}^{i-1} M_\ell M_{-\ell}) M_i h_l & \text{if } l = -i \\ M_0 \prod_{\ell=1}^{i-1} (M_\ell M_{-\ell}) M_i h_l - M_0 \prod_{\ell=1}^i (M_\ell M_{-\ell}) h_l & \text{otherwise} \end{cases}. \end{aligned}$$

Thus the series $f(\xi) = \sum_{i \geq 0} f_i(\xi)$, if convergent, satisfies (A.39).

We prove now the convergence of the series for f and ∇f . First we define, for $\theta \in \mathbb{T}^\infty$,

$$\Theta_{2i-1}^\theta := \phi_0^{\theta_0} \prod_{\ell=1}^{i-1} (\phi_\ell^{\theta_\ell} \phi_{-\ell}^{\theta_{-\ell}}) \phi_i^{\theta_i} \quad , \quad \Theta_{2i}^\theta := \phi_0^{\theta_0} \prod_{\ell=1}^i (\phi_\ell^{\theta_\ell} \phi_{-\ell}^{\theta_{-\ell}}) \quad \forall i \geq 0 \quad ,$$

then by (A.37) one has

$$f_{2i-1}(\xi) = \int_{\mathbb{T}^{2i}} \theta_i h_i(\Theta_{2i-1}^\theta \xi) d\theta^{2i} \quad , \quad f_{2i}(\xi) = \int_{\mathbb{T}^{2i+1}} \theta_{-i} h_{-i}(\Theta_{2i}^\theta \xi) d\theta^{2i+1} \quad (\text{A.40})$$

$$\nabla f_{2i-1}(\xi) = \int_{\mathbb{T}^{2i}} \Theta_{2i-1}^{-\theta} \theta_i \nabla h_i(\Theta_{2i-1}^\theta \xi) d\theta^{2i} \quad , \quad \nabla f_{2i}(\xi) = \int_{\mathbb{T}^{2i+1}} \Theta_{2i}^{-\theta} \theta_{-i} \nabla h_{-i}(\Theta_{2i}^\theta \xi) d\theta^{2i+1} \quad (\text{A.41})$$

where \mathbb{T}^{2i} (respectively \mathbb{T}^{2i+1}) is the $2i$ -dimensional torus ($2i+1$ -dimensional) and the measure $d\theta^{2i} := \frac{d\theta_0}{2\pi} (\prod_{\ell=1}^{i-1} \frac{d\theta_\ell}{2\pi} \frac{d\theta_{-\ell}}{2\pi}) \frac{d\theta_i}{2\pi}$ ($d\theta^{2i+1} = \frac{d\theta_0}{2\pi} (\prod_{\ell=1}^i \frac{d\theta_\ell}{2\pi} \frac{d\theta_{-\ell}}{2\pi})$). Now, using that

$$h_j(\xi) = \langle W(\xi), X_{-I_j}^0(\xi) \rangle = \text{Re}(-iW_j(\xi)\bar{\xi}_j) \quad , \quad \forall j \in \mathbb{Z}$$

one gets that $\forall i \geq 0$

$$\underline{f}_{2i-1}(|\xi|) \leq 2\pi \underline{h}_i(|\xi|) \leq 2\pi \underline{W}_i(|\xi|)|\xi_i| \quad , \quad \underline{f}_{2i}(|\xi|) \leq 2\pi \underline{h}_{-i}(|\xi|) \leq 2\pi \underline{W}_{-i}(|\xi|)|\xi_{-i}| \quad , \quad \forall i \geq 0$$

therefore for every $1 \leq p \leq 2$ one has

$$\underline{f}(|\xi|) \leq \sum_{k=0}^{\infty} \underline{f}_k(|\xi|) \leq 2\pi \|\underline{W}(|\xi|)\|_2 \|\xi\|_2 \leq 2\pi \|\underline{W}(|\xi|)\|_{p,w^0} \|\xi\|_{p,w^0}$$

and it follows that $|\underline{f}|_{\mu^7\rho} \leq 2\pi |\underline{W}|_{\mu^7\rho} \mu^7\rho$. This proves the convergence of the series defining f .

Consider now the gradient of h_i , whose k^{th} component is given by

$$[\nabla h_i(\xi)]_k = \text{Re} \left(-i \frac{\partial W_i(\xi)}{\partial \xi_k} \bar{\xi}_i \right) + \delta_{i,k} \text{Re}(-iW_i(\xi)) \quad .$$

Inserting the formula displayed above in (A.41) we get that ∇f_i is the sum of two terms. We begin by estimating the second one, which we denote by $(\nabla f_i)^{(2)}$. The k^{th} component ($k \in \mathbb{Z}$) of $(\nabla f)^{(2)} := \sum_l (\nabla f_l)^{(2)}$ is given by

$$\left[(\nabla f(\xi))^{(2)} \right]_k = \left[\sum_l (\nabla f_l(\xi))^{(2)} \right]_k = \begin{cases} \int_{\mathbb{T}^{2k}} \Theta_{2k-1}^{-\theta} \theta_k \text{Re}(-iW_k(\Theta_{2k-1}^\theta \xi)) d\theta^{2k} \quad , & k > 0 \\ \int_{\mathbb{T}^{2|k|+1}} \Theta_{2|k|}^{-\theta} \theta_k \text{Re}(-iW_k(\Theta_k^\theta \xi)) d\theta^{2|k|+1} \quad , & k \leq 0 \end{cases} \quad (\text{A.42})$$

thus, for any $\xi \in B_r^{p,w^0}(\mu^7\rho)$ one has $[(\nabla f(|\xi|))^{(2)}]_k \leq 2\pi \underline{W}_k(|\xi|)$, $\forall k \in \mathbb{Z}$, and therefore

$$\left\langle |(\nabla f)^{(2)}| \right\rangle_{\mu^7\rho} \leq 2\pi \langle |\underline{W}| \rangle_{\mu^7\rho} \leq \pi 2^8 \epsilon_1 \quad .$$

We come to the other term, which we denote by $(\nabla f_i)^{(1)}$. Its k^{th} component is given by

$$\begin{aligned} \left[(\nabla f_{2i-1}(\xi))^{(1)} \right]_k &= \int_{\mathbb{T}^{2i}} \Theta_{2i-1}^{-\theta} \theta_i \text{Re} \left(-i \frac{\partial W_i}{\partial \xi_k}(\Theta_{2i-1}^\theta \xi) \overline{\phi_i^{\theta_i} \xi_i} \right) d\theta^{2i} \quad , \\ \left[(\nabla f_{2i}(\xi))^{(1)} \right]_k &= \int_{\mathbb{T}^{2i}} \Theta_{2i}^{-\theta} \theta_{-i} \text{Re} \left(-i \frac{\partial W_{-i}}{\partial \xi_k}(\Theta_{2i}^\theta \xi) \overline{\phi_{-i}^{\theta_{-i}} \xi_{-i}} \right) d\theta^{2i+1} \quad . \end{aligned} \quad (\text{A.43})$$

Then

$$\begin{aligned} \left[\underline{f}_{2i-1}(|\xi|) \right]_k &\leq 2\pi \frac{\partial W_i}{\partial \xi_k}(|\xi|)|\xi_i| = 2\pi [\underline{dW}(|\xi|)]_k^i |\xi_i| \quad , \\ \left[\underline{f}_{2i}(|\xi|) \right]_k &\leq 2\pi \frac{\partial W_{-i}}{\partial \xi_k}(|\xi|)|\xi_{-i}| = 2\pi [\underline{dW}(|\xi|)]_k^{-i} |\xi_{-i}| \quad . \end{aligned}$$

It follows that the k^{th} component of the function $(\nabla f)^{(1)} := \sum_{i \geq 0} (\nabla f_i)^{(1)}$ satisfies

$$\left[(\nabla f(|\xi|))^{(1)} \right]_k \leq \left[\sum_{l \geq 0} (\nabla f_l(|\xi|))^{(1)} \right]_k \leq 2\pi \sum_{l \in \mathbb{Z}} [\underline{dW}(|\xi|)]_k^l |\xi_l| = 2\pi [\underline{dW}(|\xi|)^* |\xi|]_k .$$

Therefore $\left\langle |(\nabla f)^{(1)}| \right\rangle_{\mu_{\tau\rho}^{\tau\rho}} \leq 2\pi \|W\|_{\mathcal{S}_{\mu_{\tau\rho}^{\tau\rho}}} \leq \pi 2^8 \epsilon_1$. This is the step at which the control of the norm of the modulus \underline{dW}^* of dW^* is needed. Thus the claimed estimate for ∇f follows. \square

We can finally apply the Darboux procedure in order to construct an analytic change of coordinates ψ which satisfies (A.31) and (A.32).

Lemma A.17. *There exists a map $\psi : B_r^{p,w^0}(\mu^9\rho) \rightarrow \ell_r^{p,w^0}$ which satisfies (A.31). Moreover $\psi - \mathbb{1} \in \mathcal{N}_{\mu^9\rho}^T(B^{p,w^0} \cap \ell_r^{p,w^1}, \ell_r^{p,w^2})$, $\psi - \mathbb{1} = O(\xi^N)$ and*

$$\left\langle |\psi - \mathbb{1}| \right\rangle_{\mu^9\rho} \leq 2^{14} \epsilon_1 . \quad (\text{A.44})$$

Proof. As in [BM16, Lemma 2.24]. As anticipated just after Corollary A.14, we apply the Darboux procedure with $\Omega_0 = \omega_0$, $\Omega_1 = \hat{\omega}_1$ and f solution of (A.33). Then equation (A.17) takes the form

$$Y^t = (\mathbf{i} + t\Upsilon_{\hat{\omega}_1})^{-1}(\nabla f - W), \quad (\text{A.45})$$

where $\Upsilon_{\hat{\omega}_1}$ and W are defined in Corollary A.13. By Lemma A.16 and Corollary A.13, the vector field Y^t is of class $\mathcal{N}_{\mu^8\rho}^T(B^{p,w^0} \cap \ell_r^{p,w^1}, \ell_r^{p,w^2})$, $Y^t(\xi) = O(\xi^N)$ and

$$\sup_{t \in [0,1]} \left\langle |Y^t| \right\rangle_{\mu^8\rho} < 2(2^{11}\epsilon_1 + 2^7\epsilon_1) < 2^{13}\epsilon_1 .$$

Thus Y^t generates a flow $\psi^t : B_r^{p,w^0}(\mu^9\rho) \rightarrow \ell_r^{p,w^0}$, defined for every $t \in [0,1]$, which satisfies (cf. Lemma A.6)

$$\left\langle |\psi^t - \mathbb{1}| \right\rangle_{\mu^9\rho} \leq 2^{14}\epsilon_1, \quad \forall t \in [0,1] .$$

Thus the map $\psi := \psi^t|_{t=1}$ exists and satisfies the claimed properties. \square

Lemma A.18. *Let f be as in (A.37) and ψ^t be the flow map of the vector field Y^t defined in (A.45). Then $\forall l \geq 1$ one has $I_l(\psi^t(\xi)) = I_l(\xi)$, for each $t \in [0,1]$.*

Proof. As in [BM16, Lemma 2.25]. \square

Proof of Theorem 2.19. It follows as in [BM16]. The invertibility of $\tilde{\Psi}$ and the analytic properties stated in item v) follow by Lemma A.5 (ii). \square

B Proof of Lemma 3.8

We prove the result only for $Z_j^n(\zeta)$, since for $W_j^n(\zeta)$ the computations are analogous. We follow again the constuction of [KP10, BM16], adding more precise quantitative estimates.

To perform the Taylor expansion at every order it is convenient to proceed in the following way. Write $z_j(\zeta) = z_{j,1}(\zeta) + z_{j,2}(\zeta)$ where

$$z_{j,1}(\zeta) := ((L_0 - \lambda_j^0) f_j^-(\zeta), \imath f_j^-(\zeta))_{\mathcal{Y}}, \quad z_{j,2}(\zeta) := (V(\zeta) f_j^-(\zeta), \imath f_j^-(\zeta))_{\mathcal{Y}} . \quad (\text{B.1})$$

For $\varsigma = 1, 2$ we will write

$$z_{j,\varsigma}(\zeta) = \sum_{n=1}^{\infty} Z_{j,\varsigma}^n(\zeta) ,$$

where the $Z_{j,\zeta}^n(\zeta)$ are bounded n -homogeneous polynomials in ζ . To obtain the explicit expression for the $Z_{j,\zeta}^n$'s, we begin by expanding $f_j^+(\zeta)$ and $f_j^-(\zeta)$ in Taylor series. Since

$$f_j^\pm(\zeta) = U_j(\zeta) f_{j0}^\pm = \left(\mathbb{1} - (P_j(\zeta) - P_{j0})^2 \right)^{-1/2} \left(\mathbb{1} + (P_j(\zeta) - P_{j0}) \right) f_{j0}^\pm$$

we expand the r.h.s. above in power series of $P_j(\zeta) - P_{j0}$, getting:

$$f_j^\pm(\zeta) = \sum_{m=0}^{\infty} c_m (P_j(\zeta) - P_{j0})^m f_{j0}^\pm, \quad (\text{B.2})$$

where the c_m 's are the coefficients of the Taylor series of the function $\phi(x) = \frac{1+x}{(1-x^2)^{1/2}}$. In particular $c_{2k+1} = c_{2k} := (-1)^k \binom{-1/2}{k} \equiv \left(\frac{1}{4}\right)^k \binom{2k}{k}$. Using also that $\sqrt{2\pi n} \left(\frac{n}{e}\right)^n < n! < \sqrt{2\pi n} \left(\frac{n}{e}\right)^n e^{\frac{1}{12n}}$ one has $0 < c_m \leq 1, \forall m \geq 0$.

Now we expand $P_j(\zeta)$ in Taylor series: for ζ sufficiently small

$$P_j(\zeta) - P_{j0} = \sum_{n=1}^{\infty} (-1)^n \mathcal{P}^n(\zeta), \quad \mathcal{P}^n(\zeta) := -\frac{1}{2\pi i} \oint_{\Gamma_j} (L_0 - \lambda)^{-1} T^n(\zeta, \lambda) d\lambda \quad (\text{B.3})$$

where the Γ_j 's are defined as in (3.17), and

$$T(\zeta, \lambda) := V(\zeta) (L_0 - \lambda)^{-1}.$$

Substituting (B.3) into (B.2) we get that

$$\begin{aligned} f_j^\pm(\zeta) &= f_{j0}^\pm + \sum_{n \geq 1} \sum_{1 \leq m \leq n} c_m \sum_{\substack{\alpha = (\alpha_1, \dots, \alpha_m) \in \mathbb{N}^m \\ |\alpha| = n}} f_{j,m}^{\pm, \alpha}(\zeta), \\ f_{j,m}^{\pm, \alpha}(\zeta) &:= (-1)^{|\alpha|} \mathcal{P}^{\alpha_m}(\zeta) \circ \dots \circ \mathcal{P}^{\alpha_1}(\zeta) f_{j0}^\pm. \end{aligned} \quad (\text{B.4})$$

Consider now $Z_{j,1}^n$. It is obtained by inserting (B.4) into $z_{j,1}(\zeta)$ and collecting the terms of order n . Thus $Z_{j,1}^n(\zeta)$ equals

$$\sum_{\substack{q=(q_1, q_2) \in \mathbb{N}^2 \\ |q| \leq n}} c_{q_1} c_{q_2} \sum_{\substack{\beta=(\beta_1, \dots, \beta_{|q|}) \in \mathbb{N}^{|q|} \\ |\beta| = n}} \left((L_0 - \lambda_j^0) f_{j,q_1}^{-, (\beta_1, \dots, \beta_{q_1})}(\zeta), \iota f_{j,q_2}^{-, (\beta_{|q|}, \dots, \beta_{q_1+1})}(\zeta) \right)_{\mathcal{Y}}. \quad (\text{B.5})$$

We claim that for every $\|\zeta\|_2 < 1/8$ one has

$$\iota \mathcal{P}^n(\zeta) = \mathcal{P}^n(\zeta)^* \iota. \quad (\text{B.6})$$

Indeed by Lemma 3.5, $\Gamma_j \ni \lambda \mapsto (L_0 - \lambda)^{-1} T^n(\zeta, \lambda) \in \mathcal{L}(\mathcal{Y})$ is continuous. By (3.13)

$$\iota (L_0 - \lambda)^{-1} T^n(\zeta, \lambda) = [(L_0 - \lambda)^{-1} T^n(\zeta, \lambda)]^* \iota,$$

thus (B.6) follows by a direct computation. Using (B.6), the scalar product of (B.5) becomes

$$(-1)^n \left(\mathcal{P}^{\beta_{|q|}}(\zeta) \circ \dots \circ \mathcal{P}^{\beta_{q_1+1}}(\zeta) (L_0 - \lambda_j^0) \mathcal{P}^{\beta_{q_1}}(\zeta) \circ \dots \circ \mathcal{P}^{\beta_1}(\zeta) f_{j0}^-, \iota f_{j0}^- \right)_{\mathcal{Y}}. \quad (\text{B.7})$$

To write it explicitly remark that by (H4b)

$$(L_0 - \lambda)^{-1} f_{j0}^\pm = \frac{1}{\lambda_j^0 - \lambda} f_{j0}^\pm, \quad V(\zeta) f_{j0}^\mp = \sum_{i \in \mathbb{Z}} (V(\zeta) f_{j0}^\mp, f_{i0}^\pm)_{\mathcal{Y}} f_{i0}^\pm. \quad (\text{B.8})$$

Therefore $\forall a \in \mathbb{N}$

$$\begin{aligned} \mathcal{P}^a(\zeta) f_{j0}^- &= \\ \sum_{i_1, \dots, i_a \in \mathbb{Z}} &\left[\frac{i}{2\pi} \oint_{\Gamma_j} \frac{1}{\lambda_j^0 - \lambda} \frac{(V(\zeta) f_{j0}^-, f_{i_1 0}^+)}{\lambda_{i_1}^0 - \lambda} \frac{(V(\zeta) f_{i_1 0}^+, f_{i_2 0}^-)}{\lambda_{i_2}^0 - \lambda} \dots \frac{(V(\zeta) f_{i_{a-1} 0}^{\sigma_{a-1}}, f_{i_a 0}^{\sigma_a})}{\lambda_{i_a}^0 - \lambda} d\lambda \right] f_{i_a 0}^{\sigma_a} \end{aligned} \quad (\text{B.9})$$

where $\sigma_a = +$ if a is odd, and $\sigma_a = -$ if a is even. Using repeatedly (B.9) one gets

$$(B.7) = \sum_{i_1, \dots, i_n \in \mathbb{Z}} \mathcal{S}_{j,1}^{q,\beta}(\mathbf{i}) (V(\zeta) f_{j0}^-, f_{i_1 0}^+)_{\mathcal{Y}} (V(\zeta) f_{i_1 0}^+, f_{i_2 0}^-)_{\mathcal{Y}} \cdots (V(\zeta) f_{i_{n-1} 0}^{\sigma_{n-1}}, f_{i_n 0}^{\sigma_n})_{\mathcal{Y}} (f_{i_n 0}^{\sigma_n}, \iota f_{j0}^-)_{\mathcal{Y}},$$

with the kernel

$$\begin{aligned} \mathcal{S}_{j,1}^{q,\beta}(\mathbf{i}) &:= \left(\frac{i}{2\pi}\right)^{|q|} (-1)^n \oint_{\Gamma_j} \cdots \oint_{\Gamma_j} s_{j,1}^{q,\beta}(\mathbf{i}, \vec{\lambda}) d\lambda_1 \cdots d\lambda_{|q|}, \\ s_{j,1}^{q,\beta}(\mathbf{i}, \vec{\lambda}) &:= \prod_{m=1}^{n-1} \frac{1}{\lambda(i_m) - \mu_m} \times \prod_{\ell=1}^{|q|-1} \frac{1}{\lambda(i_{\sum_{r=1}^{\ell} \beta_r}) - \lambda_{\ell+1}} \times \frac{\lambda(i_{\beta_1+\dots+\beta_{q_1}}) - \lambda(j)}{\lambda(j) - \lambda_1} \times \frac{1}{\lambda(i_n) - \lambda_{|q|}} \end{aligned}$$

where we denoted $\lambda(a) := \lambda_a^0$, $\mathbf{i} = (i_1, \dots, i_n)$, $\vec{\lambda} = (\lambda_1, \dots, \lambda_{|q|})$, and $\mu_m = \mu_m(\lambda; q, \beta) \in \Gamma_j$.

Now remark that $\iota f_{j0}^- = f_{j0}^+$, hence $(f_{i_n 0}^{\sigma_n}, \iota f_{j0}^-)_{\mathcal{Y}} \equiv (f_{i_n 0}^{\sigma_n}, f_{j0}^+)_{\mathcal{Y}} \neq 0$ only if $i_n = j$, $\sigma_n = +$. This implies that when n is even (B.7) = 0, while when n is odd

$$(B.7) = \sum_{i_1, \dots, i_{n-1} \in \mathbb{Z}} \mathcal{S}_{j,1}^{q,\beta}(\mathbf{i}) (V(\zeta) f_{j0}^-, f_{i_1 0}^+)_{\mathcal{Y}} (V(\zeta) f_{i_1 0}^+, f_{i_2 0}^-)_{\mathcal{Y}} \cdots (V(\zeta) f_{i_{n-1} 0}^{\sigma_{n-1}}, f_{j0}^+)_{\mathcal{Y}}. \quad (B.10)$$

Altogether one has

$$\begin{aligned} Z_{j,1}^n(\zeta) &= \sum_{i_1, \dots, i_{n-1} \in \mathbb{Z}} \tilde{\mathcal{K}}_{j,1}(\mathbf{i}) (V(\zeta) f_{j0}^-, f_{i_1 0}^+)_{\mathcal{Y}} (V(\zeta) f_{i_1 0}^+, f_{i_2 0}^-)_{\mathcal{Y}} \cdots (V(\zeta) f_{i_{n-1} 0}^{\sigma_{n-1}}, f_{j0}^+)_{\mathcal{Y}}, \\ \tilde{\mathcal{K}}_{j,1}(\mathbf{i}) &:= \sum_{\substack{q=(q_1, q_2) \in \mathbb{N}^2 \\ |q| \leq n}} c_{q_1} c_{q_2} \sum_{\substack{\beta=(\beta_1, \dots, \beta_{|q|}) \in \mathbb{N}^{|q|} \\ |\beta|=n}} \mathcal{S}_{j,1}^{q,\beta}(\mathbf{i}). \end{aligned}$$

Now consider $\mathcal{S}_{j,1}^{q,\beta}(\mathbf{i})$. Recall that $\lambda_1, \dots, \lambda_{|q|} \in \Gamma_j \equiv \{\lambda \in \mathbb{C} : |\lambda - \lambda(j)| = \pi/2\}$ and that for any $\lambda \in \Gamma_j$ one has the estimate

$$4|\lambda(i) - \lambda| \geq \langle \lambda(i) - \lambda(j) \rangle \geq \langle i - j \rangle, \quad \forall i \in \mathbb{Z}, \quad \forall \lambda \in \Gamma_j;$$

this implies

$$|\mathcal{S}_{j,1}^{q,\beta}(\mathbf{i})| \leq 4^{n-1} \prod_{m=1}^{n-1} \frac{1}{\langle i_m - j \rangle} \cdot \sup_{\lambda \in \Gamma_j} \left| \frac{\lambda(i_{\beta_1+\dots+\beta_{q_1}}) - \lambda(j)}{\lambda(i_{\beta_1+\dots+\beta_{q_1}}) - \lambda} \right| \leq 4^n \prod_{m=1}^{n-1} \frac{1}{\langle i_m - j \rangle}.$$

Since the coefficients $c_q \leq 1 \forall q \in \mathbb{N}$, it follows that²

$$\begin{aligned} |\tilde{\mathcal{K}}_{j,1}^n(\mathbf{i})| &\leq 4^n \prod_{m=1}^{n-1} \frac{1}{\langle i_m - j \rangle} \sum_{\substack{q=(q_1, q_2) \in \mathbb{N}^2 \\ |q| \leq n}} c_{q_1} c_{q_2} \sum_{\substack{\beta=(\beta_1, \dots, \beta_{|q|}) \in \mathbb{N}^{|q|} \\ |\beta|=n}} 1 \\ &\leq 4^n \prod_{m=1}^{n-1} \frac{1}{\langle i_m - j \rangle} \sum_{r=1}^n \sum_{\substack{q=(q_1, q_2) \in \mathbb{N}^2 \\ |q|=r}} \sum_{\substack{\beta=(\beta_1, \dots, \beta_{|q|}) \in \mathbb{N}^{|q|} \\ |\beta|=n}} 1 \\ &\leq 4^n \prod_{m=1}^{n-1} \frac{1}{\langle i_m - j \rangle} \sum_{r=1}^n \binom{r-1}{1} \binom{n-1}{r-1} \leq 4^n \prod_{m=1}^{n-1} \frac{1}{\langle i_m - j \rangle} \sum_{r=0}^{n-1} r \binom{n-1}{r} \\ &\leq 4^n \prod_{m=1}^{n-1} \frac{1}{\langle i_m - j \rangle} (n-1) 2^{n-2} \leq 16^{n-1} \prod_{m=1}^{n-1} \frac{1}{\langle i_m - j \rangle}, \end{aligned} \quad (B.11)$$

Similar computations can be performed for $Z_{j,2}^n(\zeta)$ and one finds

$$Z_{j,2}^n(\zeta) = \sum_{i_1, \dots, i_{n-1} \in \mathbb{Z}} \tilde{\mathcal{K}}_{j,2}^n(\mathbf{i}) (V(\zeta) f_{j0}^-, f_{i_1 0}^+)_{\mathcal{Y}} (V(\zeta) f_{i_1 0}^+, f_{i_2 0}^-)_{\mathcal{Y}} \cdots (V(\zeta) f_{i_{n-1} 0}^{\sigma_{n-1}}, f_{j0}^+)_{\mathcal{Y}}$$

²recall that any $n \in \mathbb{N}$ can be written as a sum of k positive integers in $\binom{n-1}{k-1}$ possible ways.

where the kernel $\tilde{\mathcal{K}}_{j,2}(\mathbf{i})$ fulfills the same bound (B.11). In such a way we proved that

$$Z_j^n(\zeta) = \sum_{i_1, \dots, i_{n-1} \in \mathbb{Z}} \tilde{\mathcal{K}}_j^n(\mathbf{i}) (V(\zeta) f_{j0}^-, f_{i_1 0}^+)_{\mathcal{Y}} (V(\zeta) f_{i_1 0}^+, f_{i_2 0}^-)_{\mathcal{Y}} \cdots (V(\zeta) f_{i_{n-1} 0}^-, f_{j0}^+)_{\mathcal{Y}} \quad (\text{B.12})$$

with

$$|\tilde{\mathcal{K}}_j^n(\mathbf{i})| \leq 2 \cdot 16^{n-1} \prod_{m=1}^{n-1} \frac{1}{\langle i_m - j \rangle}. \quad (\text{B.13})$$

Now remark that by (H4b), $j + i_1$, $i_a + i_{a+1}$ and $i_{n-1} + j$ must be even, so define k_1, \dots, k_n by

$$2k_1 = i_1 + j, \quad 2k_\ell = (-1)^{\ell+1} (i_\ell + i_{\ell-1}), \quad 2k_n = i_{n-1} + j.$$

In this way one has

$$j = k_1 + \dots + k_n \quad (\text{B.14})$$

and by (H4b), for any $1 \leq m \leq n$, m odd

$$(V(\zeta) f_{j0}^-, f_{i_1 0}^+)_{\mathcal{Y}} = \xi_{k_1}, \quad (V(\zeta) f_{i_{m-1} 0}^-, f_{i_m 0}^+)_{\mathcal{Y}} = \xi_{k_m}, \quad (V(\zeta) f_{i_m 0}^+, f_{i_{m+1} 0}^-)_{\mathcal{Y}} = \eta_{-k_{m+1}}$$

so that (B.12) becomes

$$Z_j^n(\zeta) = \sum_{k_1 + k_2 + \dots + k_n = j} \mathcal{K}_j^n(k_1, \dots, k_n) \xi_{k_1} \eta_{-k_2} \cdots \xi_{k_n}$$

with kernel

$$\mathcal{K}_j^n(k_1, \dots, k_n) := \tilde{\mathcal{K}}_j^n((2k_1 - j), -(2k_2 + 2k_1 - j), \dots, -(2k_{n-1} + \dots + 2k_1 - j)).$$

Such kernel is supported in (B.14) and by the estimate (B.13) it fulfills (3.30).

C Technical results

First we write explicitly $\mathfrak{g}_{n,r}$ defined in (3.2). Let $1 \leq r \leq n$ be integers. For r odd we have

$$\begin{aligned} \mathfrak{g}_{n,r}(\mathbf{k}; j) &\equiv \mathbb{1}_{\mathfrak{S}_{-j}^{n,r}} \prod_{\substack{m=1 \\ m \text{ odd}}}^{r-2} \left\langle \sum_{\ell=1}^m k_\ell - k_r \right\rangle^{-1} \left\langle \sum_{\ell=1}^{m+1} k_\ell \right\rangle^{-1} \\ &\times \left\langle \sum_{\ell=1}^{r-1} k_\ell + j - k_r \right\rangle^{-1} \left\langle \sum_{\ell=1}^{r-1} k_\ell + j + k_{r+1} \right\rangle^{-1}, \\ &\times \prod_{\substack{m=r+1 \\ m \text{ odd}}}^{n-1} \left\langle \sum_{\substack{\ell=1 \\ \ell \neq r}}^m k_\ell + j - k_r \right\rangle^{-1} \left\langle \sum_{\substack{\ell=1 \\ \ell \neq r}}^{m+1} k_\ell + j \right\rangle^{-1} \end{aligned} \quad (\text{C.1})$$

where $\mathfrak{S}_a^{n,r}$ is defined in (3.41). For r even

$$\begin{aligned} \mathfrak{g}_{n,r}(\mathbf{k}; j) &\equiv \mathbb{1}_{\mathfrak{S}_{-j}^{n,r}} \prod_{\substack{m=1 \\ m \text{ odd}}}^{r-3} \left\langle \sum_{\ell=1}^m k_\ell - k_r \right\rangle^{-1} \left\langle \sum_{\ell=1}^{m+1} k_\ell \right\rangle^{-1} \\ &\times \left\langle \sum_{\ell=1}^{r-1} k_\ell - k_r \right\rangle^{-1} \left\langle \sum_{\ell=1}^{r-1} k_\ell + j \right\rangle^{-1} \\ &\times \prod_{\substack{m=r+1 \\ m \text{ odd}}}^{n-1} \left\langle \sum_{\substack{\ell=1 \\ \ell \neq r}}^m k_\ell + j - k_r \right\rangle^{-1} \left\langle \sum_{\substack{\ell=1 \\ \ell \neq r}}^{m+1} k_\ell + j \right\rangle^{-1}. \end{aligned} \quad (\text{C.2})$$

Example C.1. For $n = 3$, one has

$$\begin{aligned} f_3(k_1, k_2, k_3; j) &= \mathbb{1}_{\{k_1+k_2+k_3=j\}} (\langle k_1 - j \rangle \langle k_1 + k_2 \rangle)^{-1}, \\ g_{3,1}(k_1, k_2, k_3; j) &= \mathbb{1}_{\{-k_1+k_2+k_3=-j\}} (\langle j - k_1 \rangle \langle j + k_2 \rangle)^{-1}, \\ g_{3,2}(k_1, k_2, k_3; j) &= \mathbb{1}_{\{k_1-k_2+k_3=-j\}} (\langle k_1 - k_2 \rangle \langle k_1 + j \rangle)^{-1}, \\ g_{3,3}(k_1, k_2, k_3; j) &= \mathbb{1}_{\{k_1+k_2-k_3=-j\}} (\langle k_1 - k_3 \rangle \langle k_1 + k_2 \rangle)^{-1}, \end{aligned}$$

while for $n = 5$

$$\begin{aligned} f_5(k_1, k_2, k_3, k_4, k_5; j) &= \frac{\mathbb{1}_{\{k_1+k_2+k_3+k_4+k_5=j\}}}{\langle k_1 - j \rangle \langle k_1 + k_2 \rangle \langle k_1 + k_2 + k_3 - j \rangle \langle k_1 + k_2 + k_3 + k_4 \rangle} \\ g_{5,1}(k_1, k_2, k_3, k_4, k_5; j) &= \frac{\mathbb{1}_{\{-k_1+k_2+k_3+k_4+k_5=-j\}}}{\langle j - k_1 \rangle \langle j + k_2 \rangle \langle j + k_2 + k_3 - k_1 \rangle \langle j + k_2 + k_3 + k_4 \rangle} \\ g_{5,2}(k_1, k_2, k_3, k_4, k_5; j) &= \frac{\mathbb{1}_{\{k_1-k_2+k_3+k_4+k_5=-j\}}}{\langle k_1 - k_2 \rangle \langle k_1 + j \rangle \langle k_1 + j + k_3 - k_2 \rangle \langle k_1 + j + k_3 + k_4 \rangle} \\ g_{5,3}(k_1, k_2, k_3, k_4, k_5; j) &= \frac{\mathbb{1}_{\{k_1+k_2-k_3+k_4+k_5=-j\}}}{\langle k_1 - k_3 \rangle \langle k_1 + k_2 \rangle \langle k_1 + k_2 + j - k_3 \rangle \langle k_1 + k_2 + j + k_4 \rangle} \\ g_{5,4}(k_1, k_2, k_3, k_4, k_5; j) &= \frac{\mathbb{1}_{\{k_1+k_2+k_3-k_4+k_5=-j\}}}{\langle k_1 - k_4 \rangle \langle k_1 + k_2 \rangle \langle k_1 + k_2 + k_3 - k_4 \rangle \langle k_1 + k_2 + k_3 + j \rangle} \\ g_{5,5}(k_1, k_2, k_3, k_4, k_5; j) &= \frac{\mathbb{1}_{\{k_1+k_2+k_3+k_4-k_5=-j\}}}{\langle k_1 - k_5 \rangle \langle k_1 + k_2 \rangle \langle k_1 + k_2 + k_3 - k_5 \rangle \langle k_1 + k_2 + k_3 + k_4 \rangle} \end{aligned}$$

C.1 Proof of Lemma 3.1

Verification of (3.3). We consider $1 < p \leq 2$ and $p = 1$ separately.

Case $1 < p \leq 2$. The quantity that we have to bound is the p' root of

$$\sum_{k_1+\dots+k_n=j} \prod_{\substack{m=1 \\ m \text{ odd}}}^{n-1} \frac{1}{\langle k_1 + \dots + k_m - j \rangle^{p'}} \cdot \frac{1}{\langle k_1 + \dots + k_{m+1} \rangle^{p'}}. \quad (\text{C.3})$$

For any $j \in \mathbb{Z}$ we have (recall that n is odd, thus $m \in \{1, 3, \dots, n-2\}$)

$$\begin{aligned} (\text{C.3}) &\leq \sum_{k_1} \frac{1}{\langle k_1 - j \rangle^{p'}} \sum_{k_2} \frac{1}{\langle k_1 + k_2 \rangle^{p'}} \cdots \sum_{k_{n-2}} \frac{1}{\langle k_1 + \dots + k_{n-2} - j \rangle^{p'}} \sum_{k_{n-1}} \frac{1}{\langle k_1 + \dots + k_{n-2} + k_{n-1} \rangle^{p'}} \\ &\leq \left(\sum_k \frac{1}{\langle k \rangle^{p'}} \right)^{n-1} \equiv R_*^{p'(n-1)}. \end{aligned}$$

Thus (3.3) follows.

Case $p = 1$. In this case the quantity to estimate is

$$\sup_{k_1+\dots+k_n=j} \prod_{\substack{m=1 \\ m \text{ odd}}}^{n-1} \frac{1}{\langle k_1 + \dots + k_m - j \rangle} \cdot \frac{1}{\langle k_1 + \dots + k_{m+1} \rangle}$$

which is trivially majorated by 1, thus (3.3) follows.

Verification of (3.4). Case $1 < p \leq 2$. To majorate (3.4), it is sufficient to remark that any index k_1, \dots, k_{n-1} appears at least once in one term of $\mathbf{g}_{n,r}(k_1, \dots, k_n; j)$, as one verifies inspecting formulas (C.1), (C.2). Then one gets again that $\|\mathbf{g}_{n,r}(\cdot; j)\|_{\ell^{p'}(\mathbb{Z}^n)}^{p'} \leq \left(\sum_k \frac{1}{\langle k \rangle^{p'}} \right)^{n-1} \equiv R_*^{p'(n-1)}$ and (3.4) is fulfilled.

Case $p = 1$. The argument is similar to the previous case, and we skip it.

C.2 Proof of Proposition 3.2 (i)

Verification of (3.5). We treat separately the case $1 < p \leq 2$ and $p = 1$.

Case $1 < p \leq 2$. For every $j \in \mathbb{Z}$, $a \geq 0$, $0 < b \leq 1$ the l.h.s. of (3.5) is the p^{th} -root of the supremum over j of

$$\langle j \rangle^{sp'} e^{p'a|j|^b} \sum_{k_1+\dots+k_n=j} \frac{1}{n^{p'}} \frac{1}{\prod_{i=1}^n \langle k_i \rangle^{sp'} e^{p'a|k_i|^b}} \prod_{\substack{m=1 \\ m \text{ odd}}}^{n-1} \frac{1}{\langle k_1 + \dots + k_m - j \rangle^{p'}} \cdot \frac{1}{\langle k_1 + \dots + k_{m+1} \rangle^{p'}}$$

which is majorated by (C.3) since for any $j = k_1 + \dots + k_n$

$$e^{p'a|j|^b} \leq e^{p'a|k_1|^b} \dots e^{p'a|k_n|^b}, \quad \langle j \rangle^{sp'} \leq \langle k_1 \rangle^{sp'} \dots \langle k_n \rangle^{sp'}. \quad (\text{C.4})$$

Then the result follows from Lemma 3.1, with $R_0 = 1$, $R_1 = R_*$.

Case $p = 1$. In this case the quantity to estimate is

$$\langle j \rangle^s e^{a|j|^b} \sup_{k_1+\dots+k_n=j} \frac{1}{n} \frac{1}{\prod_{i=1}^n \langle k_i \rangle^s e^{a|k_i|^b}} \prod_{\substack{m=1 \\ m \text{ odd}}}^{n-1} \frac{1}{\langle k_1 + \dots + k_m - j \rangle} \cdot \frac{1}{\langle k_1 + \dots + k_{m+1} \rangle}$$

which is trivially majorated by 1, using (C.4). Thus (3.5) holds with $R_0 = R_1 = 1$.

Verification of (3.6). Using (C.4), one is brought back to estimate (3.4).

C.3 Proof of Proposition 3.2 (ii)

Verification of (3.5). As in the previous case we treat two cases:

Case $1 < p \leq 2$. We must estimate the p' -root of

$$\langle j \rangle^{p's} \sum_{k_1+\dots+k_n=j} \frac{1}{(\sum_{l=1}^n \langle k_l \rangle^s)^{p'}} \prod_{\substack{m=1 \\ m \text{ odd}}}^{n-1} \frac{1}{\langle k_1 + \dots + k_m - j \rangle^{p'}} \cdot \frac{1}{\langle k_1 + \dots + k_{m+1} \rangle^{p'}} \quad (\text{C.5})$$

For $j = k_1 + \dots + k_n$, one has the inequalities

$$\langle j \rangle^a \leq \left(\sum_{l=1}^n \langle k_l \rangle \right)^a \leq n^{a-1} \sum_{l=1}^n \langle k_l \rangle^a, \quad \forall a \geq 1, \quad (\text{C.6})$$

which yields

$$\begin{aligned} (\text{C.5}) &\leq n^{p's-1} \sum_{k_1+\dots+k_n=j} \prod_{\substack{m=1 \\ m \text{ odd}}}^{n-1} \frac{1}{\langle k_1 + \dots + k_m - j \rangle^{p'}} \cdot \frac{1}{\langle k_1 + \dots + k_{m+1} \rangle^{p'}} \\ &\leq (2^{p's-1})^{n-1} \cdot (\text{C.3}) \leq \left(2^{p's} \sum_k \frac{1}{\langle k \rangle^{p'}} \right)^{n-1} \end{aligned}$$

where in the last line we used that $\forall a > 1, \forall 3 \leq n \in \mathbb{N}$, one has $n^a \leq 2^{a(n-1)}$. Thus (3.5) is fulfilled with $R_0 = 1$, $R_1 = 2^s \left(\sum_{k \in \mathbb{Z}} \frac{1}{\langle k \rangle^{p'}} \right)^{1/p'}$.

Case $p = 1$. One has to bound the quantity

$$\langle j \rangle^s \sup_{k_1+\dots+k_n=j} \frac{1}{\sum_{l=1}^n \langle k_l \rangle^s} \prod_{\substack{m=1 \\ m \text{ odd}}}^{n-1} \frac{1}{\langle k_1 + \dots + k_m - j \rangle} \cdot \frac{1}{\langle k_1 + \dots + k_{m+1} \rangle}; \quad (\text{C.7})$$

using (C.6) one gets the desired bound easily.

Verification of (3.6). One uses inequality (C.6) and proceed as in the proof of Proposition 3.2(i); in turn (3.6) is fulfilled with $R_0 = 1$, $R_1 = 2^s \left(\frac{1}{\langle k \rangle^{p'}} \right)^{1/p'}$.

C.4 Proof of Proposition 3.2 (iii)

Verification of (3.5). As above we treat two cases:

Case $1 < p \leq 2$. We must estimate the p' -root of the supremum over j of

$$\langle j \rangle^{p'(s+1)} \sum_{k_1+\dots+k_n=j} \frac{1}{\left(\sum_{l=1}^n \langle k_l \rangle^s \prod_{m \neq l} \langle k_m \rangle \right)^{p'}} \prod_{\substack{m=1 \\ m \text{ odd}}}^{n-1} \frac{1}{\langle k_1 + \dots + k_m - j \rangle^{p'}} \cdot \frac{1}{\langle k_1 + \dots + k_{m+1} \rangle^{p'}} \quad (\text{C.8})$$

Using (C.6), expression (C.8) is majorated by $n^{p'(s+1)-1} \sum_{l=1}^n \mathcal{M}_{l,n}$,

$$\mathcal{M}_{l,n} := \sum_{k_1+\dots+k_n=j} \frac{\langle k_l \rangle^{p'}}{\prod_{m \neq l} \langle k_m \rangle^{p'}} \prod_{\substack{m=1 \\ m \text{ odd}}}^{n-1} \frac{1}{\langle k_1 + \dots + k_m - j \rangle^{p'}} \cdot \frac{1}{\langle k_1 + \dots + k_{m+1} \rangle^{p'}} \quad (\text{C.9})$$

We claim that $\forall 1 \leq l \leq n$

$$\mathcal{M}_{l,n} \leq R_b^{n-1}, \quad R_b := 2^{p'} R_*^{p'}. \quad (\text{C.10})$$

To prove (C.10) we consider separately the case l even and l odd; if l is even then

$$\mathcal{M}_{l,n} \leq \langle k_l \rangle^{p'} \sum_{k_m: m \neq l} \frac{1}{\prod_{m \neq l} \langle k_m \rangle^{p'}} \cdot \frac{1}{\langle k_1 + \dots + k_{l-1} + k_l \rangle^{p'}} \leq R_b^{n-1},$$

which follows using repeatedly the inequalities

$$\sum_{k \in \mathbb{Z}} \frac{1}{\langle k \rangle^{p'} \langle k - j \rangle^{p'}} \leq \frac{R_b}{\langle j \rangle^{p'}}, \quad \sum_{k \in \mathbb{Z}} \frac{1}{\langle k \rangle^{p'}} < R_b. \quad (\text{C.11})$$

Similarly, if l is odd and $l \neq n$, then

$$\mathcal{M}_{l,n} \leq \langle k_l \rangle^{p'} \sum_{k_m: m \neq l} \frac{1}{\prod_{m \neq l} \langle k_m \rangle^{p'}} \cdot \frac{1}{\langle k_1 + \dots + k_l + k_{l+1} \rangle^{p'}} \leq R_b^{n-1},$$

where once again we used (C.11) iteratively. Finally consider the case $l = n$: using that $j - k_1 = k_2 + \dots + k_n$,

$$\mathcal{M}_{n,n} \leq \langle k_n \rangle^{p'} \sum_{k_m: m \neq n} \frac{1}{\prod_{m \neq n} \langle k_m \rangle^{p'}} \cdot \frac{1}{\langle k_2 + \dots + k_{n-1} + k_n \rangle^{p'}} \leq R_b^{n-1}.$$

All together we proved (C.10), consequently (C.8) $\leq n^{p'(s+1)-1} \sum_{l=1}^n \mathcal{M}_{l,n} \leq (2^{s+2} R_*)^{p'(n-1)}$. Thus (3.5) holds with $R_0 = 1$, $R_1 = 2^{s+2} R_*$.

Case $p = 1$. One has to bound

$$\langle j \rangle^{s+1} \sup_{k_1+\dots+k_n=j} \left| \frac{1}{\sum_{l=1}^n \langle k_l \rangle^s \prod_{m \neq l} \langle k_m \rangle} \prod_{\substack{m=1 \\ m \text{ odd}}}^{n-1} \frac{1}{\langle k_1 + \dots + k_m - j \rangle} \cdot \frac{1}{\langle k_1 + \dots + k_{m+1} \rangle} \right|, \quad (\text{C.12})$$

which is majorated by $n^s \sum_{l=1}^n \widetilde{\mathcal{M}}_{l,n}$,

$$\widetilde{\mathcal{M}}_{l,n} := \sup_{k_1+\dots+k_n=j} \left| \frac{\langle k_l \rangle}{\prod_{m \neq l} \langle k_m \rangle} \prod_{\substack{m=1 \\ m \text{ odd}}}^{n-1} \frac{1}{\langle k_1 + \dots + k_m - j \rangle} \cdot \frac{1}{\langle k_1 + \dots + k_{m+1} \rangle} \right|. \quad (\text{C.13})$$

We claim that $\forall 1 \leq l \leq n$

$$\widetilde{\mathcal{M}}_{l,n} \leq 1 ; \quad (\text{C.14})$$

this is proved exactly as (C.10) using iteratively the inequality

$$\sup_{k \in \mathbb{Z}} \frac{1}{\langle k \rangle \langle k-j \rangle} \leq \frac{1}{\langle j \rangle} \quad (\text{C.15})$$

in place of (C.11). Thus (3.5) holds with $R_0 = 1$, $R_1 = 2^{s+1}$.

Verification of (3.6). We consider two cases.

Case $1 < p \leq 2$. The quantity that we must estimate is the p^{th} -root of

$$\langle j \rangle^{p'(s+1)} \sum_{\mathbf{k} \in \mathfrak{S}_{-j}^{n,r}} \frac{1}{\left(\sum_{l=1}^n \langle k_l \rangle^s \prod_{m \neq l} \langle k_m \rangle \right)^{p'}} \mathfrak{g}_{n,r}(k_1, \dots, k_n; j)^{p'} \quad (\text{C.16})$$

for every possible choice of $1 \leq r \leq n$ and $n \geq 3$, n odd. First remark that (C.16) is majorated by $n^{p'(s+1)-1} \sum_{l=1}^n \mathcal{N}_{n,r}^l$,

$$\mathcal{N}_{n,r}^l := \sum_{\mathbf{k} \in \mathfrak{S}_{-j}^{n,r}} \frac{\langle k_l \rangle^{p'}}{\prod_{m \neq l} \langle k_m \rangle^{p'}} \mathfrak{g}_{n,r}(k_1, \dots, k_n; j)^{p'} . \quad (\text{C.17})$$

Once again we claim that $\forall 1 \leq l, r \leq n$,

$$\mathcal{N}_{n,r}^l \leq R_b^{n-1} . \quad (\text{C.18})$$

First let r be odd. In this case $\mathfrak{g}_{n,r}$ is given by (C.1). If $l \leq r-1$, then we have

$$\mathcal{N}_{n,r}^l \leq \langle k_l \rangle^{p'} \sum_{k_i : i \neq l} \frac{1}{\prod_{i \neq l} \langle k_i \rangle^{p'}} \cdot \frac{1}{\left\langle \sum_{\ell=1}^{r-1} k_\ell \right\rangle^{p'}} \leq R_b^{n-1} ,$$

using estimates (C.11) iteratively. If $l = r$, the term $\sum_{\ell=1}^{r-1} k_\ell + j + k_{r+1}$ equals $-\sum_{\ell=r+2}^n k_\ell + k_r$ (due to the condition $\mathbf{k} \in \mathfrak{S}_{-j}^{n,r}$), thus using again (C.11)

$$\mathcal{N}_{n,r}^r \leq \langle k_r \rangle^{p'} \sum_{k_i : i \neq r} \frac{1}{\prod_{i \neq r} \langle k_i \rangle^{p'}} \cdot \frac{1}{\left\langle \sum_{\ell=r+2}^n k_\ell - k_r \right\rangle^{p'}} \leq R_b^{n-1} .$$

Finally if $l \geq r+1$, we use that the term $\sum_{\ell=1}^{r-1} k_\ell + j - k_r$ equals $-\sum_{\ell=r+1}^n k_\ell$ (due to the condition $\mathbf{k} \in \mathfrak{S}_{-j}^{n,r}$), thus again (C.11)

$$\mathcal{N}_{n,r}^l \leq \langle k_l \rangle^{p'} \sum_{k_i : i \neq l} \frac{1}{\prod_{i \neq l} \langle k_i \rangle^{p'}} \cdot \frac{1}{\left\langle \sum_{\ell=r+1}^n k_\ell \right\rangle^{p'}} \leq R_b^{n-1} .$$

Now take r even; the relevant formula for $\mathfrak{g}_{n,r}$ is (C.2). In case $l \leq r-1$ we have that

$$\mathcal{N}_{n,r}^l \leq \langle k_l \rangle^{p'} \sum_{k_i : i \neq l} \frac{1}{\prod_{i \neq l} \langle k_i \rangle^{p'}} \cdot \frac{1}{\left\langle \sum_{\ell=1}^{r-1} k_\ell - k_r \right\rangle^{p'}} \leq R_b^{n-1} .$$

In case $l \geq r$ we use that the term $\sum_{\ell=1}^{r-1} k_\ell + j$ equals $k_r - \sum_{\ell=r+1}^n k_\ell$ (due to the condition $\mathbf{k} \in \mathfrak{S}_{-j}^{n,r}$), we have that

$$\mathcal{N}_{n,r}^l \leq \langle k_l \rangle^{p'} \sum_{k_i : i \neq l} \frac{1}{\prod_{i \neq r} \langle k_i \rangle^{p'}} \cdot \frac{1}{\left\langle k_r - \sum_{\ell=r+1}^n k_\ell \right\rangle^{p'}} \leq R_b^{n-1} .$$

All together we have proved (C.18), thus

$$(C.16) \leq n^{p'(s+1)-1} \sum_{l=1}^n \mathcal{N}_{n,r}^l \leq (2^{p'(s+1)} R_b)^{n-1} \leq (2^{s+2} R_*)^{p'(n-1)},$$

and (3.6) follows with $R_0 = 1$, $R_1 = 2^{s+2} R_*$.

Case $p = 1$. One proceeds as in the case $1 < p \leq 2$ treating separately r odd and even. One verifies that (3.6) holds with $R_0 = 1$, $R_1 = 2^{s+2}$.

References

- [BG06] D. Bambusi and B. Grébert. Birkhoff normal form for partial differential equations with tame modulus. *Duke Math. J.*, 135(3):507–567, 2006.
- [BGMR1] D. Bambusi, B. Grébert, A. Maspero and D. Robert. Reducibility of the quantum harmonic oscillator in d -dimensions with polynomial time-dependent perturbation. *Anal. PDE*, 11(3), 775–799, 2018. DOI: 10.2140/apde.2018.11.775
- [BGMR2] D. Bambusi, B. Grébert, A. Maspero and D. Robert. Growth of Sobolev norms for abstract linear Schrödinger Equations. *ArXiv e-print*, arXiv:1706.09708, 2017.
- [BM16] D. Bambusi and A. Maspero. Birkhoff coordinates for the Toda Lattice in the limit of infinitely many particles with an application to FPU. *Jour. Funct. Anal.*, 270(5), 1818–1887, 2016.
- [BM16b] D. Bambusi and A. Maspero. Freezing of energy of a soliton in an external potential. *Commun. Math. Phys.* 344(1): 155–191, 2016.
- [BGGK] D. Bättig, B. Grébert, J.-C. Guillot and T. Kappeler. Foliation of the phase space for the cubic nonlinear Schrödinger equation. *Comp. Math*, 85: 163–199, 1993.
- [BBP10] M. Berti, P. Bolle, M. Procesi. An abstract Nash-Moser theorem with parameters and applications to PDEs. *Ann. I. H. Poincaré – AN*, 27: 377 – 399, 2010.
- [BBP13] M. Berti, L. Biasco and M. Procesi. KAM theory for the Hamiltonian derivative wave equation. *Annales Scientifiques de l’ENS*, 46(2), 299–371, 2013.
- [Bou] J. Bourgain. Global solutions of nonlinear Schrödinger equations. American Mathematical Society Colloquium Publications, vol. 46. American Mathematical Society, Providence, RI, 1999.
- [CLY] H. Cong, J. Liu, and X. Yuan. Stability of KAM tori for nonlinear Schrödinger equation. *Mem. Amer. Math. Soc.*, 239 no. 1134, 2016.
- [Eli90] H. Eliasson. Normal forms for Hamiltonian systems with Poisson commuting integrals—elliptic case. *Comment. Math. Helv.*, 65(1):4–35, 1990.
- [EK09] H. Eliasson and S. Kuksin. On reducibility of Schrödinger equations with quasiperiodic in time potentials. *Comm. Math. Phys.*, 286(1): 125–135, 2009.
- [GG10] P. Gérard and S. Grellier. The cubic Szegő equation. *Ann. Sci. Ec. Norm. Supér. (4)*, 43(5):761–810, 2010.
- [GG12] P. Gérard and S. Grellier. Invariant tori for the cubic Szegő equation. *Invent. Math.*, 187(3):707—754, 2012.
- [GG17] P. Gérard and S. Grellier. The cubic Szegő equation and Hankel operators. *Astérisque*, 389, 2017.

- [GK14] B. Grébert and T. Kappeler. The defocusing NLS equation and its normal form. *European Math. Soc.*, 2014.
- [Kat66] T. Kato. *Perturbation theory for linear operators*. Die Grundlehren der mathematischen Wissenschaften, Band 132. Springer-Verlag New York, Inc., New York, 1966.
- [KMMT] T. Kappeler, A. Maspero, J. Molnar, and P. Topalov. On the convexity of the KdV Hamiltonian, *Comm. Math. Phys.*, 346(1): 191–236, 2016.
- [KM16] T. Kappeler and R. Montalto. Canonical coordinates with tame estimates for the defocusing NLS equation on the circle. *Int Math Res Notices* 2016. DOI: 10.1093/imrn/rnw233
- [KP09] T. Kappeler and J. Pöschel. On the periodic KdV equation in weighted Sobolev spaces. *Ann. I. H. Poincaré* AN26: 841–853, 2009.
- [KST13] T. Kappeler, B. Schaad, and P. Topalov. Qualitative features of periodic solutions of KdV. *Comm. Partial Differential Equations*, 38(9):1626–1673, 2013.
- [KSTa] T. Kappeler, B. Schaad, and P. Topalov. Semi-linearity of the nonlinear Fourier transform of the defocusing NLS equation. *Int. Math. Res. Notices*, (2016), DOI: 10.1093/imrn/rnv397.
- [KSTb] T. Kappeler, B. Schaad and P. Topalov. Scattering-like phenomena of the periodic defocusing NLS equation. *Arxiv e-print* arXiv:1505.07394, 2016.
- [KP10] S. Kuksin and G. Perelman. Vey theorem in infinite dimensions and its application to KdV. *Discrete Contin. Dyn. Syst.*, 27(1):1–24, 2010.
- [MP17] A. Maspero and M. Procesi. Long time stability of small finite gap solutions of the cubic Nonlinear Schrödinger equation on \mathbb{T}^2 . *ArXiv e-prints*, arXiv:1710.08168, 2017.
- [MR17] A. Maspero and D. Robert. On time dependent Schrödinger equations: global well-posedness and growth of Sobolev norms. *Jour. Funct. Anal.*, 273(2):721–781, 2017.
- [MS16] A. Maspero and B. Schaad. One smoothing property of the scattering map of the KdV on \mathbb{R} . *Discrete Contin. Dyn. Syst.*, 36(3):1493–1537, 2016.
- [Mo15] J. Molnar. New estimates of the nonlinear Fourier transform for the defocusing NLS equation. *Int. Math. Res. Not.*, 2015(17), 8309–8352, 2015.
- [Mo16] J. Molnar. Features of the nonlinear Fourier transform for the dNLS equation. PhD Thesis, University of Zurich, 2016. <http://janbernloehr.de/Download/fs16/diss>
- [Muj86] J. Mujica. *Complex analysis in Banach spaces*, volume 120 of *North-Holland Mathematics Studies*. North-Holland Publishing Co., Amsterdam, 1986.
- [Nik86] N. Nikolenko. The method of Poincaré normal forms in problems of integrability of equations of evolution type. *Uspekhi Mat. Nauk*, 41(5(251)):109–152, 263, 1986.
- [PTV] F. Planchon, N. Tzvetkov and N. Visciglia. On the growth of Sobolev norms for NLS on 2- and 3-dimensional manifolds. *Anal. PDE*, 10(5), 1123–1147, 2017.
- [Soh97] V. Sohinger. Bounds on the growth of high Sobolev norms of solutions to non-linear Schrödinger equations on S^1 . *Differential and Integral Equations*, 24: 653–718, 2011.
- [Vey78] J. Vey. Sur certains systèmes dynamiques séparables. *Amer. J. Math.*, 100(3):591–614, 1978.
- [ZM74] V. Zaharov and S. Manakov. The complete integrability of the nonlinear Schrödinger equation. *Teoret. Mat. Fiz.*, **19** (1974), 332–343.
- [ZS71] V. Zakharov and A. Shabat. Exact theory of two-dimensional self-focusing and one-dimensional self-modulation of waves in nonlinear media. *Ž. Èksper. Teoret. Fiz.*, **61** (1971), 118–134.